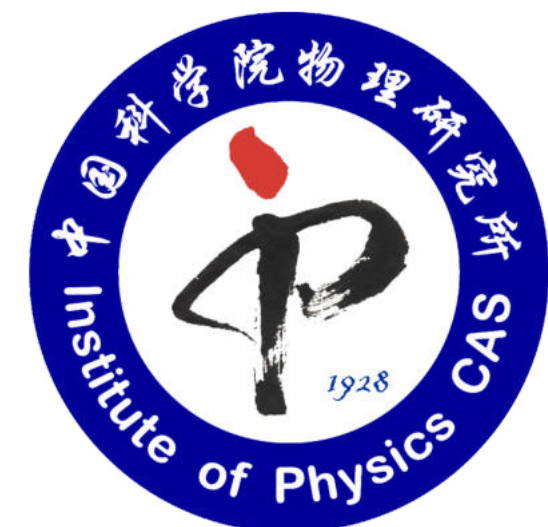


A physicist's perspective on generative models

Lei Wang (王磊)

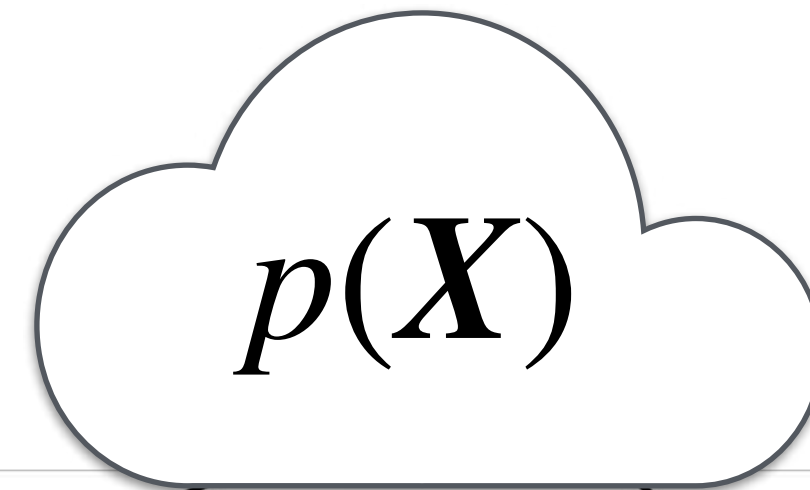
Institute of Physics, CAS

<https://wangleiphy.github.io>



Generative models and their physics genes

Goodfellow,
NIPS tutorial, 1701.00160



Explicit density

Implicit density

Direct
GAN

Tractable density

Approximate density

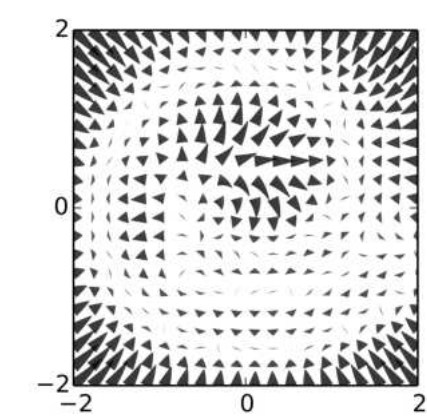
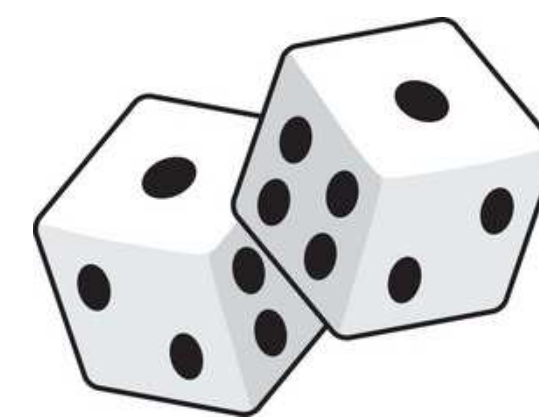
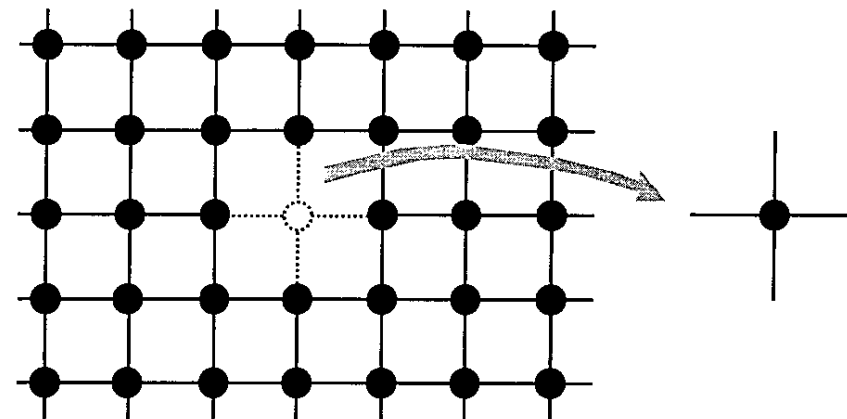
Markov Chain
GSN

Variational

Markov Chain

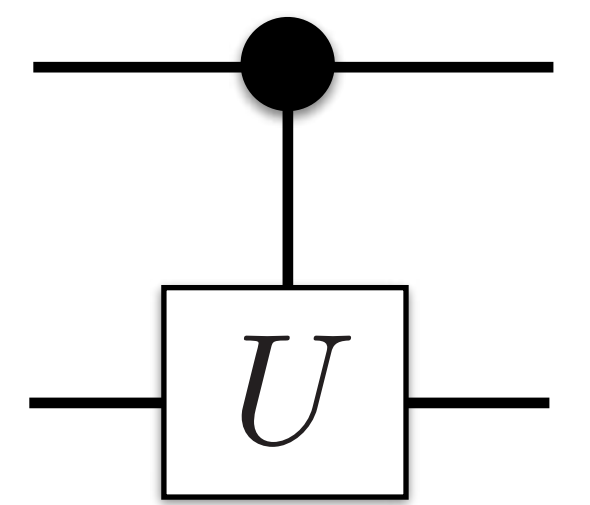
- Fully visible belief nets
- NADE
- MADE
- PixelRNN
- Change of variables models (nonlinear ICA)

Variational autoencoder Boltzmann machine + **Diffusion models**



Tensor Networks

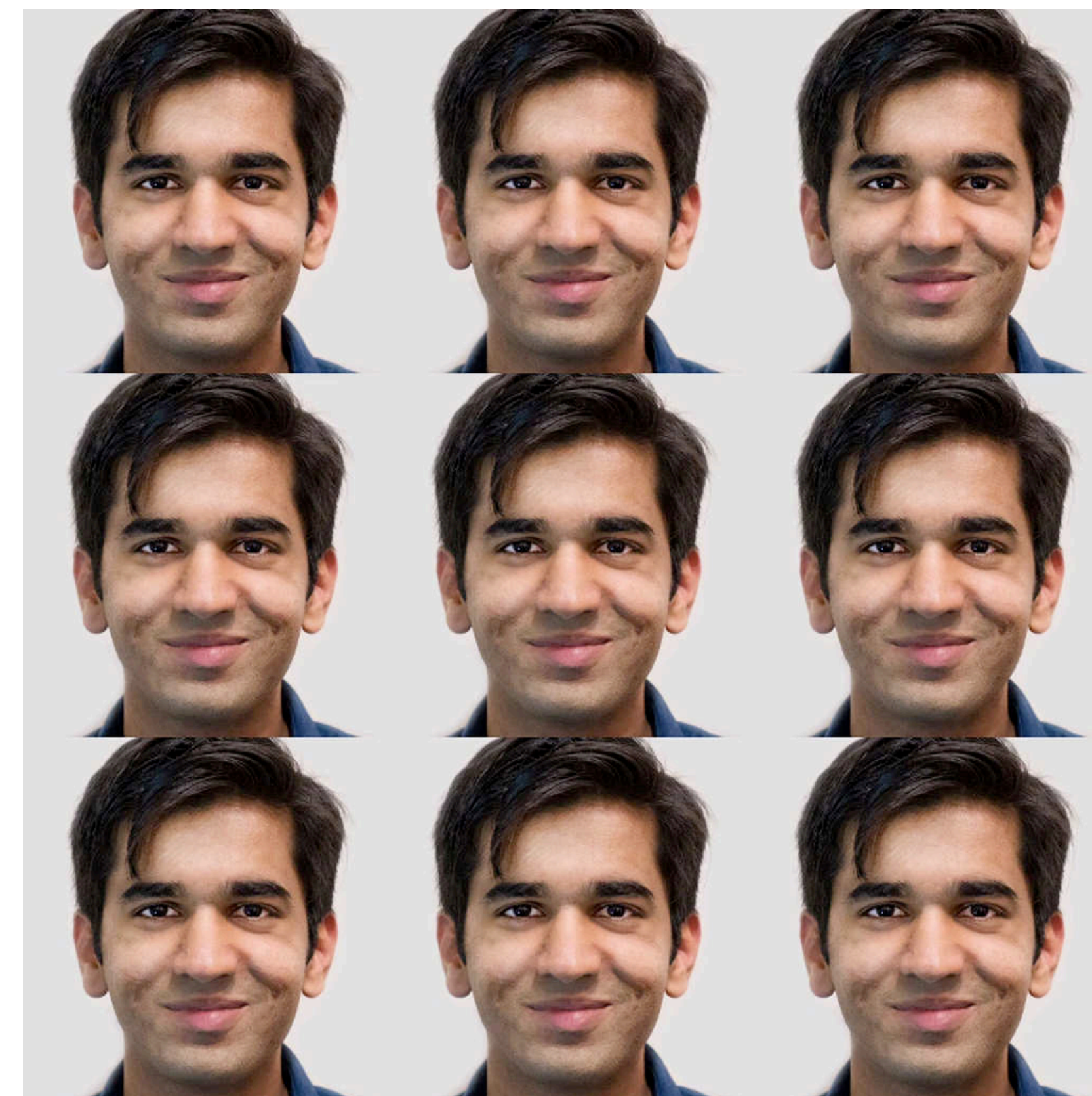
Han et al, PRX '18



Quantum Circuits

Liu et al PRA '18

Normalizing flows



 Parallel WaveNet 1711.10433

<https://deepmind.com/blog/high-fidelity-speech-synthesis-wavenet/>

 Glow 1807.03039

<https://blog.openai.com/glow/>

Normalizing flows



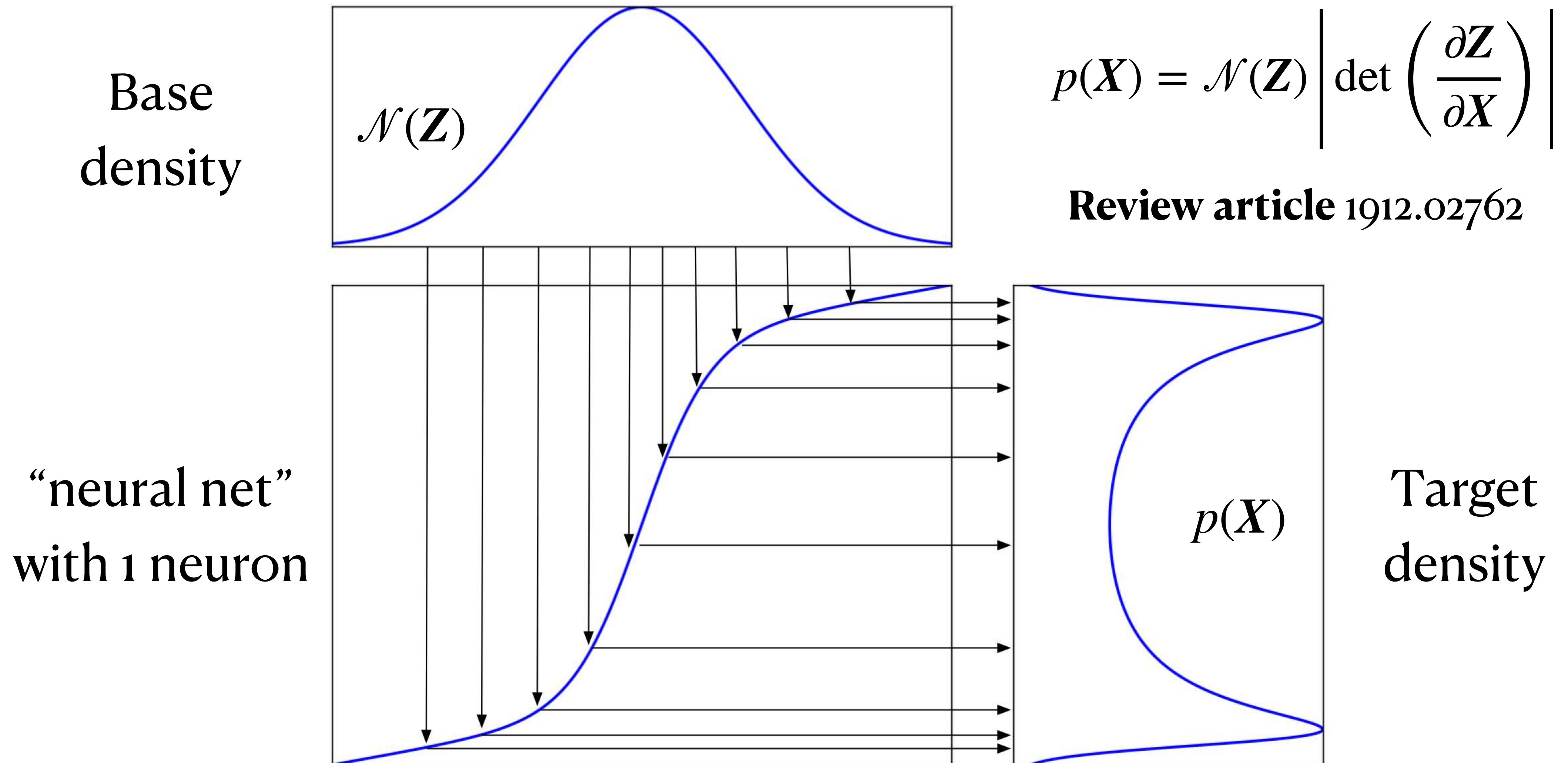
 Parallel WaveNet 1711.10433

<https://deepmind.com/blog/high-fidelity-speech-synthesis-wavenet/>

 Glow 1807.03039

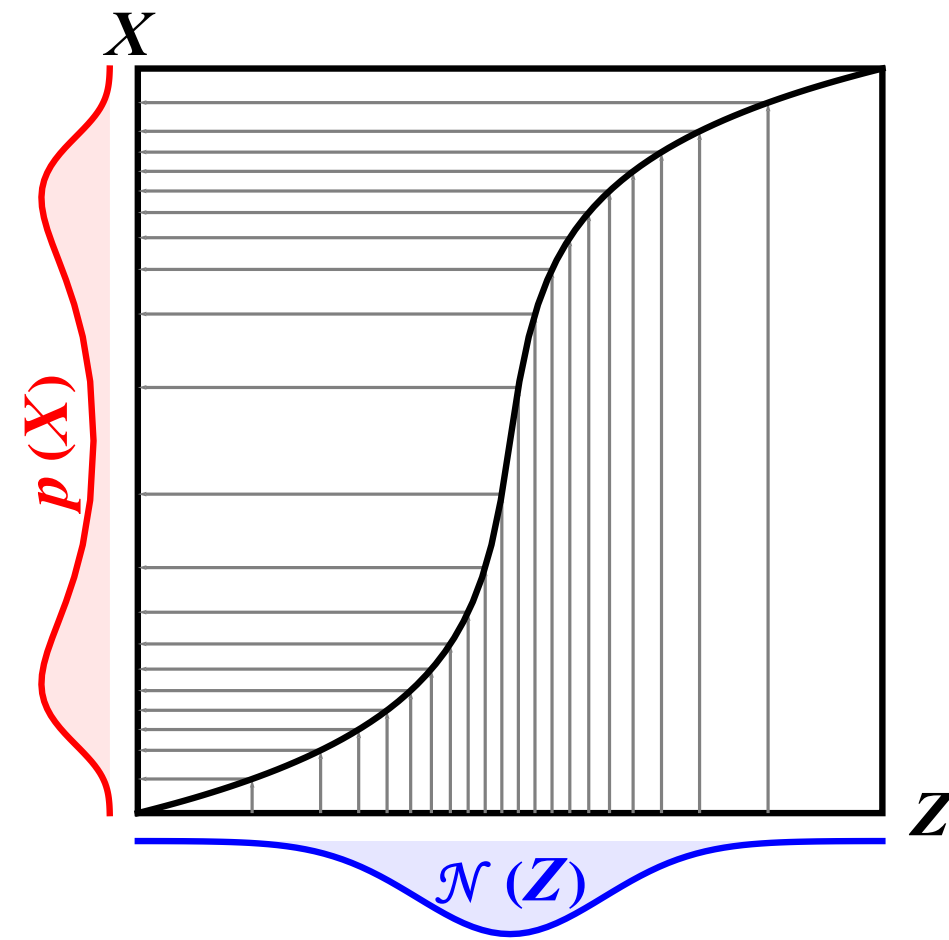
<https://blog.openai.com/glow/>

Normalizing flow in a nutshell



Flow model

$$p(\mathbf{X}) = \mathcal{N}(\mathbf{Z}) \left| \det \left(\frac{\partial \mathbf{Z}}{\partial \mathbf{X}} \right) \right|$$



Sampling

$$\mathbf{Z} \sim \mathcal{N}(\mathbf{Z})$$

$$\mathbf{Z} \mapsto \mathbf{X}$$

Normalization

$$\int d\mathbf{X} \left| \det \left(\frac{\partial \mathbf{Z}}{\partial \mathbf{X}} \right) \right| = \int d\mathbf{Z}$$

From flow to diffusion, and back

A unified perspective to transportation-based generative model

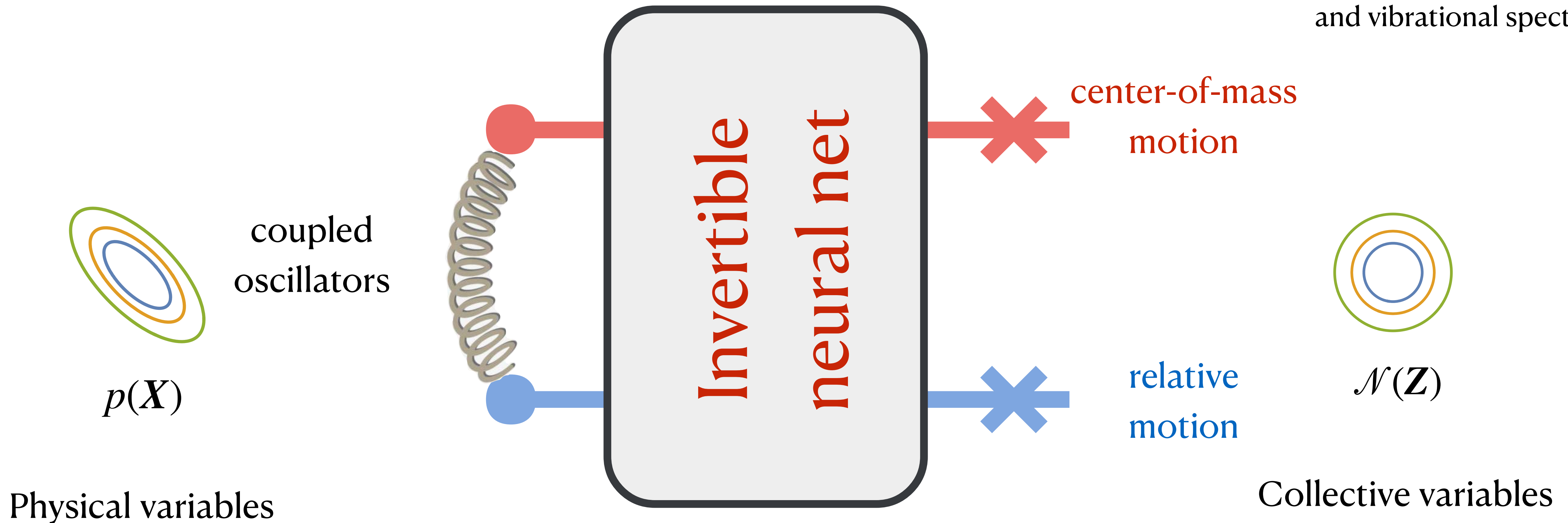
Physics intuition of normalizing flow

Li and LW, PRL '18

Li, Dong, Zhang, LW, PRX '20

Zhang, Wang, LW, JCP '24

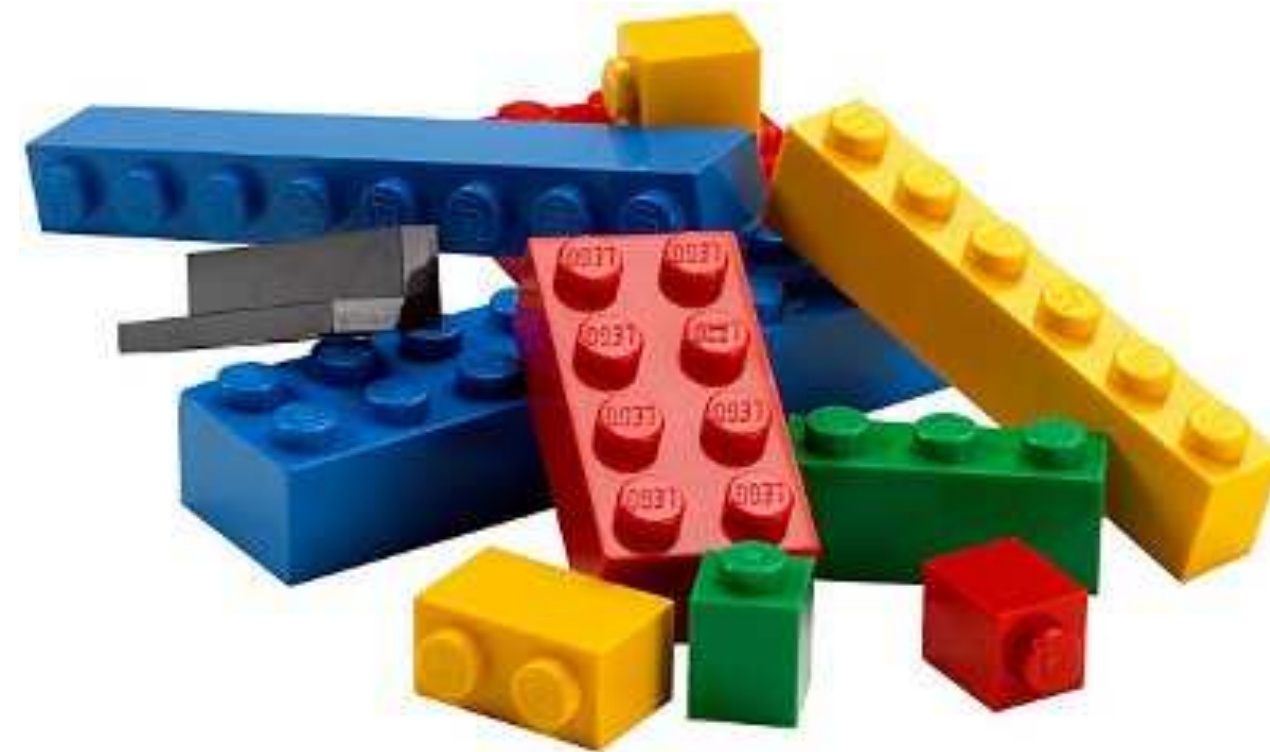
for RG, canonical transformations,
and vibrational spectra



High-dimensional, composable, learnable, nonlinear transformations

Flow architecture design

Composability

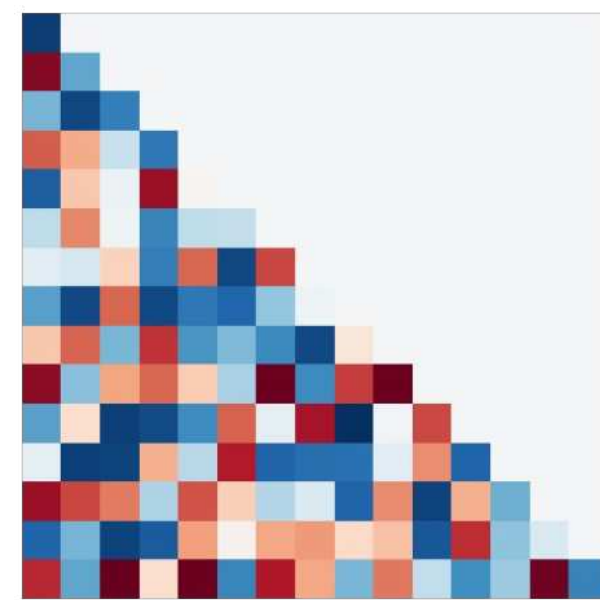


$$Z = \mathcal{T}(X)$$

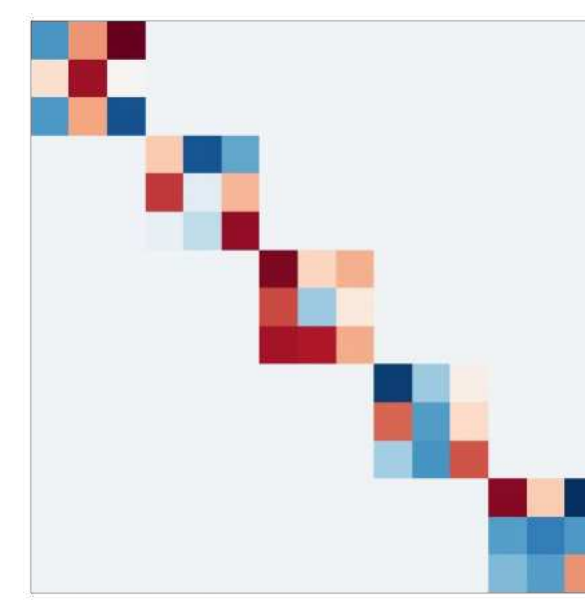
$$\mathcal{T} = \mathcal{T}_1 \circ \mathcal{T}_2 \circ \mathcal{T}_3 \circ \dots$$

**Balanced
efficiency &
inductive bias**

$$\left| \det \left(\frac{\partial Z}{\partial X} \right) \right|$$



Autoregressive



Blockwise

$$\frac{\partial p(X, t)}{\partial t} + \nabla \cdot [p(X, t)\mathbf{v}] = 0$$

Continuous flow

Example of a building block

Forward

$$\begin{cases} \mathbf{x}_{<} = \mathbf{z}_{<} \\ \mathbf{x}_{>} = \mathbf{z}_{>} \odot e^{s(\mathbf{z}_{<})} + t(\mathbf{z}_{<}) \end{cases}$$

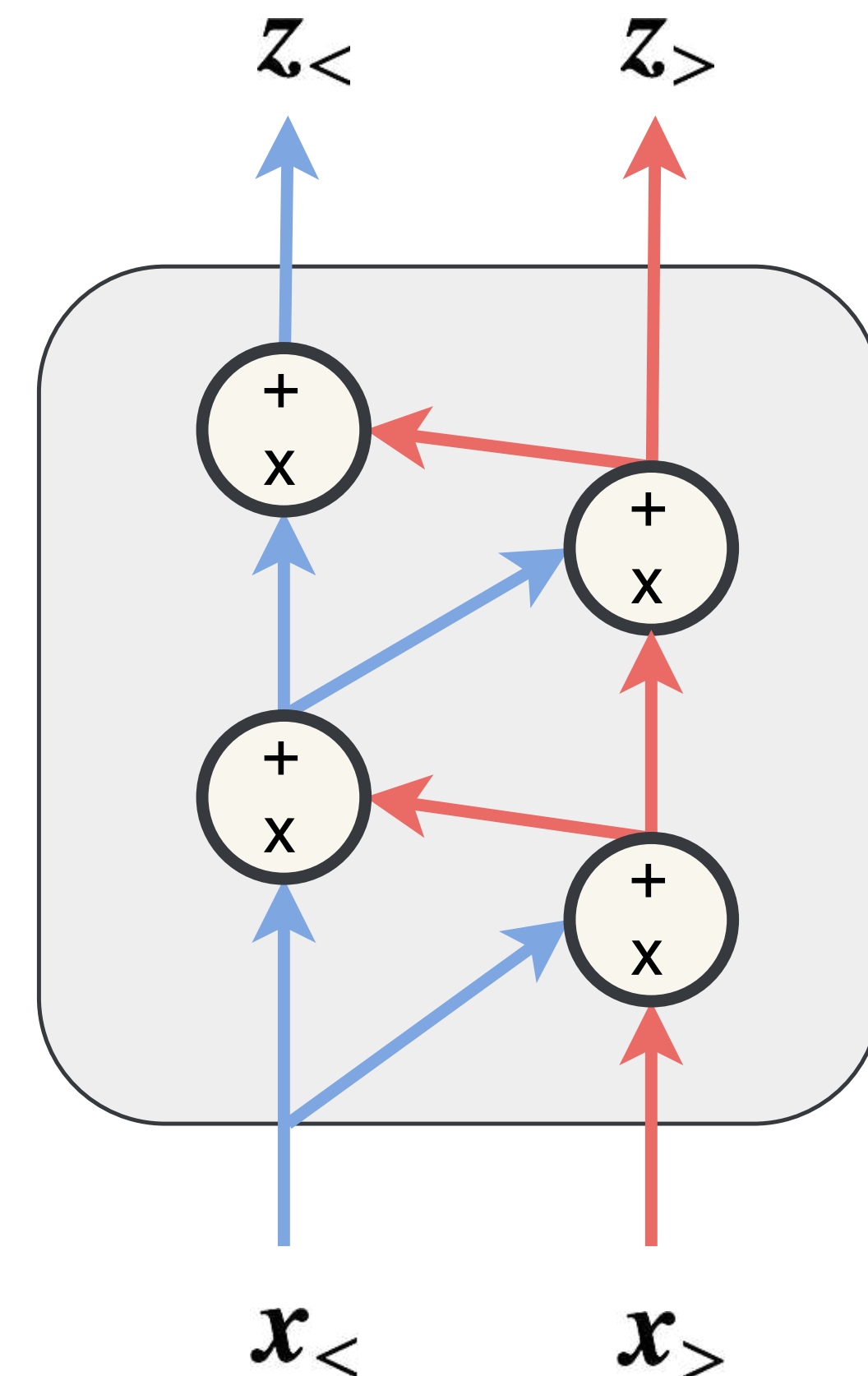
arbitrary
neural nets

Inverse

$$\begin{cases} \mathbf{z}_{<} = \mathbf{x}_{<} \\ \mathbf{z}_{>} = (\mathbf{x}_{>} - t(\mathbf{x}_{<})) \odot e^{-s(\mathbf{x}_{<})} \end{cases}$$

Log-Abs-Jacobian-Det

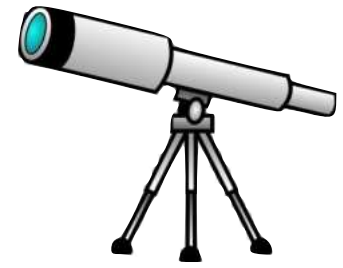
$$\ln \left| \det \left(\frac{\partial \mathbf{x}}{\partial \mathbf{z}} \right) \right| = \sum_i [s(\mathbf{z}_{<})]_i$$



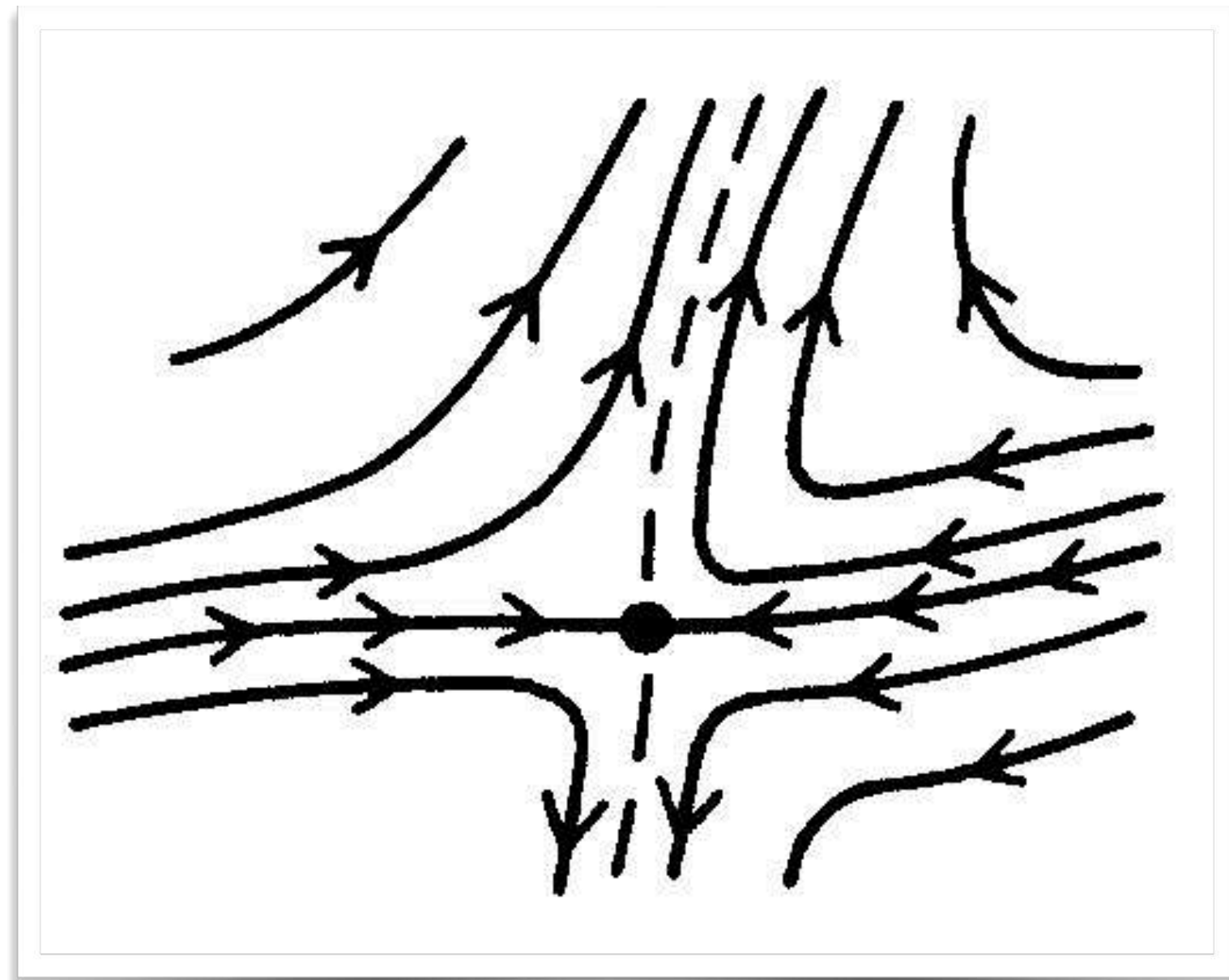
Real NVP, Dinh et al, 1605.08803

Turns out to have surprising connection Störmer–Verlet integration

Why is flow useful for physics?



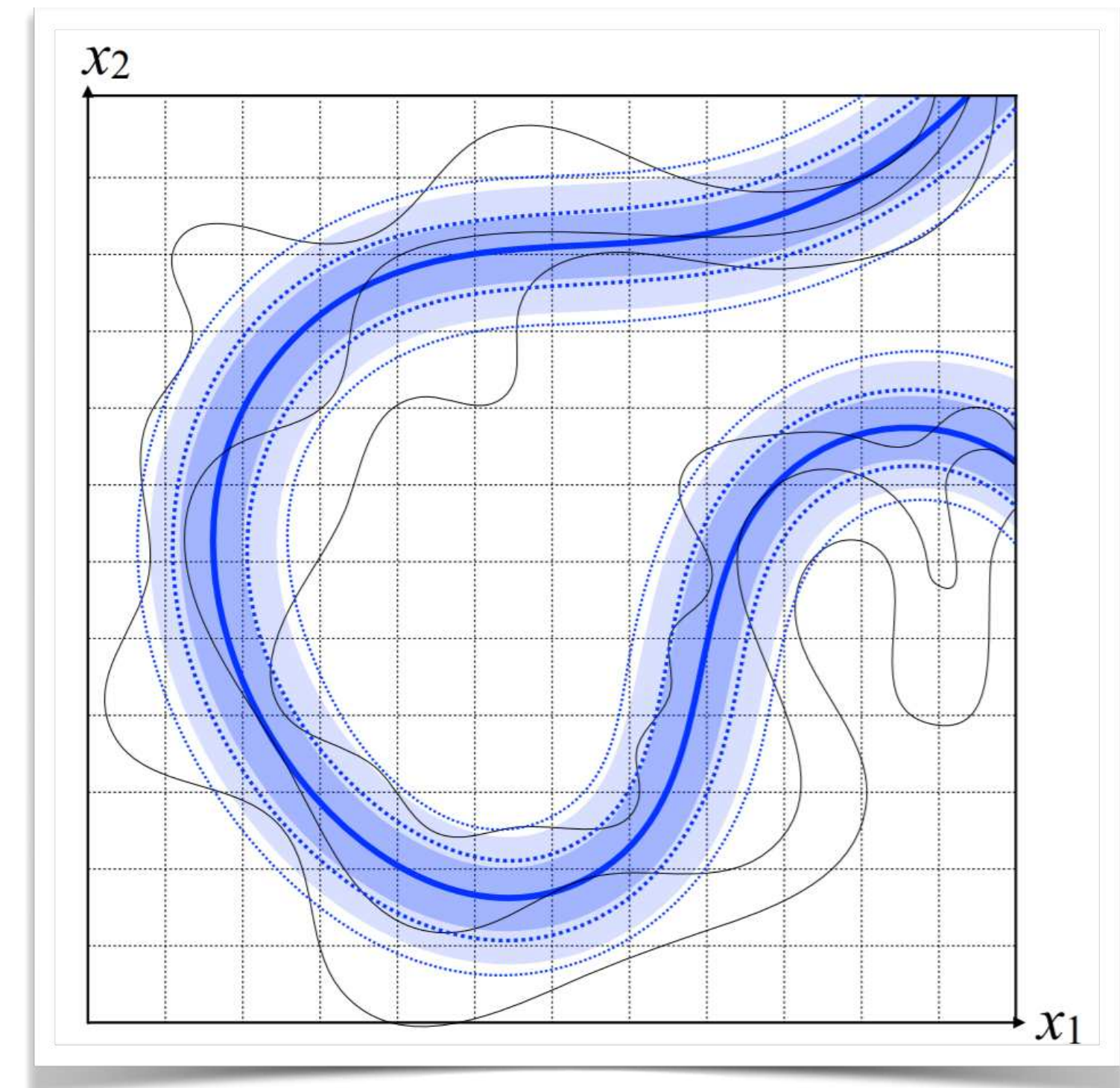
Renormalization group



Effective theory emerges upon transformation of the variables



Monte Carlo update

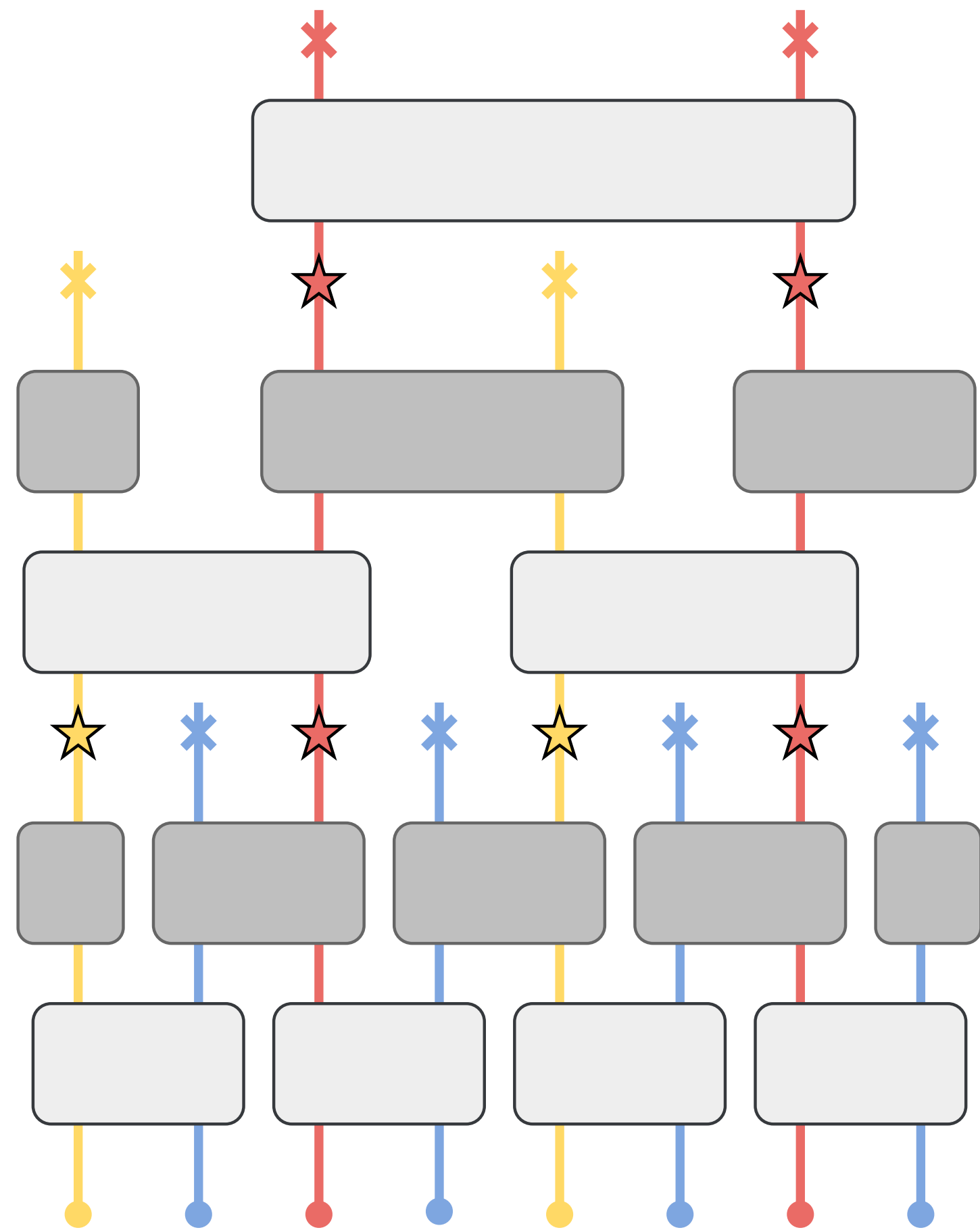


Physics happens on a manifold
Train neural nets to unfold that manifold

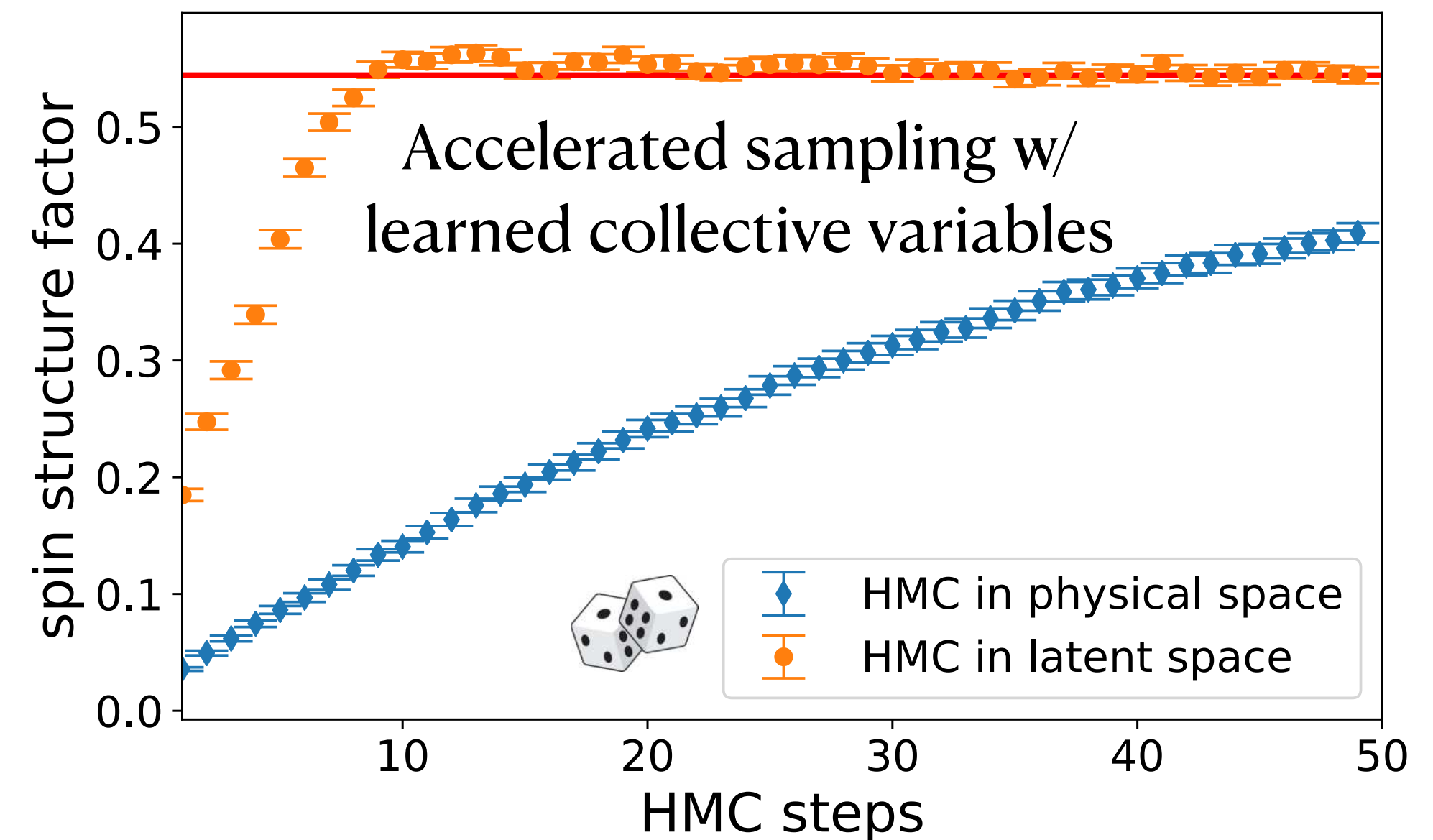
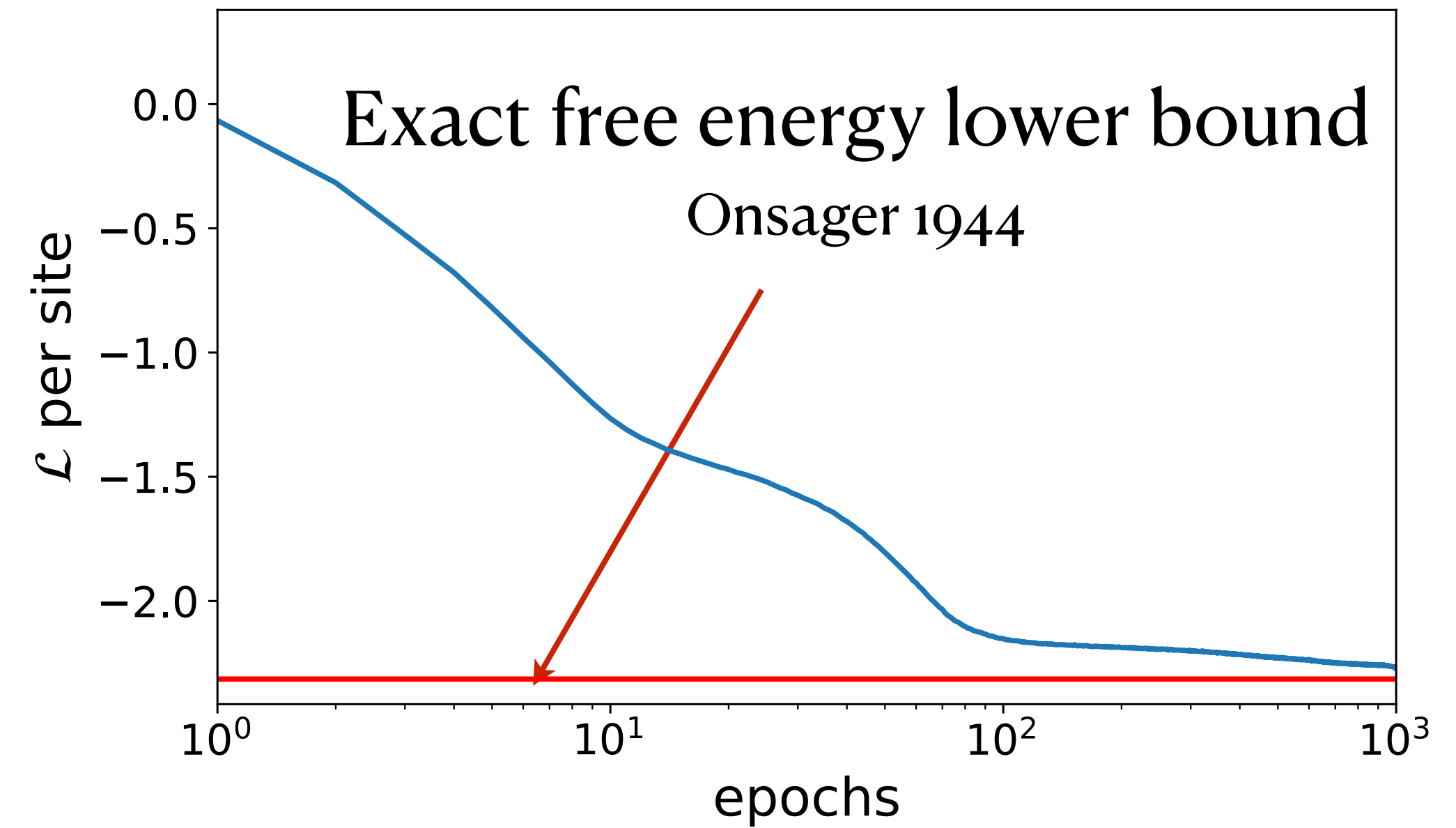
Neural network renormalization group

Li, LW, PRL '18 [li012589/NeuralRG](https://arxiv.org/abs/1801.01258)

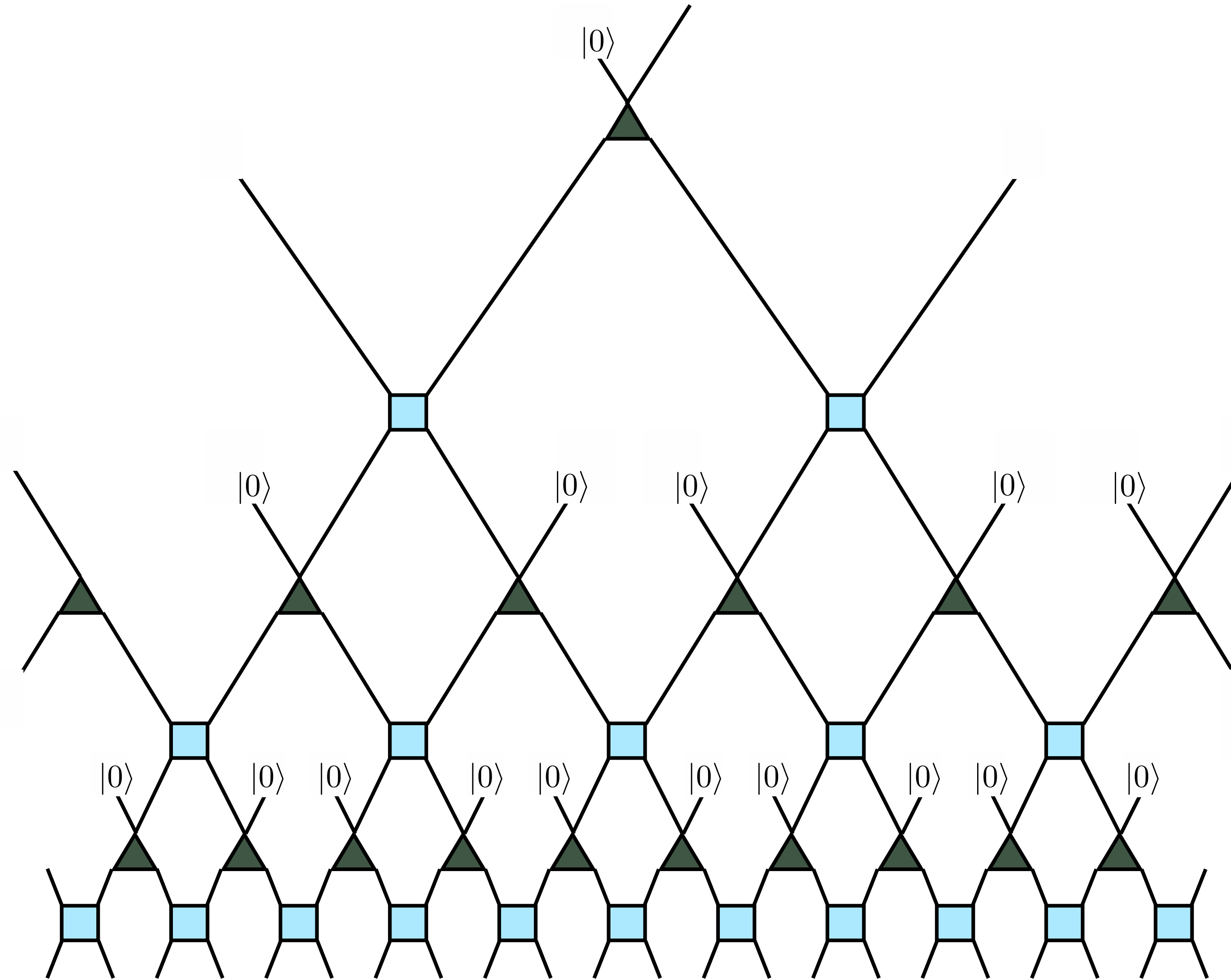
Collective variables



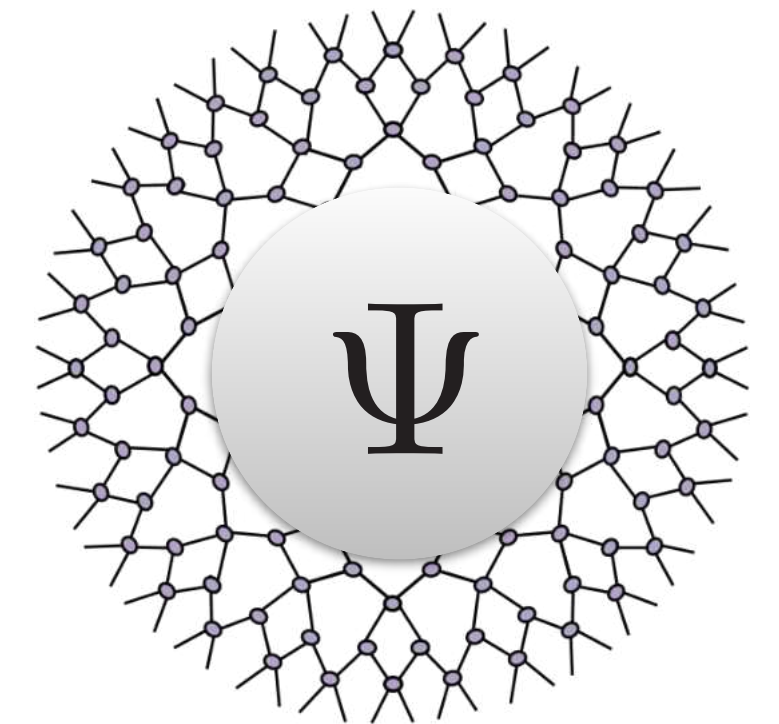
Physical variables



Quantum version of the architecture

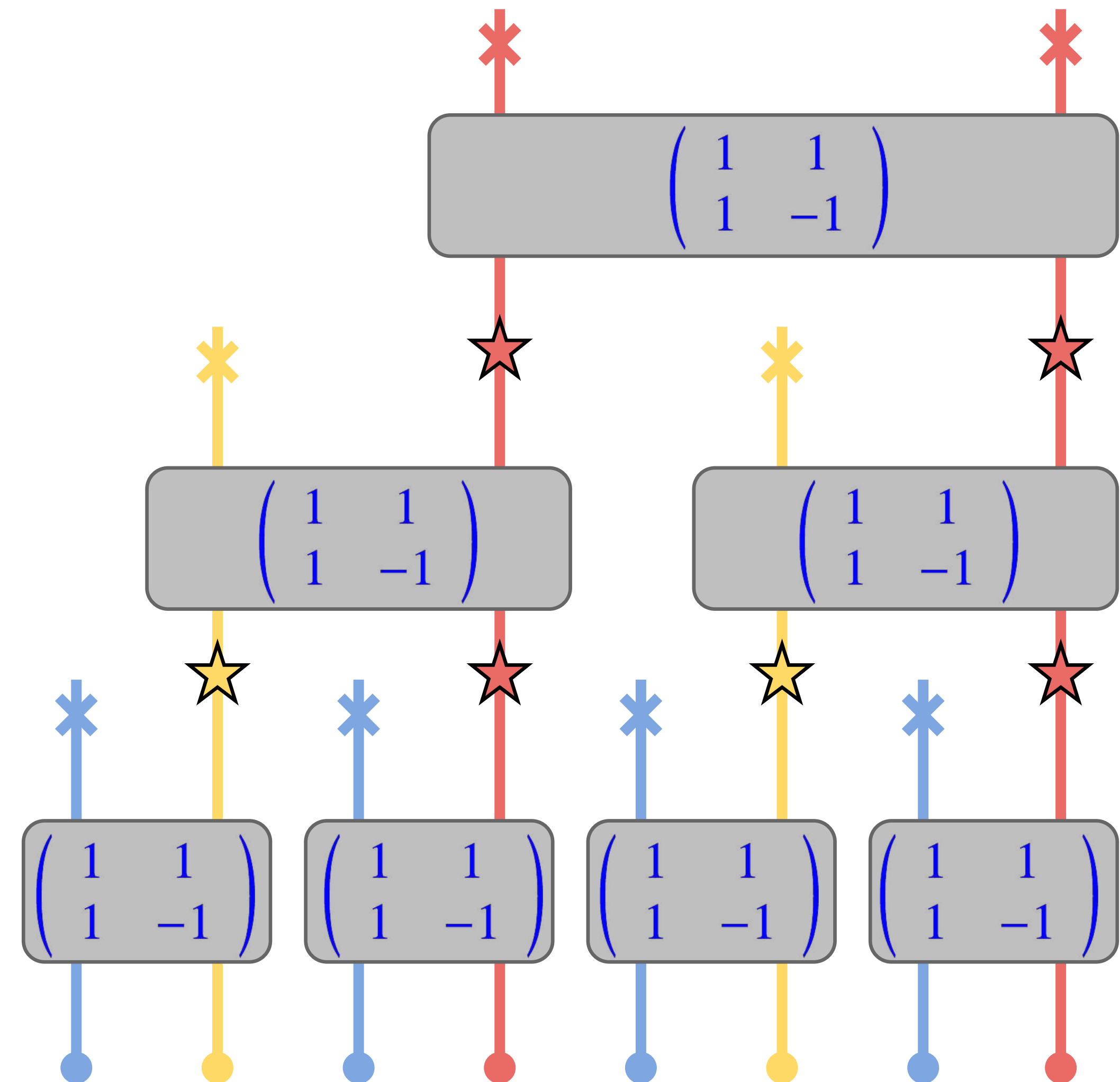
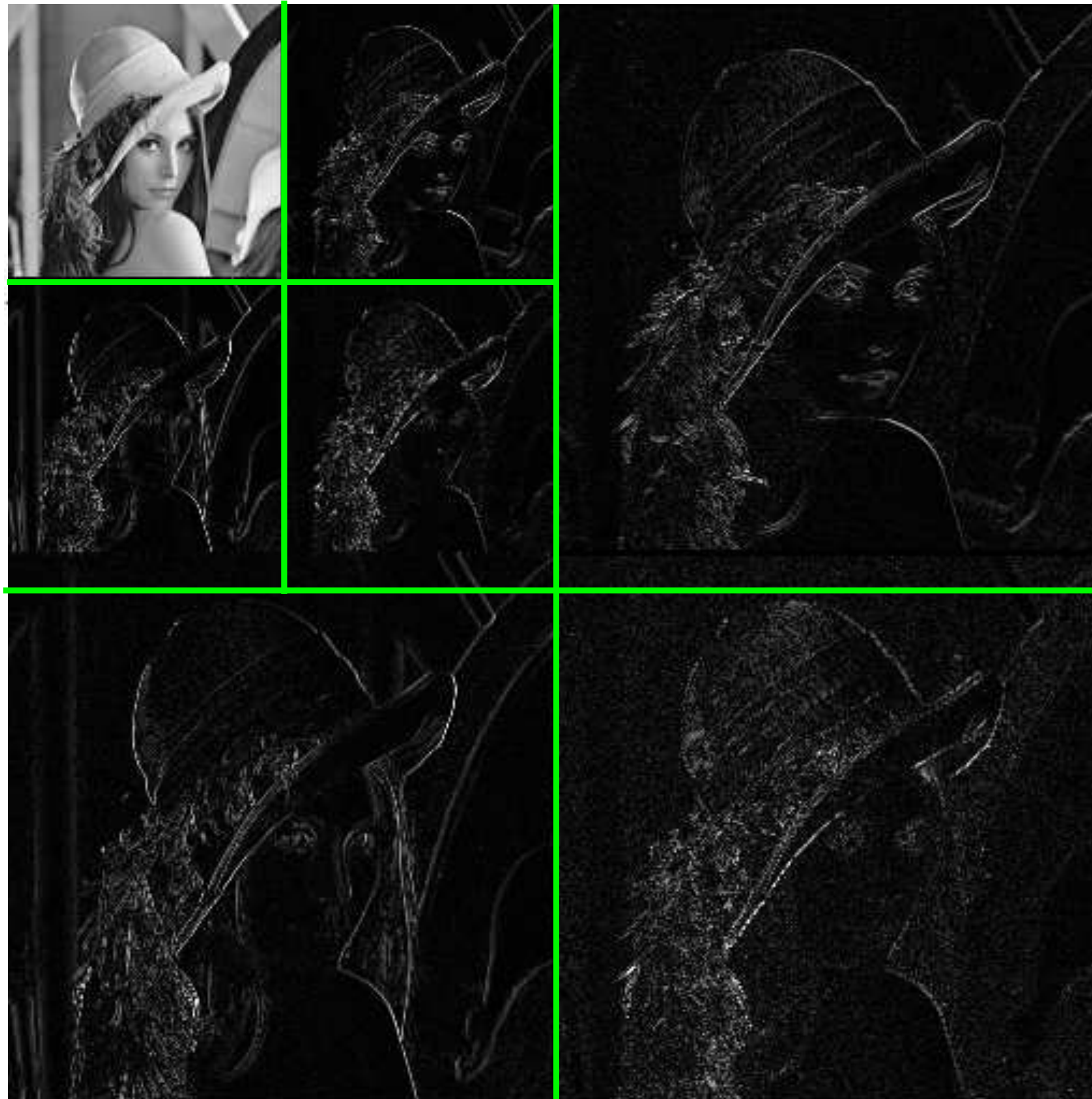


Entangled qubits



**Multi-Scale
Entanglement
Renormalization
Ansatz**

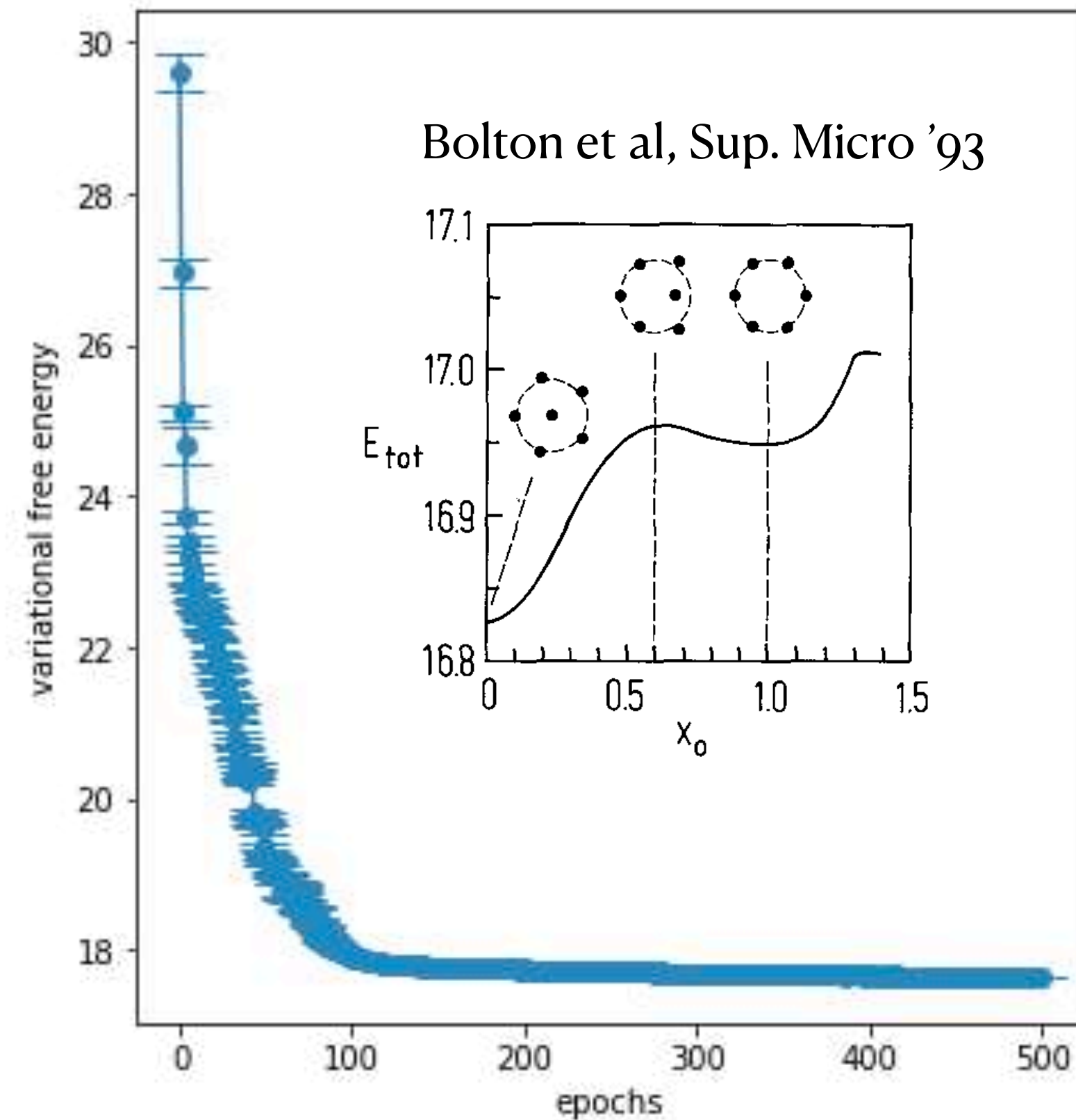
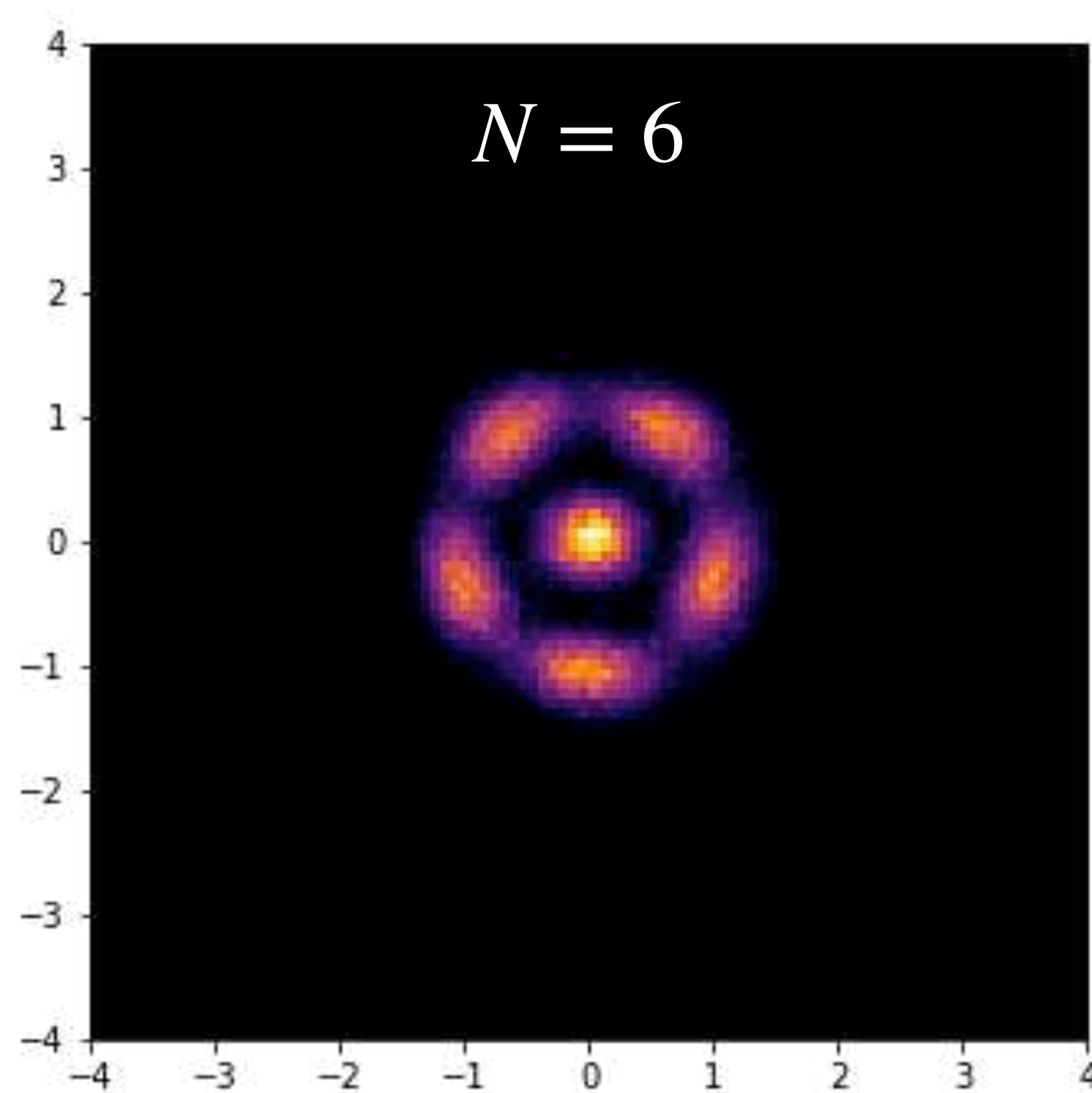
Connection to wavelets



Nonlinear & adaptive generalizations of wavelets

Demo: Classical Coulomb gas in a harmonic trap

$$E = \sum_{i < j} \frac{1}{|x_i - x_j|} + \sum_i^N x_i^2$$



Continuous normalizing flows

$$\ln p(\mathbf{X}) = \ln \mathcal{N}(\mathbf{Z}) - \ln \left| \det \left(\frac{\partial \mathbf{X}}{\partial \mathbf{Z}} \right) \right|$$

Consider infinitesimal change-of-variables [Chen et al 1806.07366](#)

$$\mathbf{X} = \mathbf{Z} + \varepsilon \mathbf{v} \quad \ln p(\mathbf{X}) - \ln \mathcal{N}(\mathbf{Z}) = - \ln \left| \det \left(1 + \varepsilon \frac{\partial \mathbf{v}}{\partial \mathbf{Z}} \right) \right|$$

$$\varepsilon \rightarrow 0$$

$$\frac{d\mathbf{X}}{dt} = \mathbf{v}$$

$$\frac{d \ln p(\mathbf{X}, t)}{dt} = - \nabla \cdot \mathbf{v}$$

Continuous normalizing flows

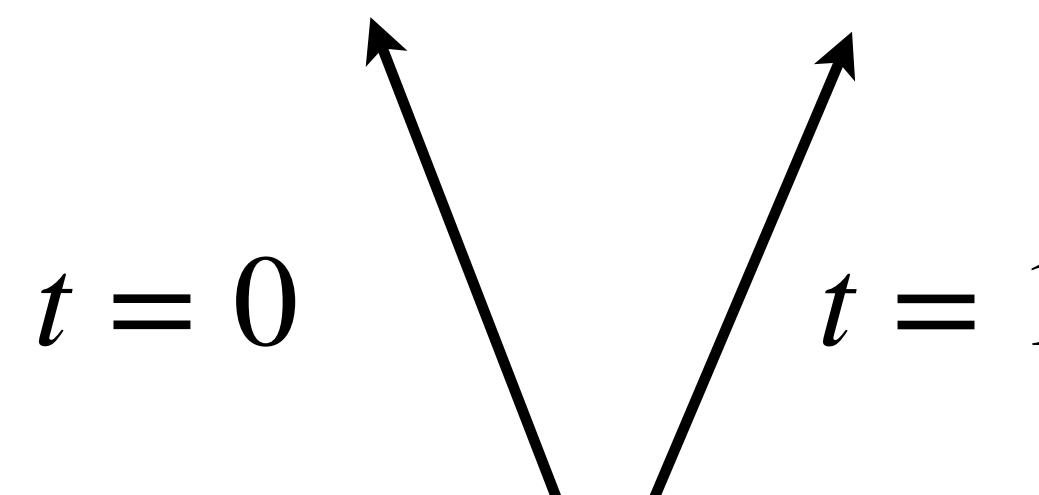
$$\ln p(\mathbf{X}) = \ln \mathcal{N}(\mathbf{Z}) - \ln \left| \det \left(\frac{\partial \mathbf{X}}{\partial \mathbf{Z}} \right) \right|$$

Consider infinitesimal change-of-variables Chen et al 1806.07366

$$\mathbf{X} = \mathbf{Z} + \varepsilon \mathbf{v} \quad \ln p(\mathbf{X}) - \ln \mathcal{N}(\mathbf{Z}) = - \ln \left| \det \left(1 + \varepsilon \frac{\partial \mathbf{v}}{\partial \mathbf{Z}} \right) \right|$$

$\varepsilon \rightarrow 0$

$$\frac{d\mathbf{X}}{dt} = \mathbf{v}$$

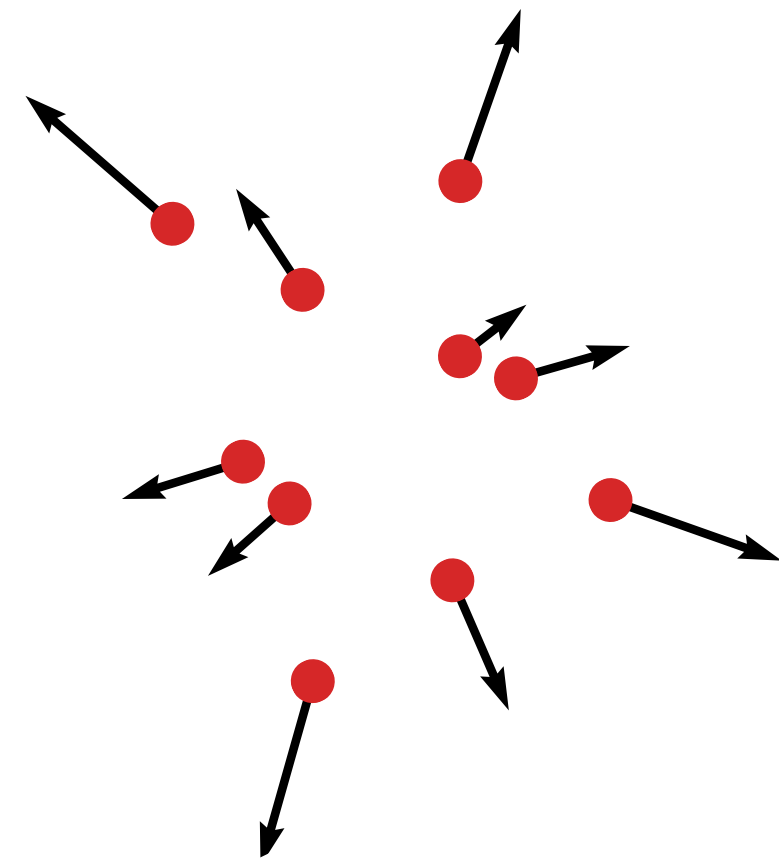
$$\frac{d \ln p(\mathbf{X}, t)}{dt} = - \nabla \cdot \mathbf{v}$$


Fluid physics behind flows

Zhang, E, LW 1809.10188

[wangleiphy/MongeAmpereFlow](https://arxiv.org/abs/1809.10188)

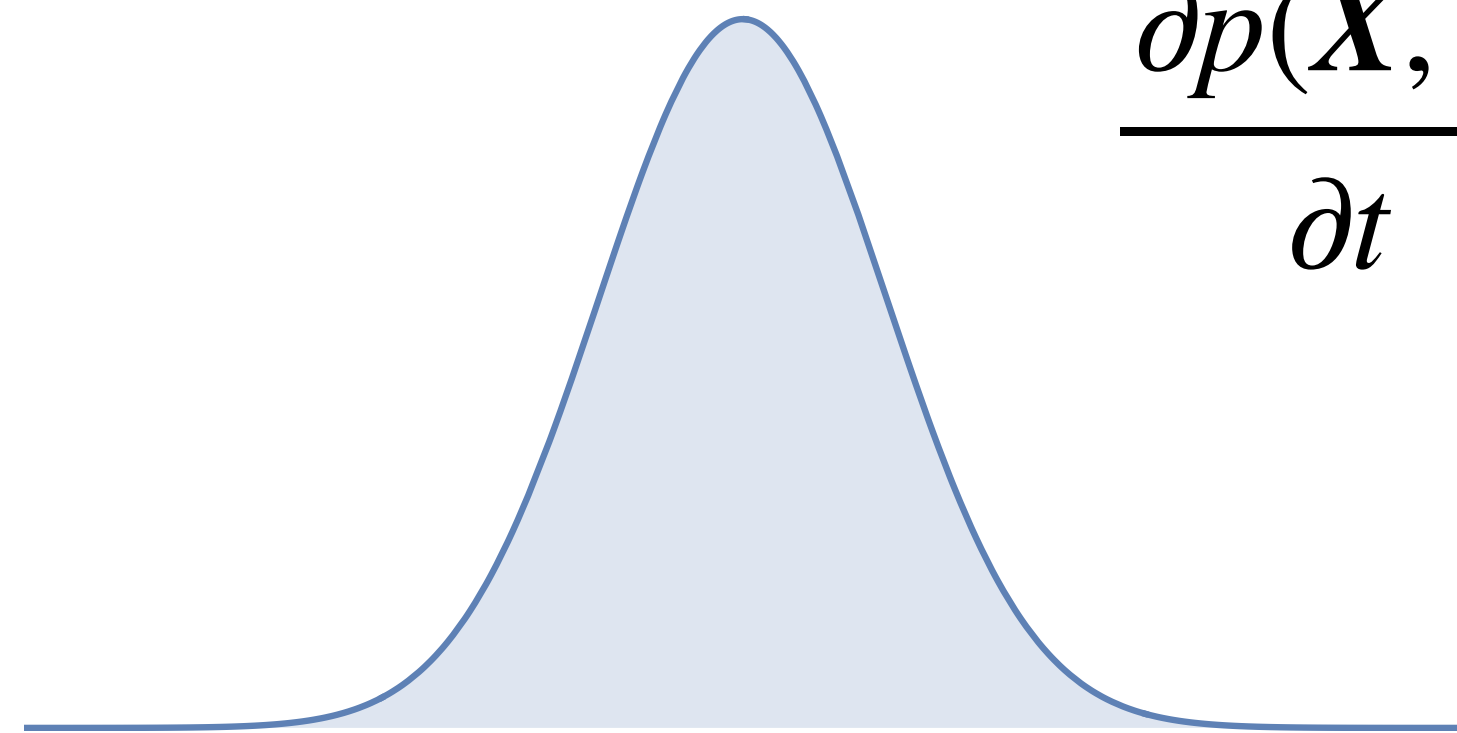
$$\frac{dX}{dt} = \mathbf{v}$$
$$\frac{d \ln p(X, t)}{dt} = - \nabla \cdot \mathbf{v}$$



$$\frac{d}{dt} = \frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla$$

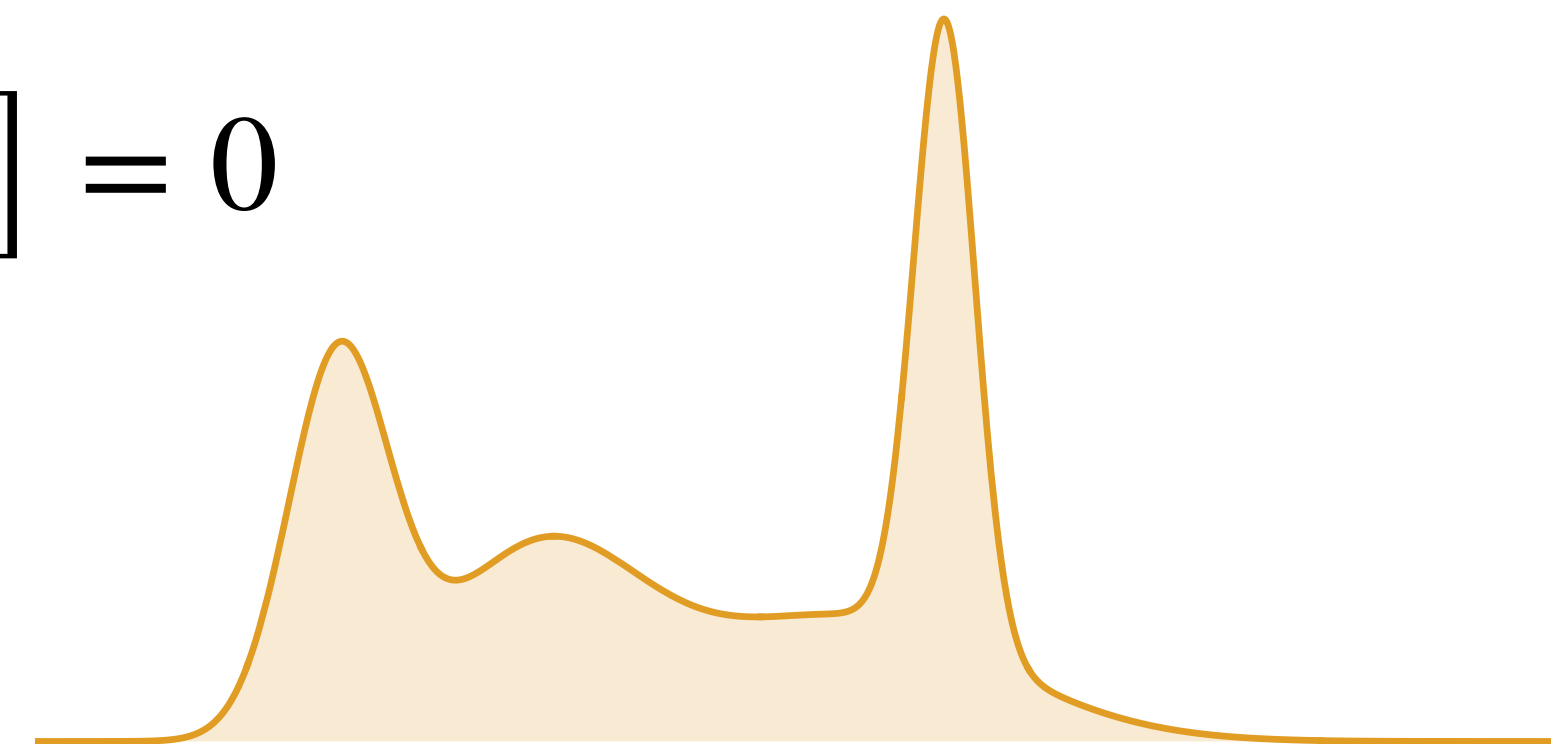
“material derivative”

Lagrangian v.s. Euler approach to fluid mechanics



Simple density

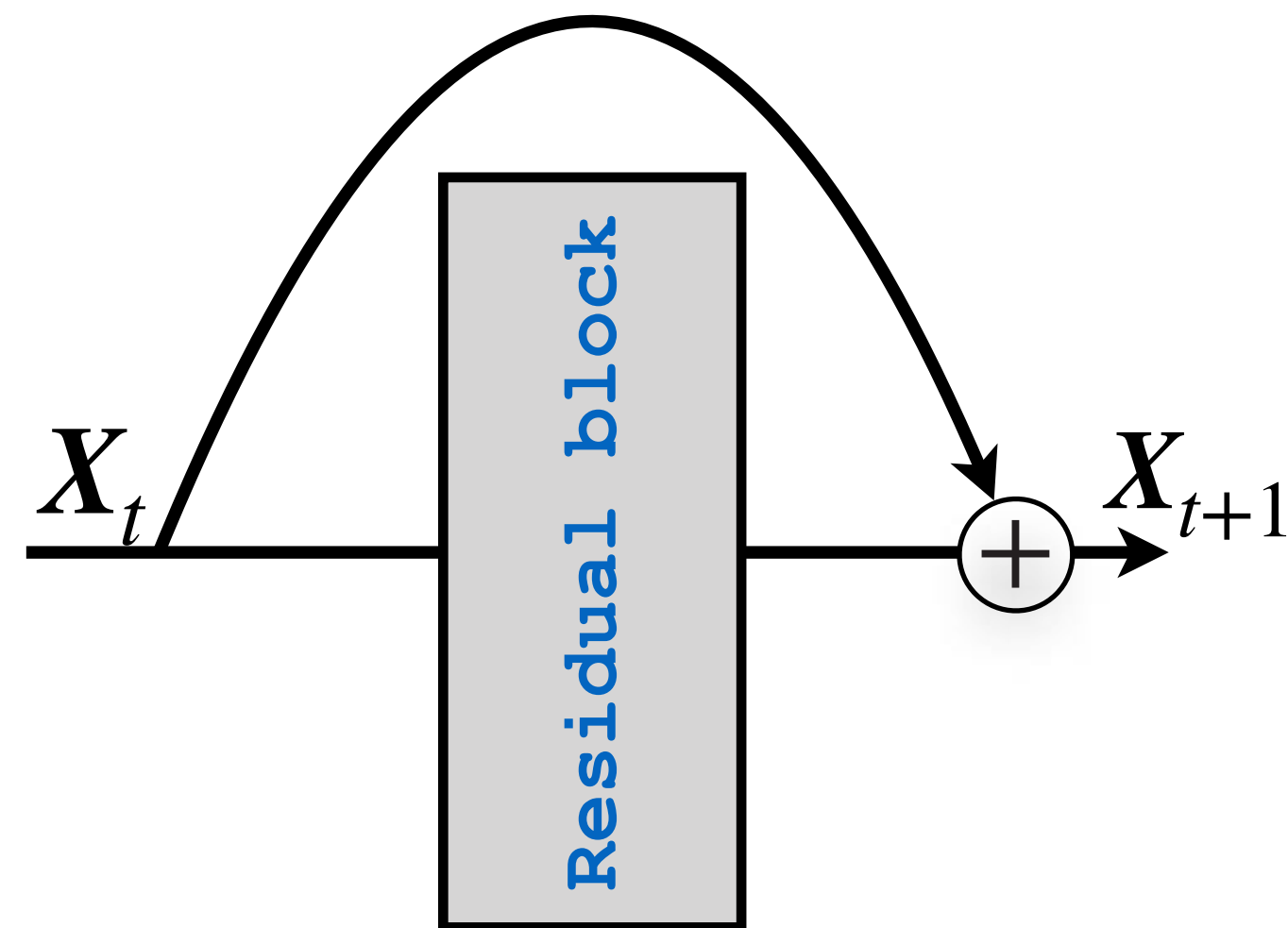
$$\frac{\partial p(X, t)}{\partial t} + \nabla \cdot [p(X, t)\mathbf{v}] = 0$$



Complex density

Neural Ordinary Differential Equations

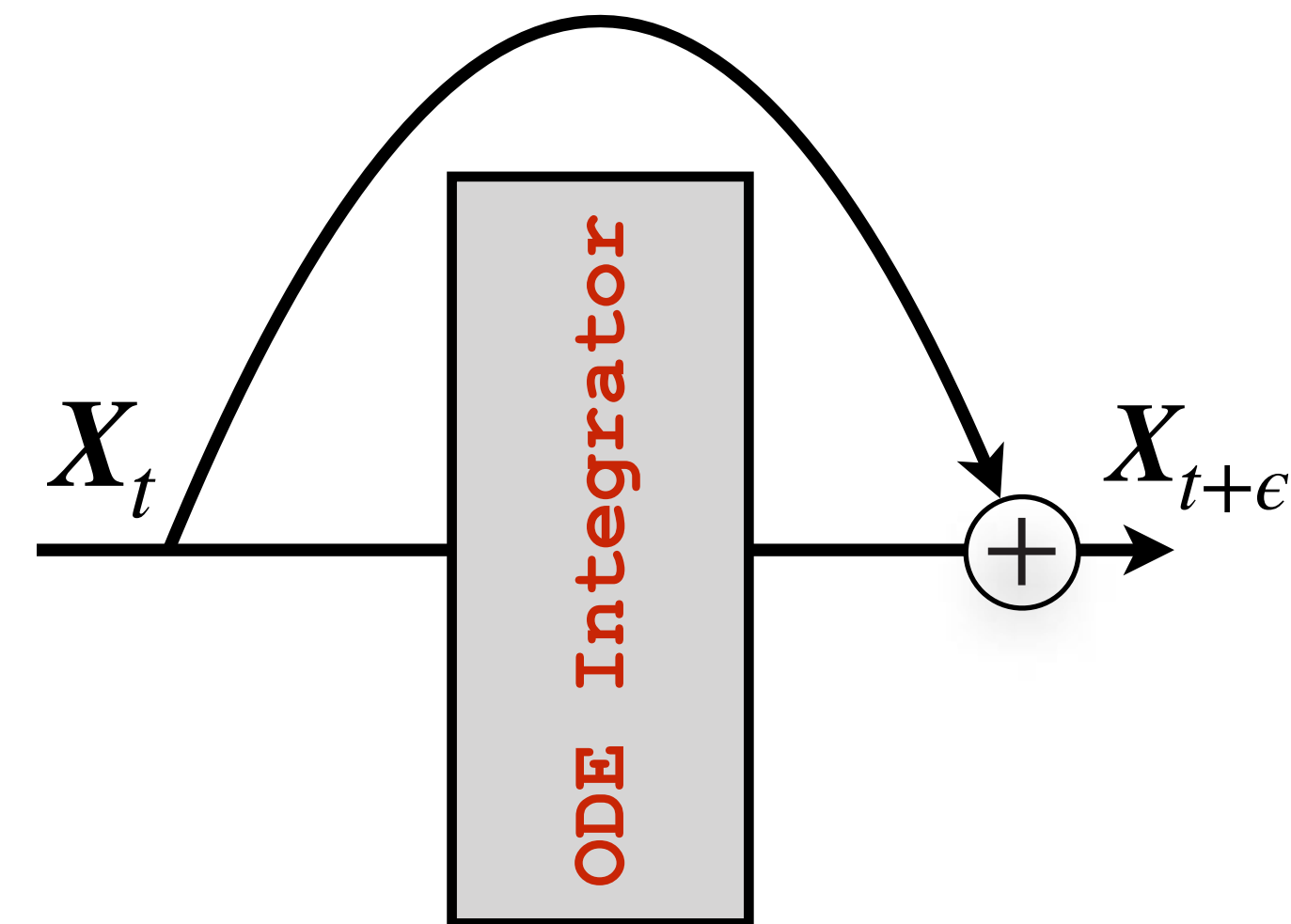
Residual network



$$X_{t+1} = X_t + v(X_t)$$

Chen et al, 1806.07366

ODE integration

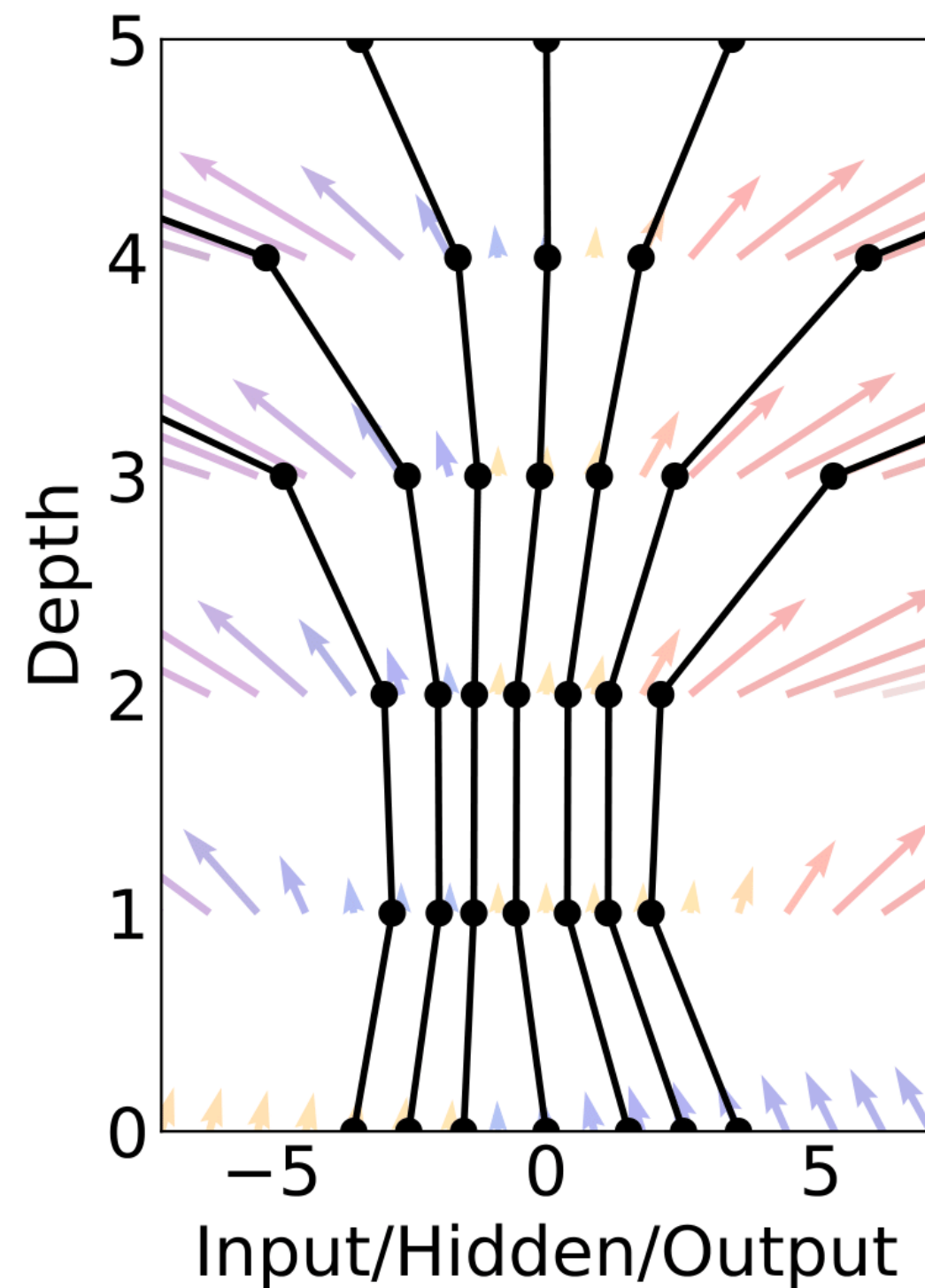


$$dX/dt = v(X)$$

Harbor et al 1705.03341
Lu et al 1710.10121,
E Commun. Math. Stat 17'...

Neural Ordinary Differential Equations

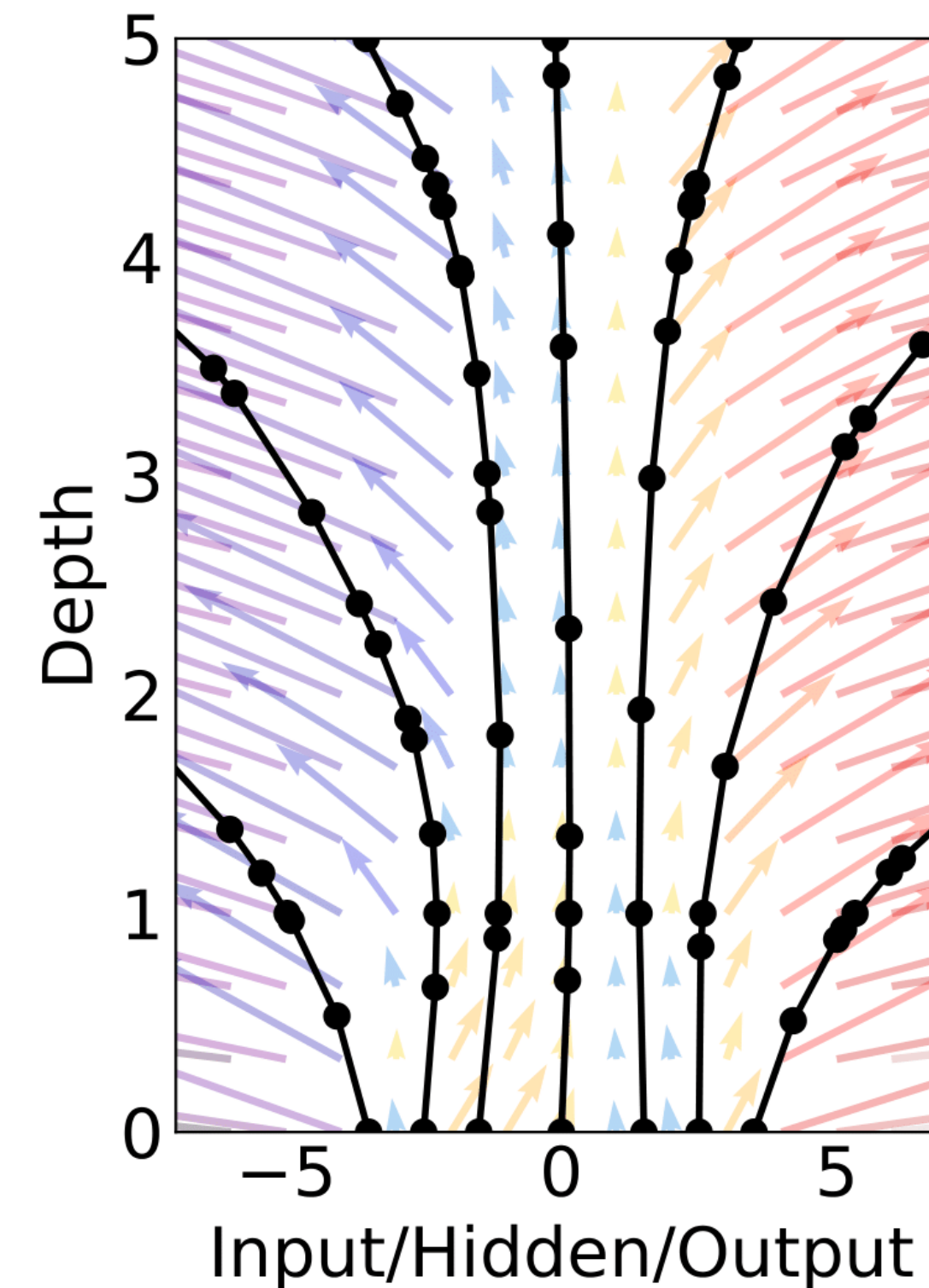
Residual network



$$X_{t+1} = X_t + \nu(X_t)$$

Chen et al, 1806.07366

ODE integration



$$dX/dt = \nu(X)$$

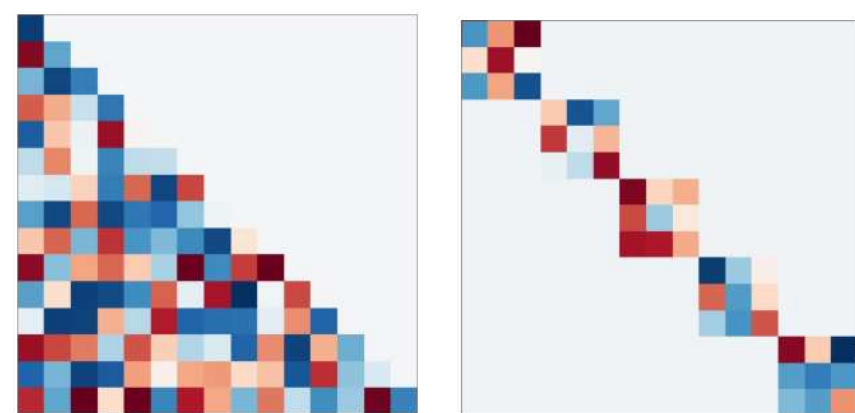
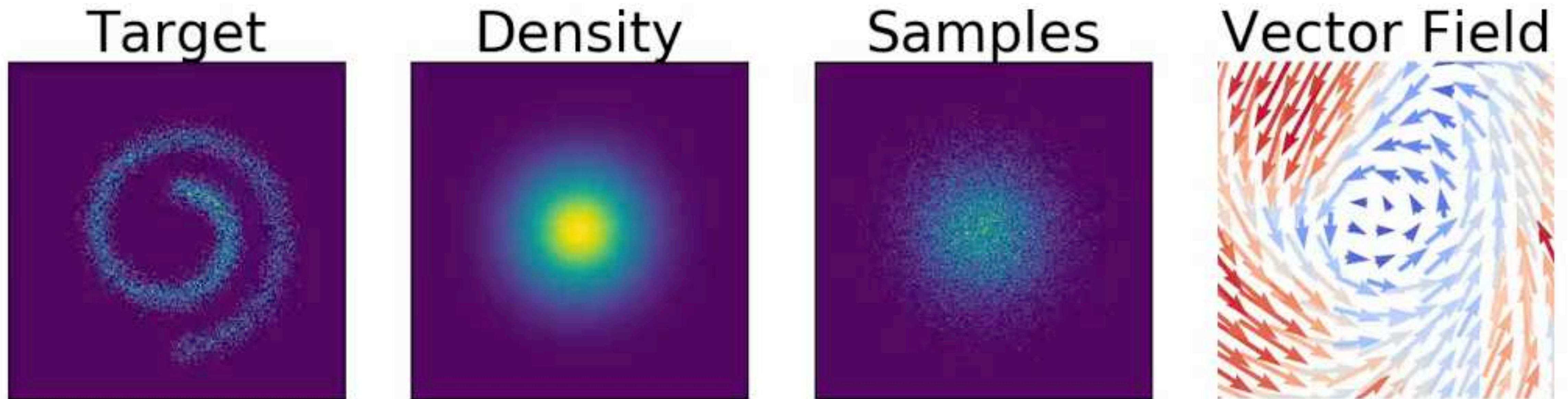
Harbor et al 1705.03341

Lu et al 1710.10121,

E Commun. Math. Stat 17'...

Continuous normalizing flows implemented with NeuralODE

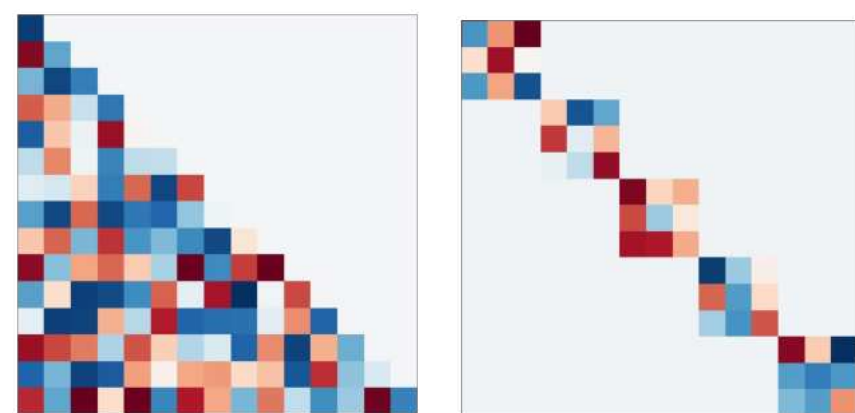
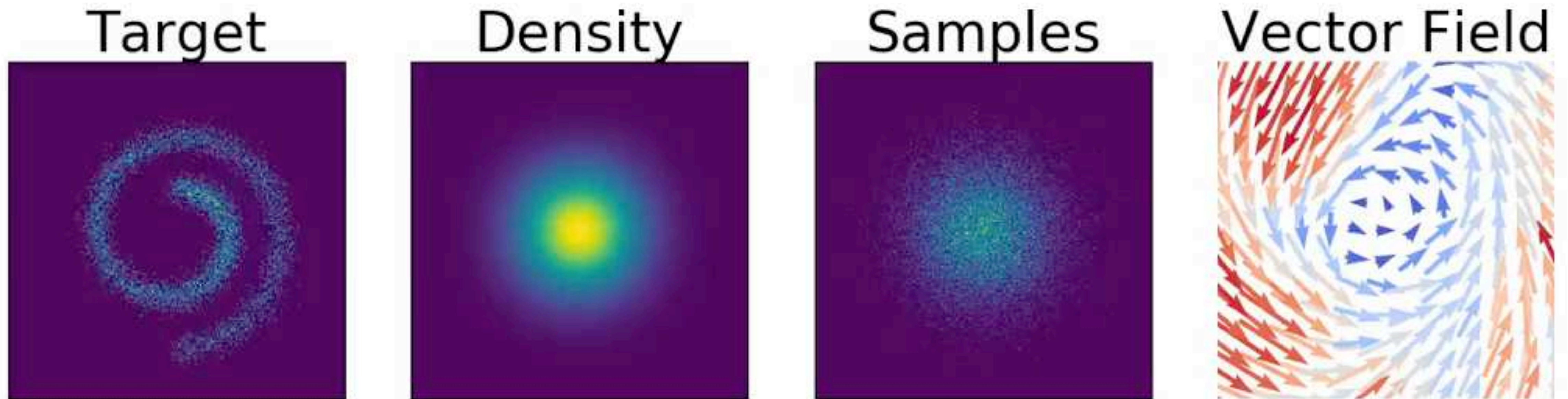
Chen et al, 1806.07366, Grathwohl et al 1810.01367



Continuous normalizing flow have no structural constraints on the transformation Jacobian

Continuous normalizing flows implemented with NeuralODE

Chen et al, 1806.07366, Grathwohl et al 1810.01367

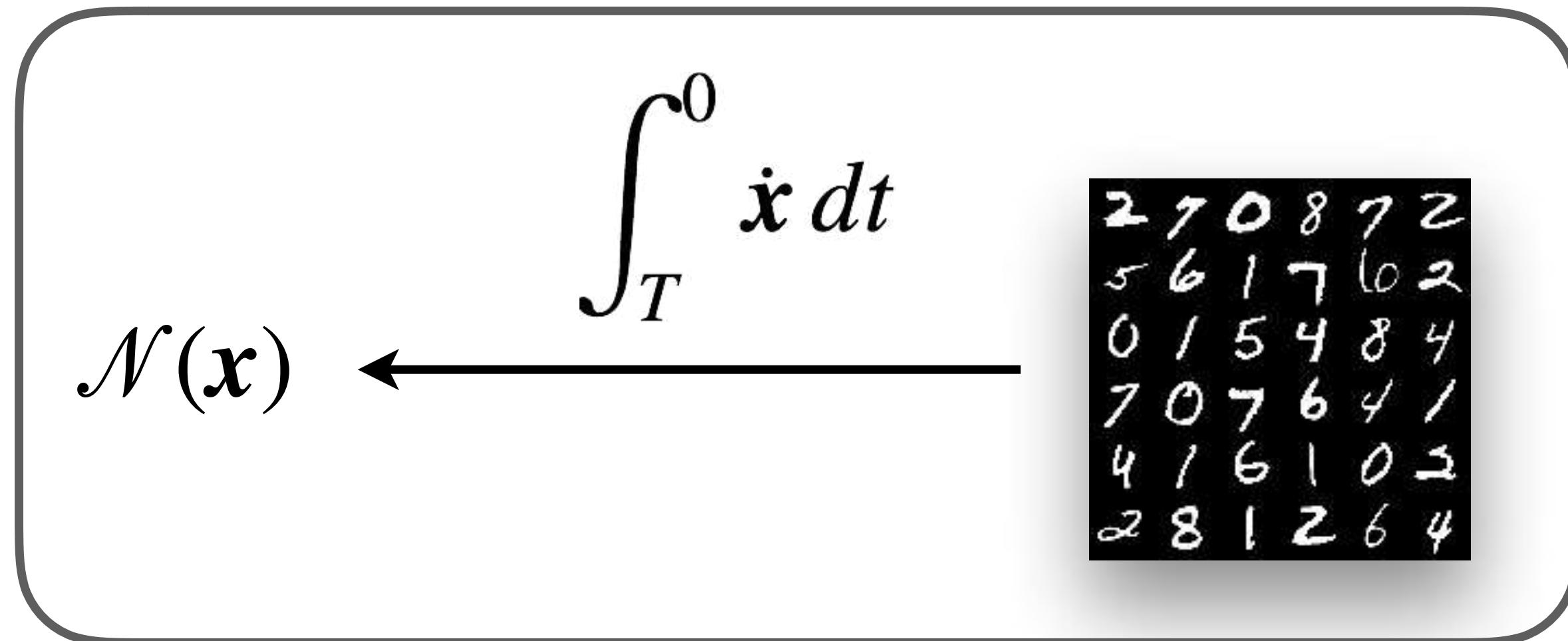


Continuous normalizing flow have no structural constraints on the transformation Jacobian

The two use cases

Zhang, E, LW, 1809.10188

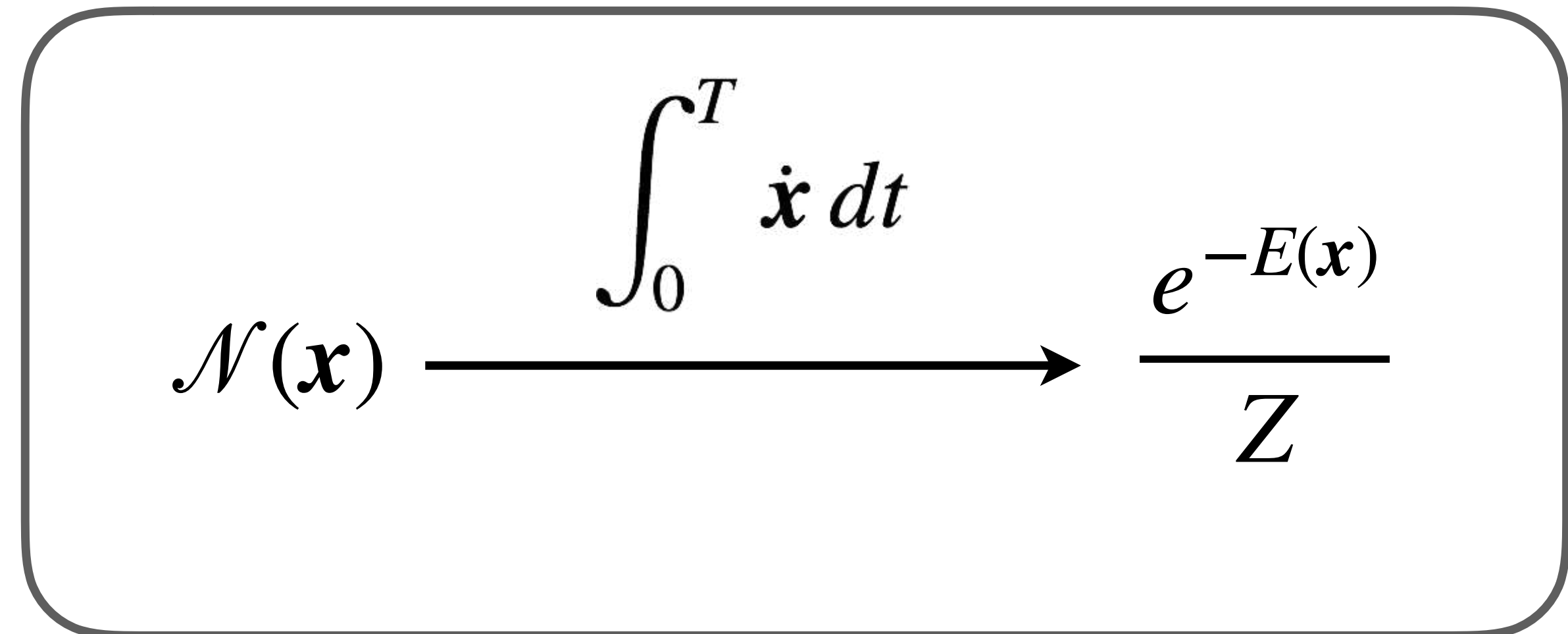
Maximum likelihood estimation



“learn from data”

$$\mathcal{L} = - \mathbb{E}_{\mathbf{x} \sim \text{data}} [\ln p(\mathbf{x})]$$

Variational free energy



“learn from Energy”

$$F = \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x})} [E(\mathbf{x}) + k_B T \ln p(\mathbf{x})]$$

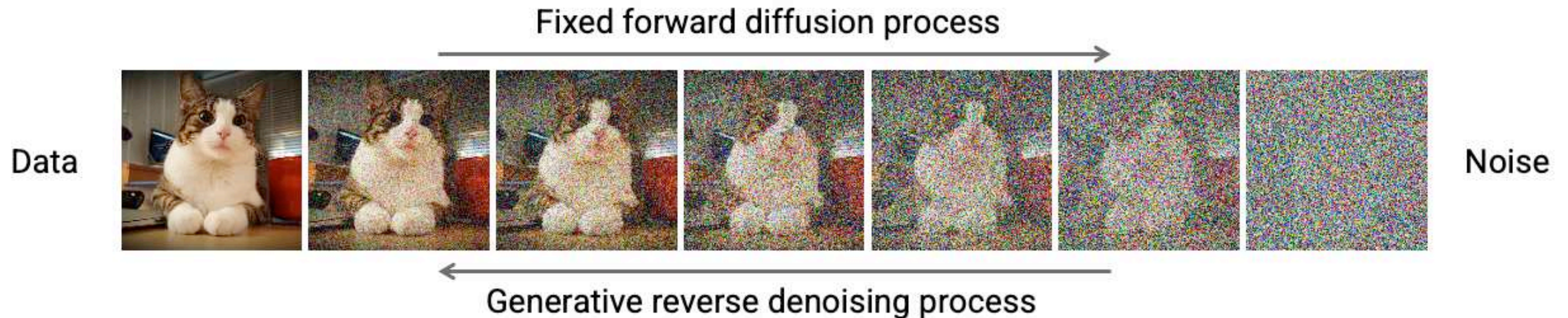
From flow to diffusion model

Continuity equation

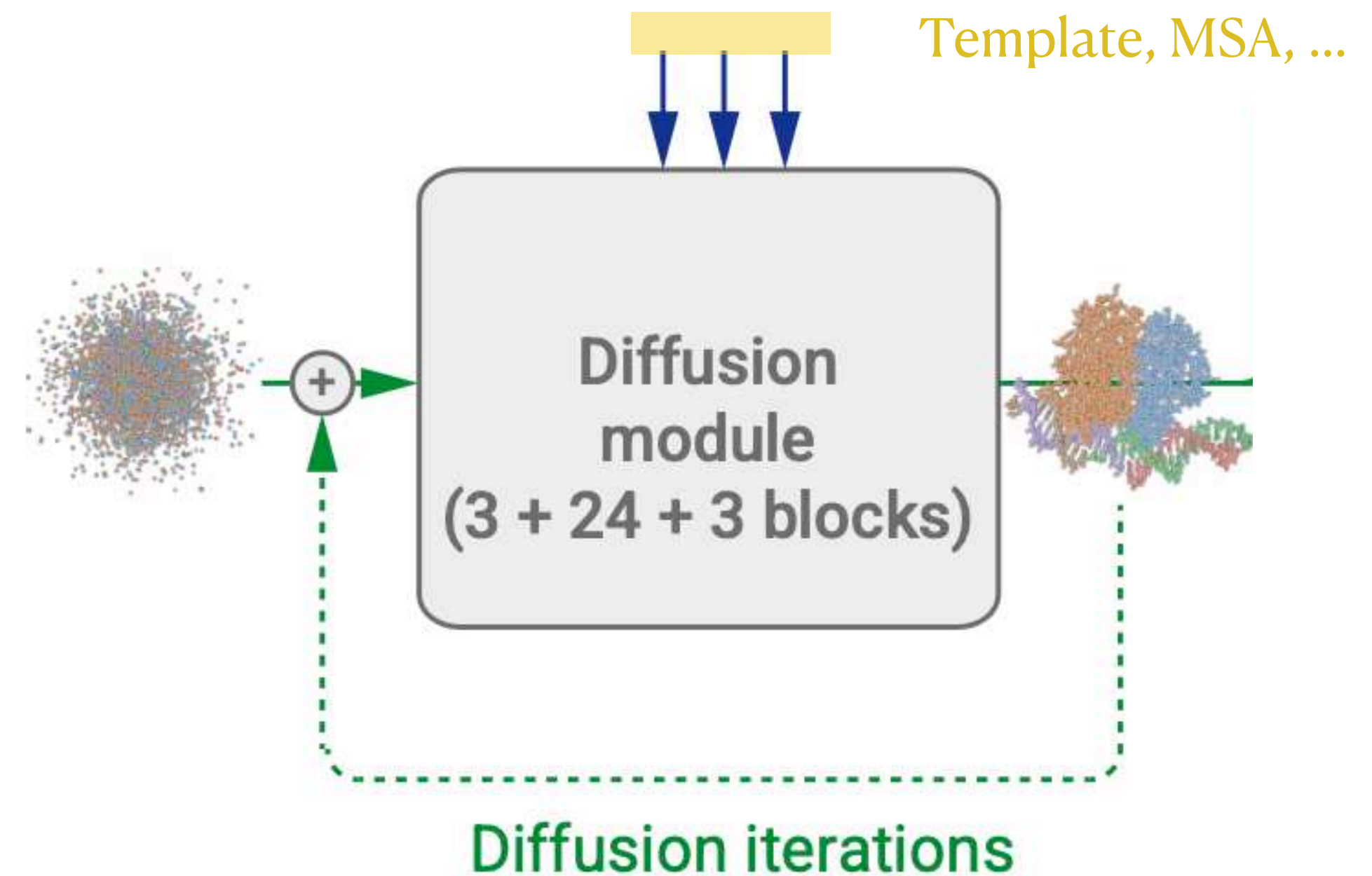
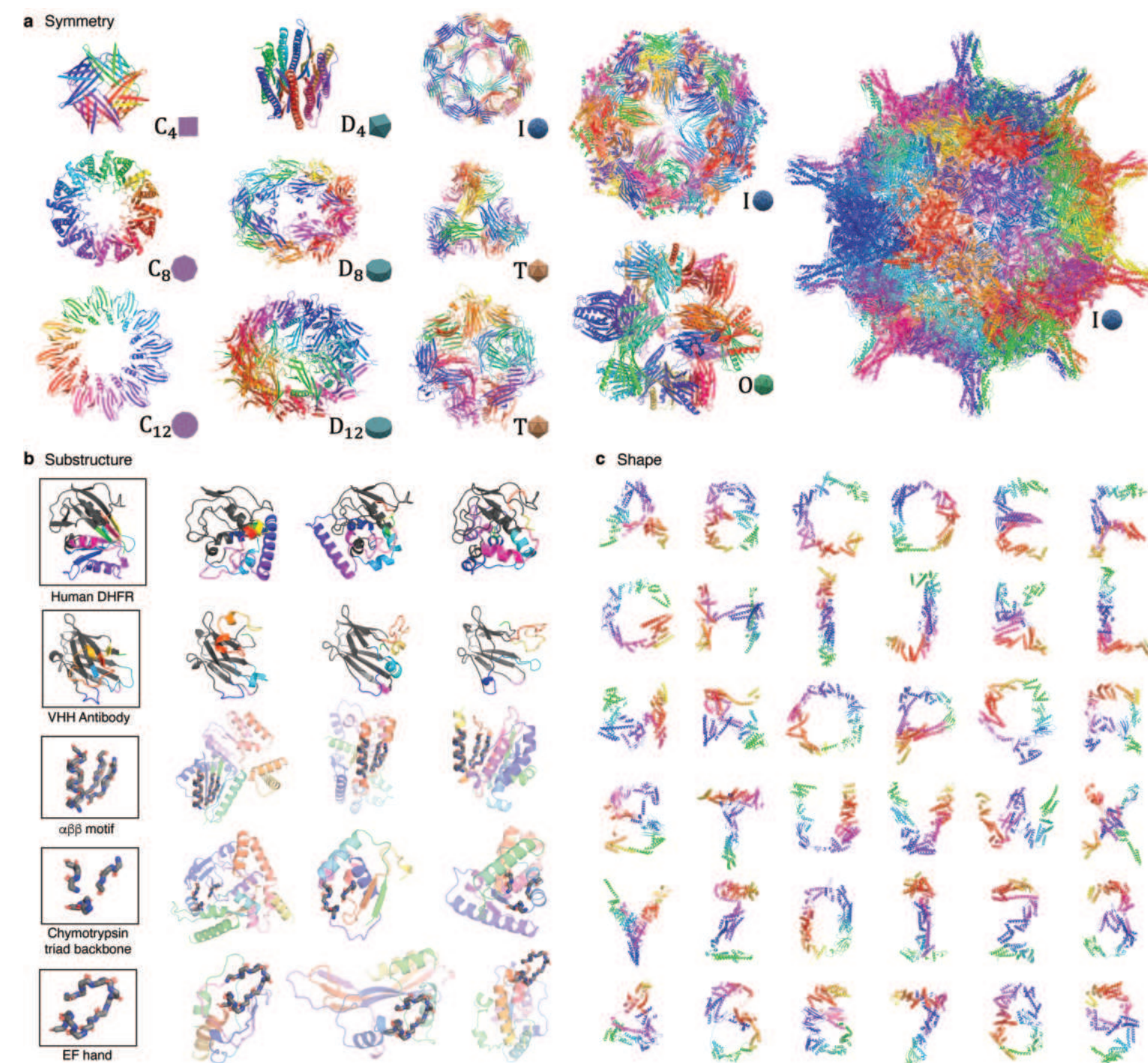
$$\frac{\partial p(\mathbf{X}, t)}{\partial t} + \nabla \cdot [p(\mathbf{X}, t)\mathbf{v}] = 0$$

Fokker-Planck equation

$$\frac{\partial p(\mathbf{X}, t)}{\partial t} + \nabla \cdot [p(\mathbf{X}, t)\mathbf{f}] - \nabla^2 p(\mathbf{X}, t) = 0$$



Diffusion models for protein structure prediction and design



Ingraham et al, Chroma, Nature 2023
<https://generatebiomedicines.com/chroma>

Abramson et al, AlphaFold3, Nature 2024
<https://deepmind.google/technologies/alphafold/>

Do they understand physics?



<https://openai.com/index/sora/>



<http://ai.ruc.edu.cn/newslist/newsdetail/20240326001.html>

“What I can not create, I do not understand”
—Richard Feynman

Do they understand physics?



<https://openai.com/index/sora/>



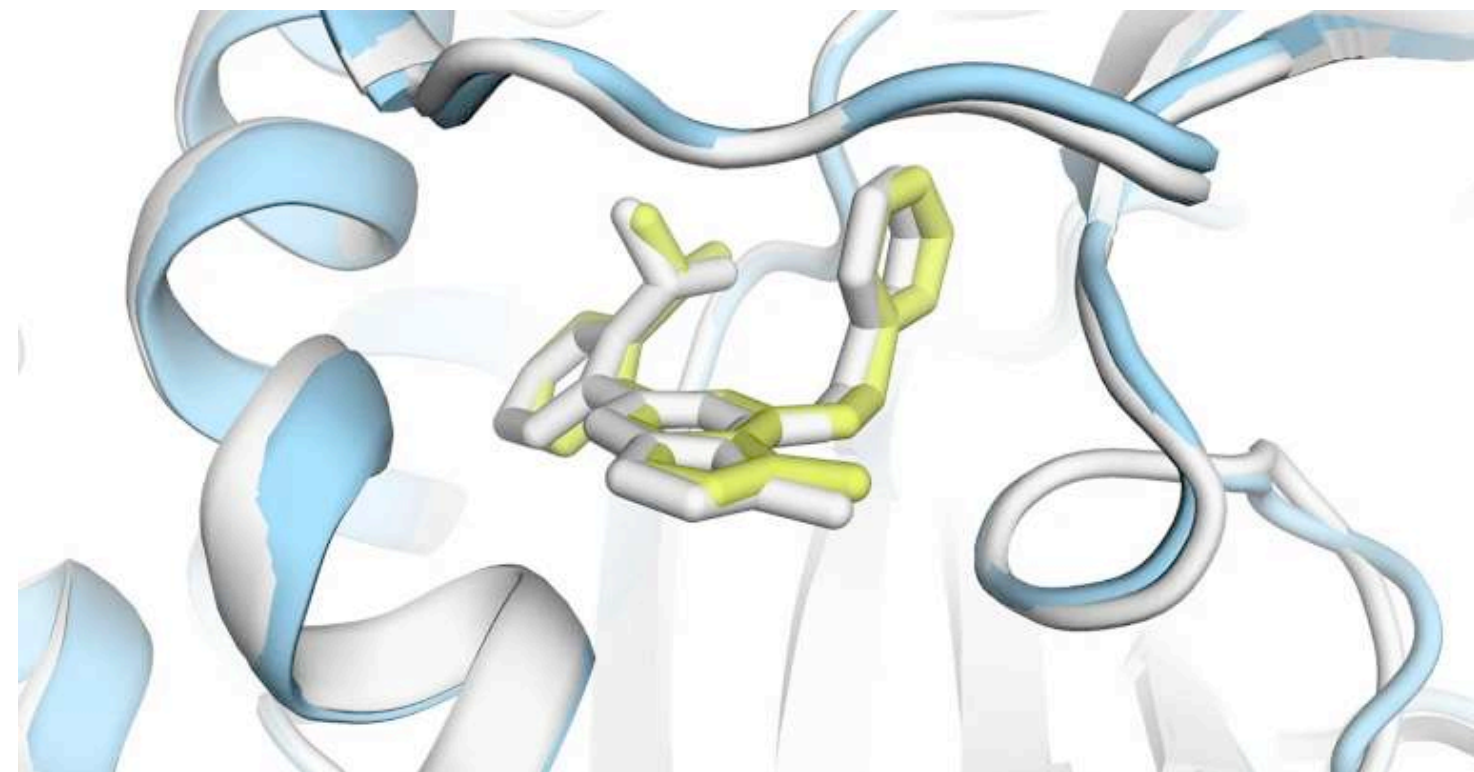
<http://ai.ruc.edu.cn/newslist/newsdetail/20240326001.html>

“What I can not create, I do not understand”
—Richard Feynman

Fold by intuition vs fold by equation

Both integrate Langevin dynamics $X_{t+1} = X_t + \frac{\eta_t}{2}s(X_t, t) + \sqrt{\eta_t}\epsilon$

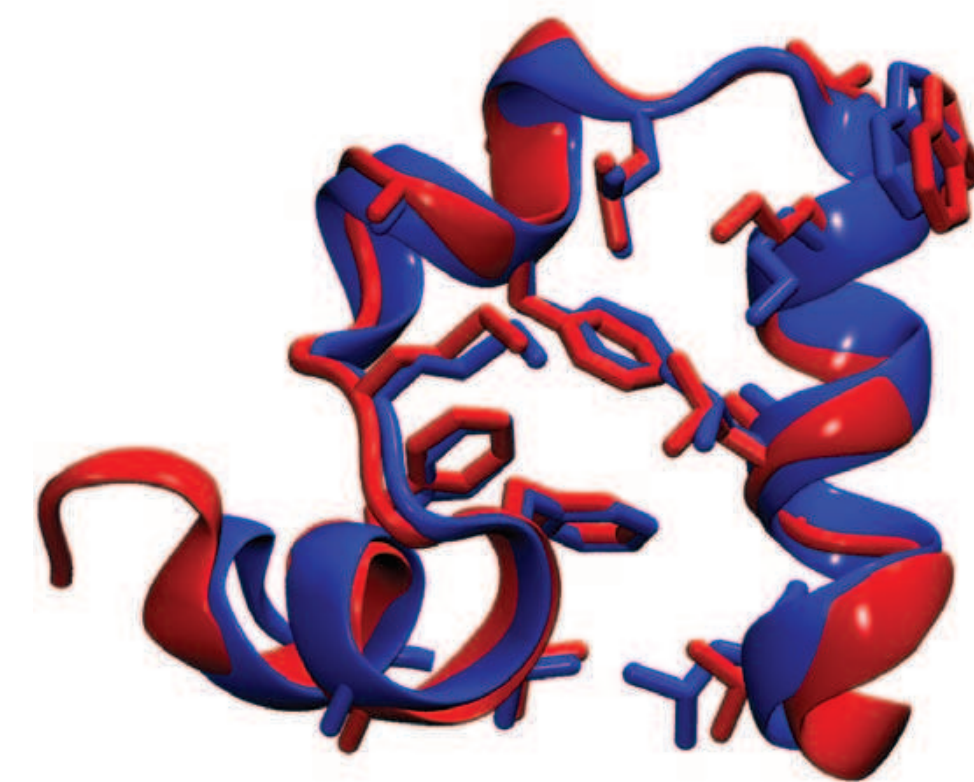
Data-driven generation AlphaFold3



Abramson et al, Nature 2024

The diffusion model may generate right conformations via unphysical pathways

Physics-based molecular dynamics



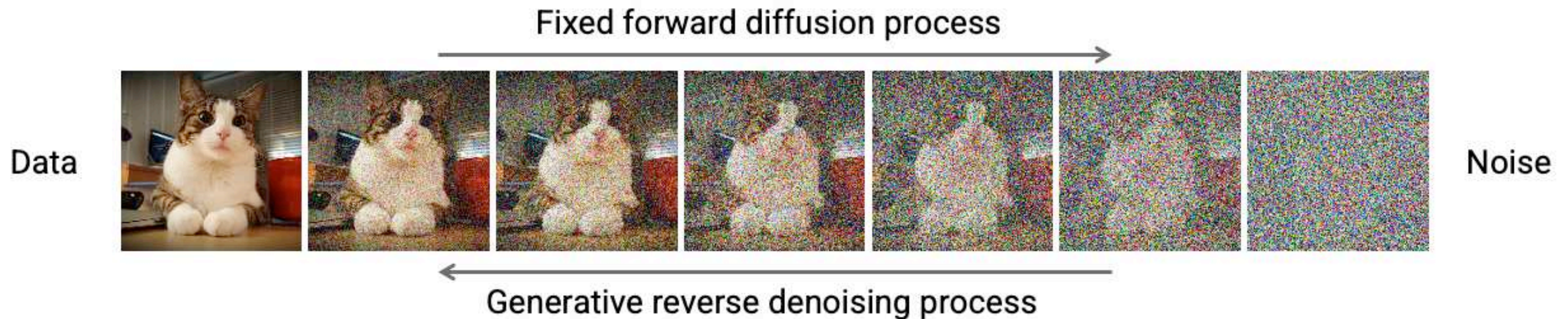
Shaw et al, Science 2010

Physical force fields may face difficulties in sampling rough energy landscapes

From flow to diffusion model, and back

Continuity equation
$$\frac{\partial p(\mathbf{X}, t)}{\partial t} + \nabla \cdot [p(\mathbf{X}, t)\mathbf{v}] = 0$$

Fokker-Planck equation
$$\frac{\partial p(\mathbf{X}, t)}{\partial t} + \nabla \cdot \left[p(\mathbf{X}, t) (\mathbf{f} - \nabla \ln p(\mathbf{X}, t)) \right] = 0$$



Maoutsa et al, 2006.00702, Song et al, 2011.13456

Liu et al 2209.03003, Albergo et al, 2209.15571, Lipman et al, 2210.02747

A tale of three equations

Langevin equation (SDE)

$$\mathbf{x}_{t+dt} = \mathbf{x}_t + \mathbf{f}dt + \sqrt{2dt}\mathcal{N}(0,I)$$

Fokker-Planck equation (PDE)

$$\frac{\partial p(\mathbf{x}, t)}{\partial t} + \nabla \cdot [p(\mathbf{x}, t)\mathbf{f}] - \nabla^2 p(\mathbf{x}, t) = 0$$

“Particle method” (ODE)

$$\frac{d\mathbf{x}}{dt} = \mathbf{f} - \nabla \ln p(\mathbf{x}, t) \equiv \mathbf{v}$$

(Another way to reverse the diffusion is via the reverse-time SDE Anderson 1982)

Maoutsa et al, 2006.00702
Song et al, 2011.13456

A tale of three equations

Langevin equation (SDE)

$$\mathbf{x}_{t+dt} = \mathbf{x}_t + \mathbf{f}dt + \sqrt{2dt} \mathcal{N}(0, I)$$

Fokker-Planck equation (PDE)

$$\frac{\partial p(\mathbf{x}, t)}{\partial t} + \nabla \cdot \left[p(\mathbf{x}, t) (\mathbf{f} - \nabla \ln p(\mathbf{x}, t)) \right] = 0$$

“Particle method” (ODE)

$$\frac{d\mathbf{x}}{dt} = \mathbf{f} - \nabla \ln p(\mathbf{x}, t) \equiv \mathbf{v}$$

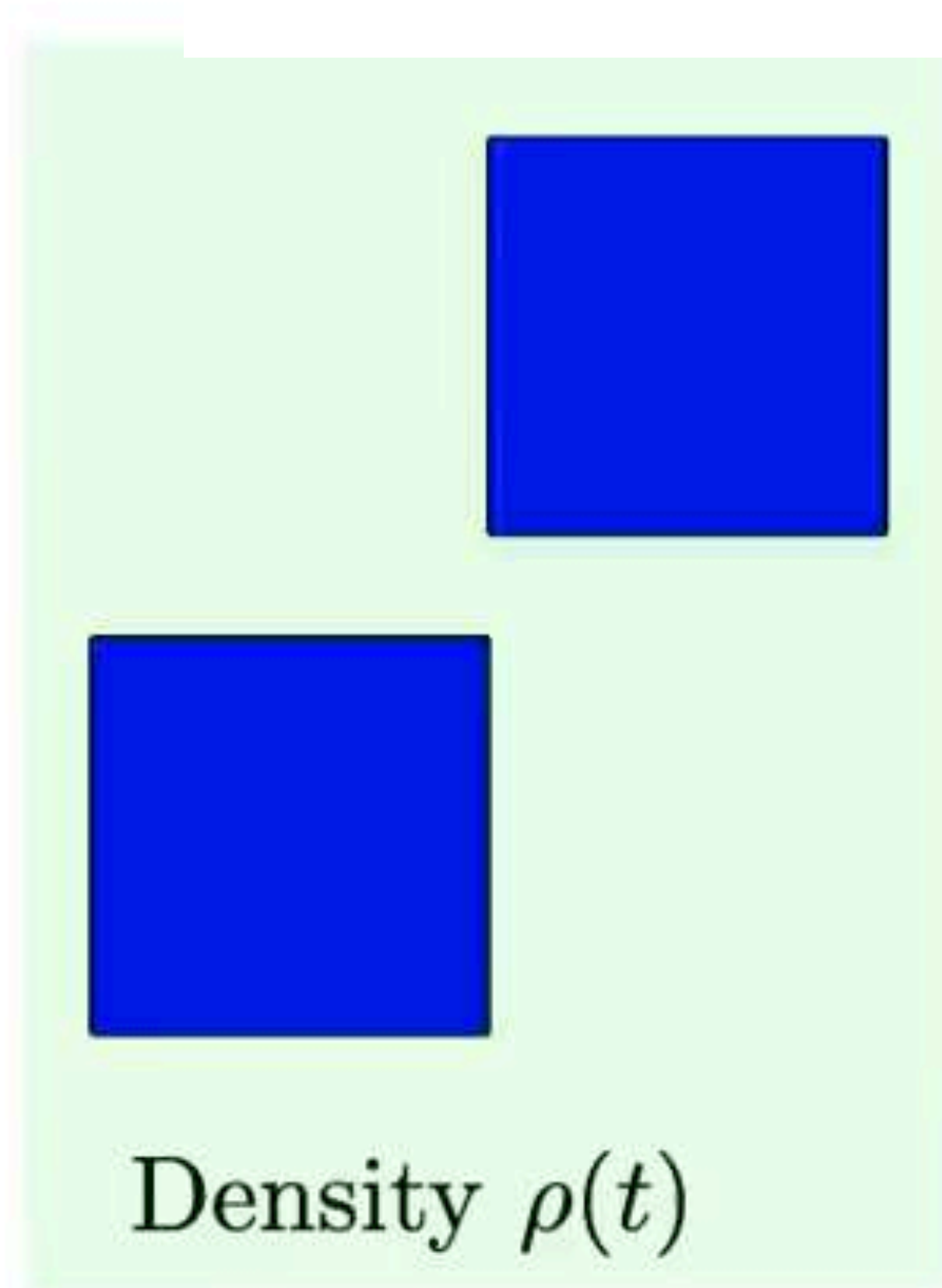
(Another way to reverse the diffusion is via the reverse-time SDE Anderson 1982)

Maoutsa et al, 2006.00702
Song et al, 2011.13456

A tale of three equations

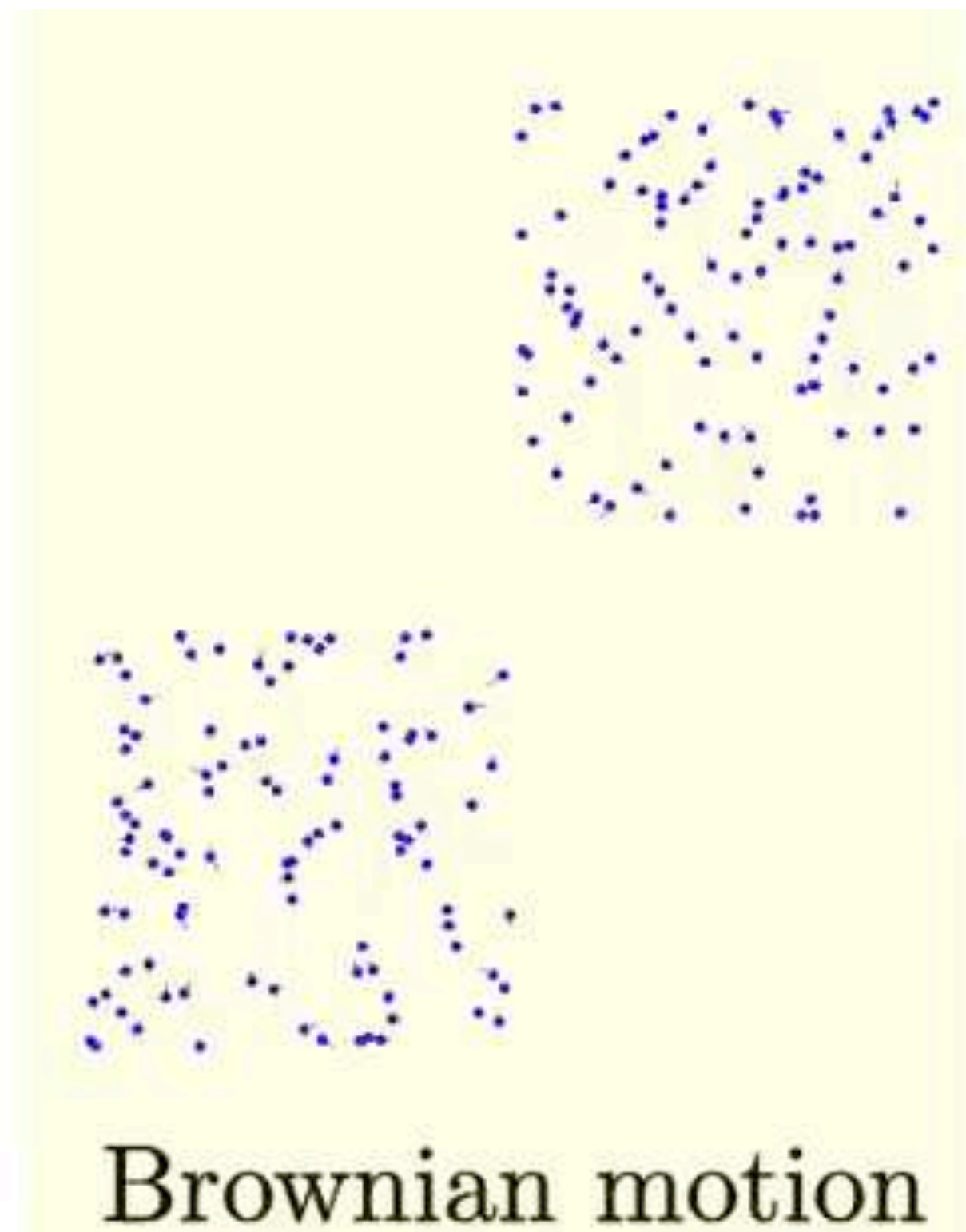
PDE

Fokker-Planck



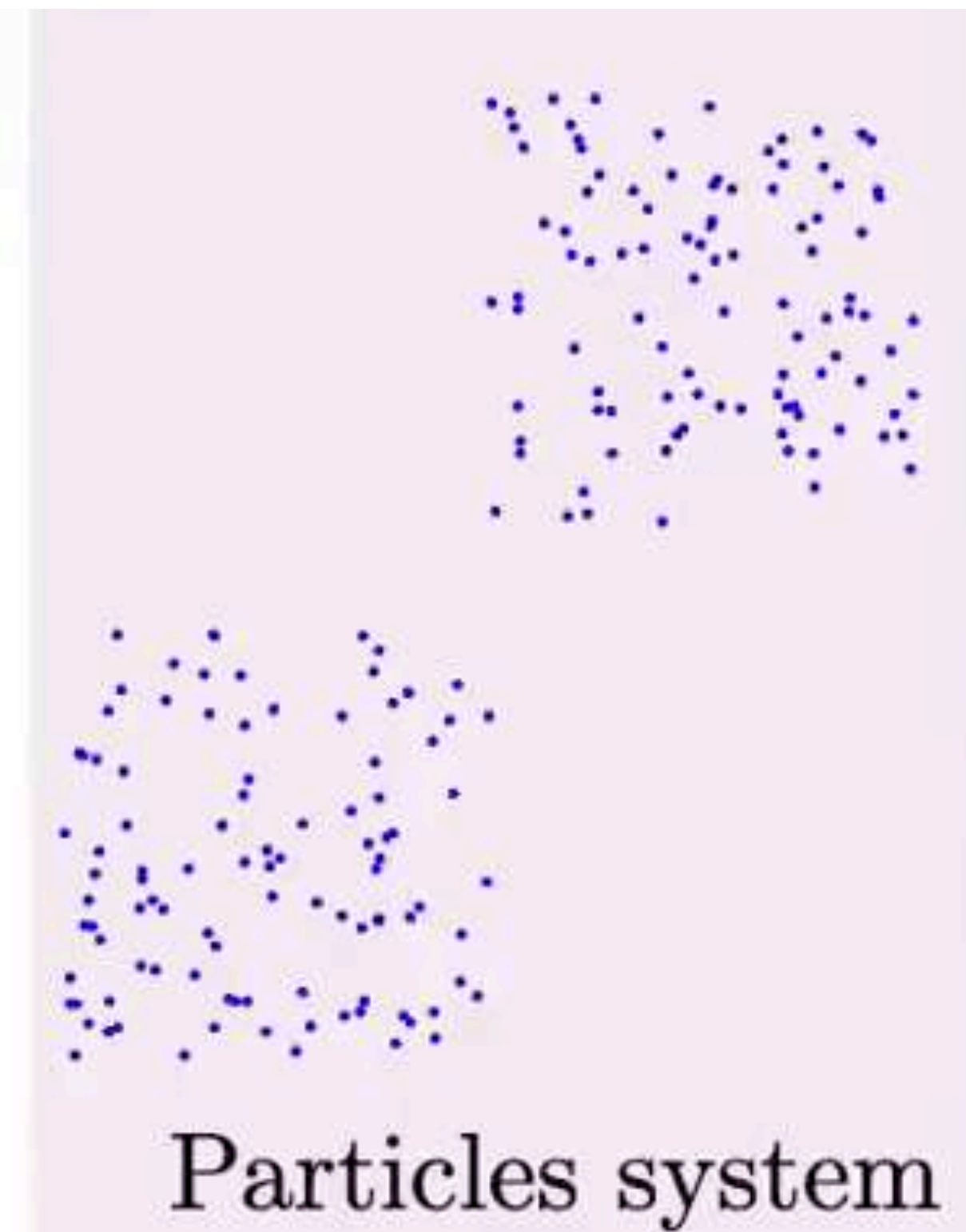
SDE

Diffusion model



ODE

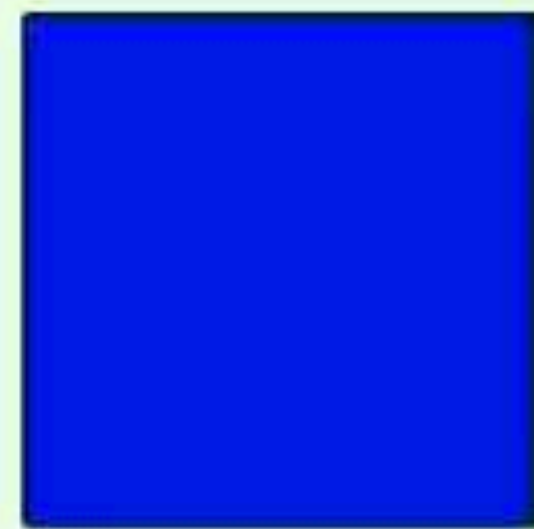
Flow model



A tale of three equations

PDE

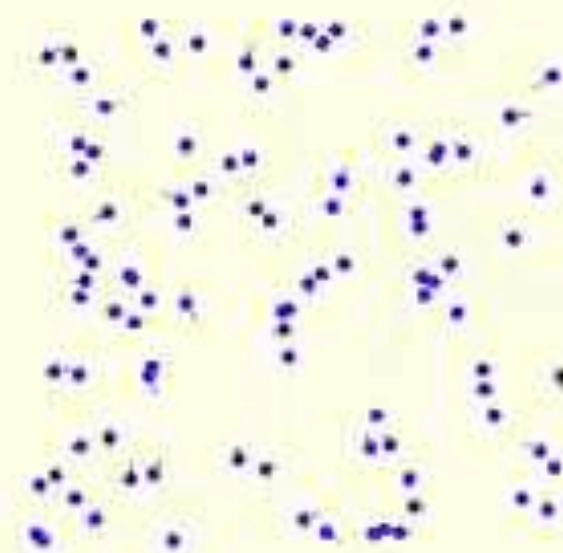
Fokker-Planck



Density $\rho(t)$

SDE

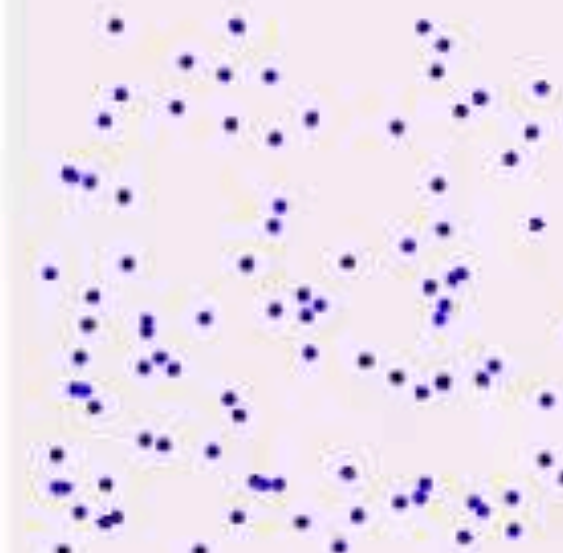
Diffusion model



Brownian motion

ODE

Flow model



Particles system

Lessons from diffusion models

Continuous normalizing flow still has GREAT potential!

Going beyond maximum likelihood estimation training (even if we can)

Break the loss into small pieces, sample them (kind of layer-wise training) <https://blog.alexalemi.com/diffusion.html>

The conditional trick (originated from denoising score matching Vincent 2011)

$$\min_{\theta} \underbrace{\mathbb{E}_{t \sim \mathcal{U}(0, T)}}_{\text{diffusion time } t} \underbrace{\mathbb{E}_{\mathbf{x}_0 \sim q_0(\mathbf{x}_0)}}_{\text{data sample } \mathbf{x}_0} \underbrace{\mathbb{E}_{\mathbf{x}_t \sim q_t(\mathbf{x}_t | \mathbf{x}_0)}}_{\text{diffused data sample } \mathbf{x}_t} \underbrace{\|\mathbf{s}_{\theta}(\mathbf{x}_t, t) - \nabla_{\mathbf{x}_t} \log q_t(\mathbf{x}_t | \mathbf{x}_0)\|_2^2}_{\text{score of diffused data sample}}$$

<https://cvpr2022-tutorial-diffusion-models.github.io/>

Lessons from diffusion models

Continuous normalizing flow still has GREAT potential!

Going beyond maximum likelihood estimation training (even if we can)

Break the loss into small pieces, sample them (kind of layer-wise training)

<https://blog.alexalemi.com/diffusion.html>

The conditional trick (originated from denoising score matching Vincent 2011)

$$\min_{\theta} \underbrace{\mathbb{E}_{t \sim \mathcal{U}(0, T)}}_{\text{diffusion time } t} \underbrace{\mathbb{E}_{\mathbf{x}_t \sim q_t(\mathbf{x}_t)}}_{\text{diffused data } \mathbf{x}_t} \underbrace{\|\mathbf{s}_{\theta}(\mathbf{x}_t, t) - \nabla_{\mathbf{x}_t} \log q_t(\mathbf{x}_t)\|_2^2}_{\text{score of diffused data (marginal)}}$$

<https://cvpr2022-tutorial-diffusion-models.github.io/>

Lessons from diffusion models

Continuous normalizing flow still has GREAT potential!

Going beyond maximum likelihood estimation training (even if we can)

Break the loss into small pieces, sample them (kind of layer-wise training) <https://blog.alexalemi.com/diffusion.html>

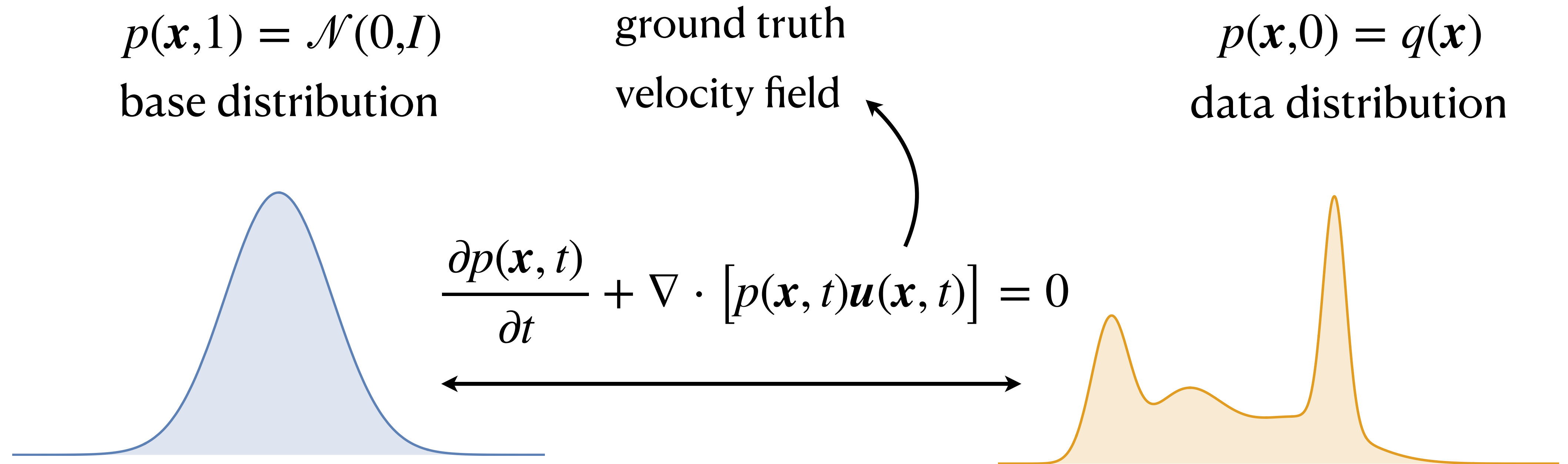
The conditional trick (originated from denoising score matching Vincent 2011)

$$\mathbb{E}_{x \sim q(x)} |s(x) - \nabla_x \ln q(x)|^2 = \mathbb{E}_{x \sim q(x|x_0)} \mathbb{E}_{x_0 \sim q_0(x_0)} |s(x) - \nabla_x \ln q(x|x_0)|^2 + \text{const}.$$

$$q(x) = \int q(x|x_0)q_0(x_0)dx_0$$

Vincent 2011

Flow matching



$$\mathcal{L} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x}, t)} \left| \mathbf{v}_\theta(\mathbf{x}, t) - \mathbf{u}(\mathbf{x}, t) \right|^2$$

Liu et al 2209.03003, Albergo et al, 2209.15571, Lipman et al, 2210.02747

<https://neurips.cc/virtual/2024/tutorial/99531>

The “conditional” trick

Given a conditional continuity equation

$$\frac{\partial p(\mathbf{x} | \mathbf{x}_0, t)}{\partial t} + \nabla \cdot [p(\mathbf{x} | \mathbf{x}_0, t) \mathbf{u}(\mathbf{x} | \mathbf{x}_0, t)] = 0$$

Then, up to a constant, we have

$$\mathcal{L} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x}_0 \sim q(\mathbf{x}_0)} \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x} | \mathbf{x}_0, t)} \left| \mathbf{v}_\theta(\mathbf{x}, t) - \mathbf{u}(\mathbf{x} | \mathbf{x}_0, t) \right|^2$$

We can learn the ground truth velocity by regressing on the conditional velocity

Claim: $\mathcal{L}_{\text{CFM}} = \mathcal{L}_{\text{FM}} + \text{const}.$

where $\mathcal{L}_{\text{FM}} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x},t)} \left| \mathbf{v}_\theta(\mathbf{x}, t) - \mathbf{u}(\mathbf{x}, t) \right|^2$

$\mathcal{L}_{\text{CFM}} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x}_0 \sim q(\mathbf{x}_0)} \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x}|\mathbf{x}_0,t)} \left| \mathbf{v}_\theta(\mathbf{x}, t) - \mathbf{u}(\mathbf{x} | \mathbf{x}_0, t) \right|^2$

$$p(\mathbf{x}, t) = \int p(\mathbf{x} | \mathbf{x}_0, t) q(\mathbf{x}_0) d\mathbf{x}_0 \quad p(\mathbf{x}, t) \mathbf{u}(\mathbf{x}, t) = \int p(\mathbf{x} | \mathbf{x}_0, t) \mathbf{u}(\mathbf{x} | \mathbf{x}_0, t) q(\mathbf{x}_0) d\mathbf{x}_0$$

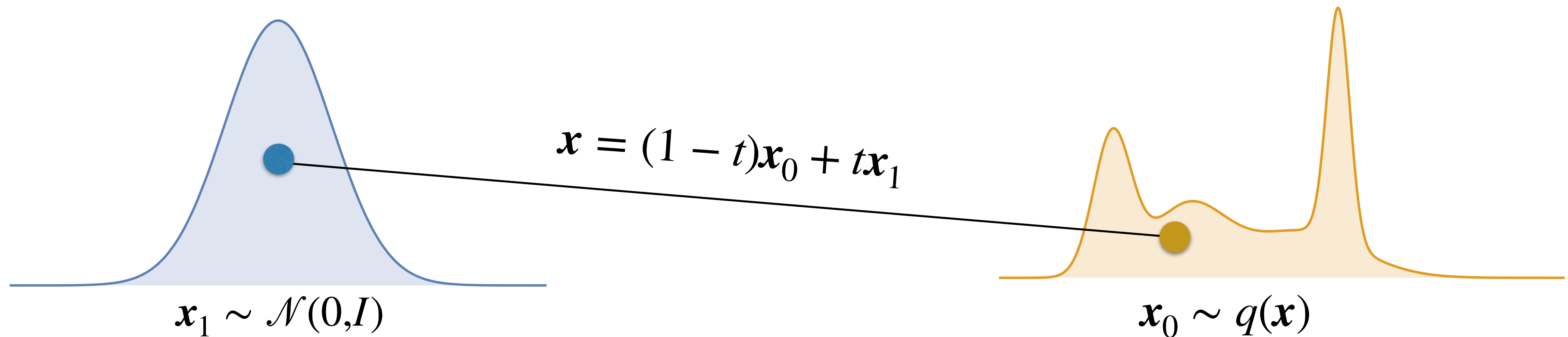
Proof:

$$\mathbb{E}_{\mathbf{x}_0 \sim q(\mathbf{x}_0)} \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x}|\mathbf{x}_0,t)} \left| \mathbf{v}_\theta \right|^2 = \int d\mathbf{x}_0 \int d\mathbf{x} q(\mathbf{x}_0) p(\mathbf{x} | \mathbf{x}_0, t) \left| \mathbf{v}_\theta \right|^2 = \int d\mathbf{x} p(\mathbf{x}, t) \left| \mathbf{v}_\theta \right|^2 = \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x},t)} \left| \mathbf{v}_\theta \right|^2$$

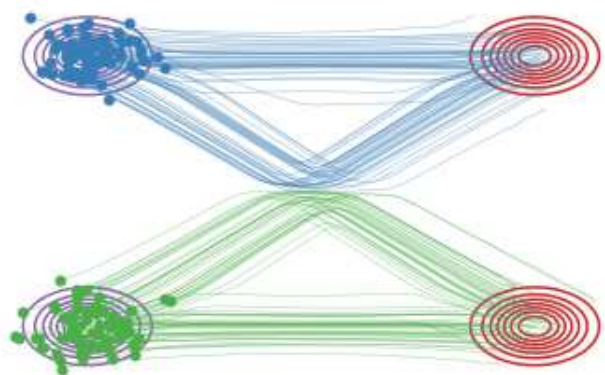
$$\begin{aligned} \mathbb{E}_{\mathbf{x}_0 \sim q(\mathbf{x}_0)} \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x}|\mathbf{x}_0,t)} \left[\mathbf{v}_\theta \cdot \mathbf{u}(\mathbf{x} | \mathbf{x}_0, t) \right] &= \int d\mathbf{x}_0 \int d\mathbf{x} q(\mathbf{x}_0) p(\mathbf{x} | \mathbf{x}_0, t) \left[\mathbf{v}_\theta \cdot \mathbf{u}(\mathbf{x} | \mathbf{x}_0, t) \right] \\ &= \int d\mathbf{x} p(\mathbf{x}, t) \mathbf{v}_\theta \cdot \mathbf{u}(\mathbf{x}, t) = \mathbb{E}_{\mathbf{x} \sim p(\mathbf{x},t)} \left[\mathbf{v}_\theta \cdot \mathbf{u}(\mathbf{x}, t) \right] \end{aligned}$$

Examples of flow matching

$$p(\mathbf{x} | \mathbf{x}_0, t) = \mathcal{N}((1-t)\mathbf{x}_0, t^2) \quad u(\mathbf{x} | \mathbf{x}_0, t) = \frac{d\mathbf{x}}{dt} = \mathbf{x}_1 - \mathbf{x}_0$$



$$\mathcal{L} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x}_0 \sim q(\mathbf{x}_0)} \mathbb{E}_{\mathbf{x}_1 \sim \mathcal{N}(0, I)} \left| \mathbf{v}_\theta(\mathbf{x}, t) - (\mathbf{x}_1 - \mathbf{x}_0) \right|^2$$

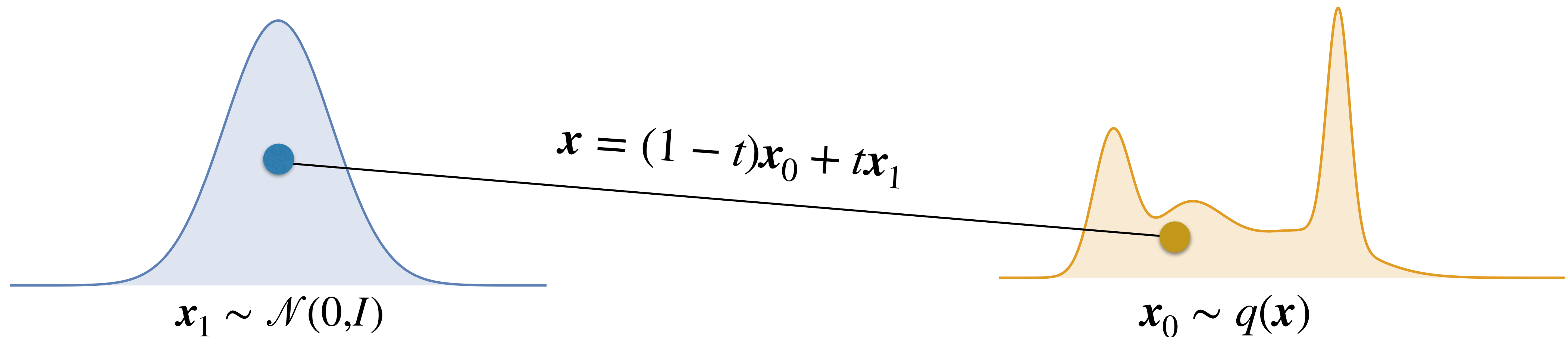


Causalizing linear interpolation with rectified flow 2209.03003

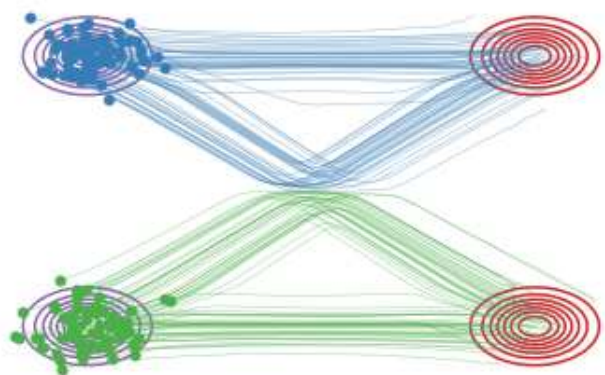
<https://www.cs.utexas.edu/~lqiang/rectflow/html/intro.html>

Examples of flow matching

$$p(\mathbf{x} | \mathbf{x}_0, t) = \mathcal{N}((1-t)\mathbf{x}_0, t^2) \quad u(\mathbf{x} | \mathbf{x}_0, t) = \frac{d\mathbf{x}}{dt} = \mathbf{x}_1 - \mathbf{x}_0$$



$$\mathcal{L} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x}_0 \sim q(\mathbf{x}_0)} \mathbb{E}_{\mathbf{x}_1 \sim \mathcal{N}(0, I)} \left| \mathbf{v}_\theta(\mathbf{x}, t) - (\mathbf{x}_1 - \mathbf{x}_0) \right|^2$$



Causalizing linear interpolation with rectified flow 2209.03003

<https://www.cs.utexas.edu/~lqiang/rectflow/html/intro.html>

Flow matching is all you need!

This framework contains various diffusion models as special cases

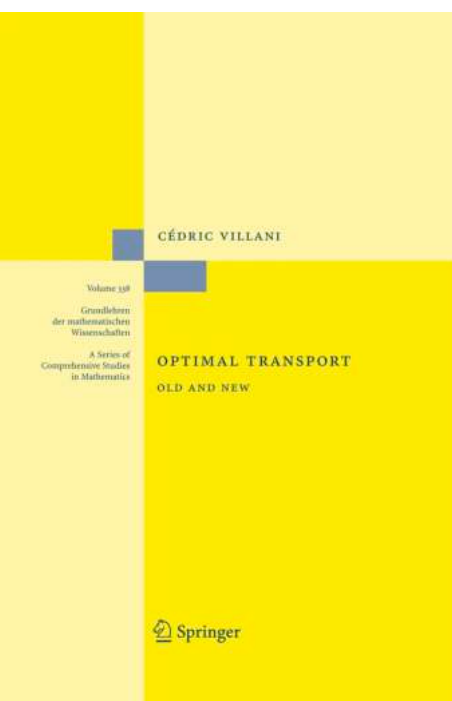
The base distribution does not have to be Gaussian

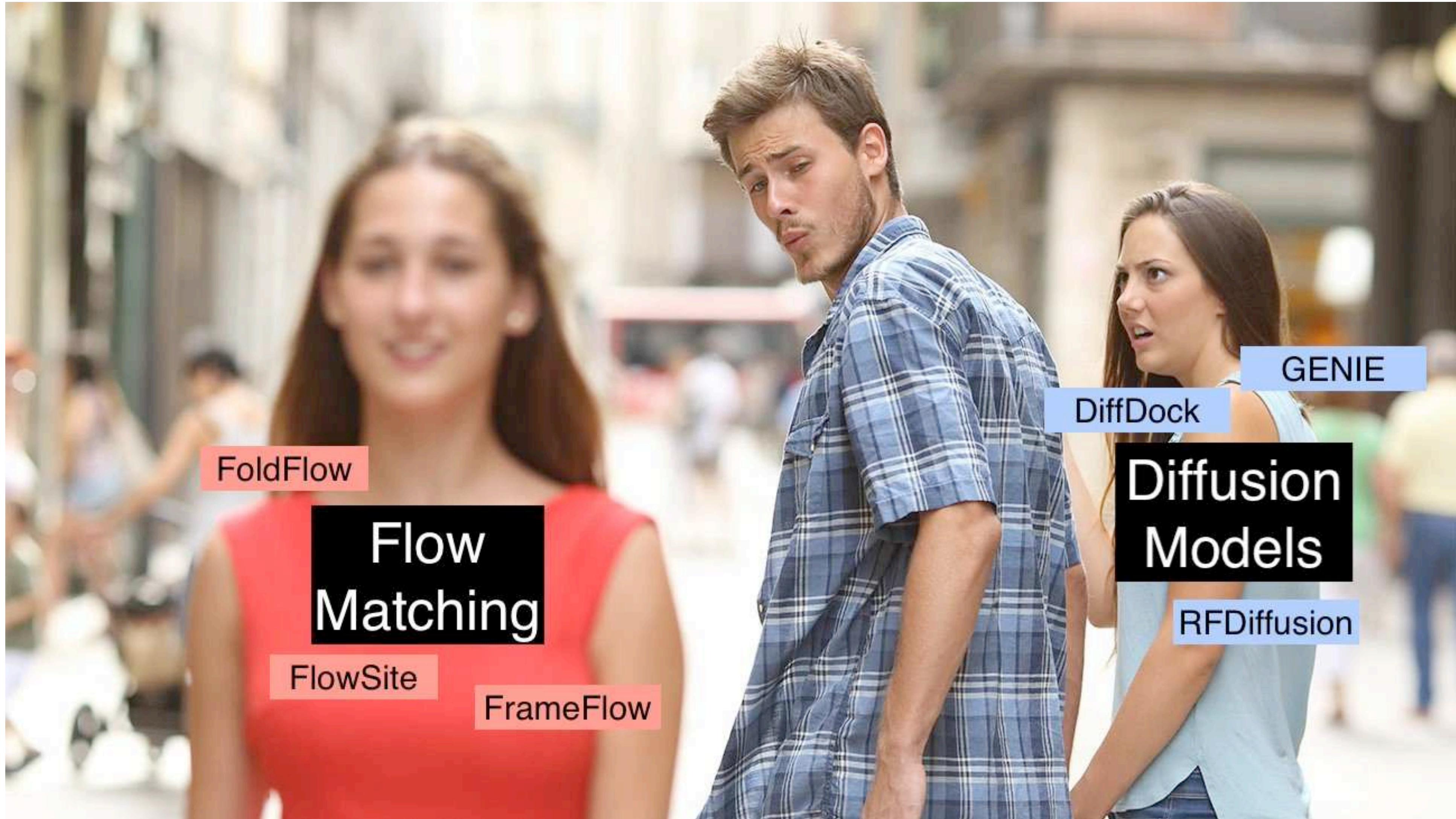
Fast generation with rectified transportation path (Liu et al 2209.03003)

400x speedup compared to continuous normalizing flow (Albergo et al, 2209.15571)

Surpasses diffusion model on Imagenet in likelihood and sample quality
(Lipman et al, 2210.02747)

Generalization to flow on Riemannian manifolds (Chen et al, 2302.03660)



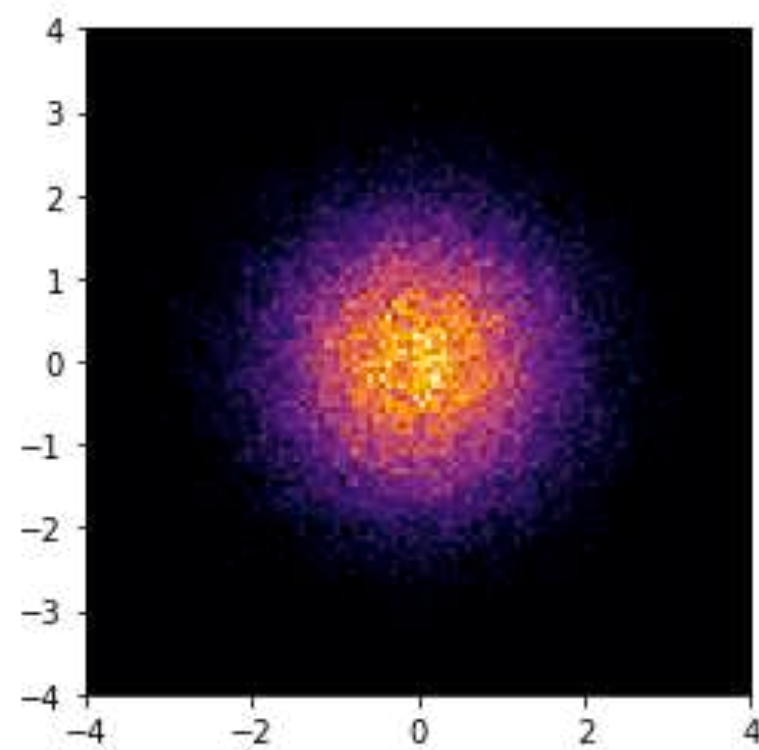


https://twitter.com/michael_galkin/status/1711845455817261409

Demo: bounding free energy of classical Coulomb gas

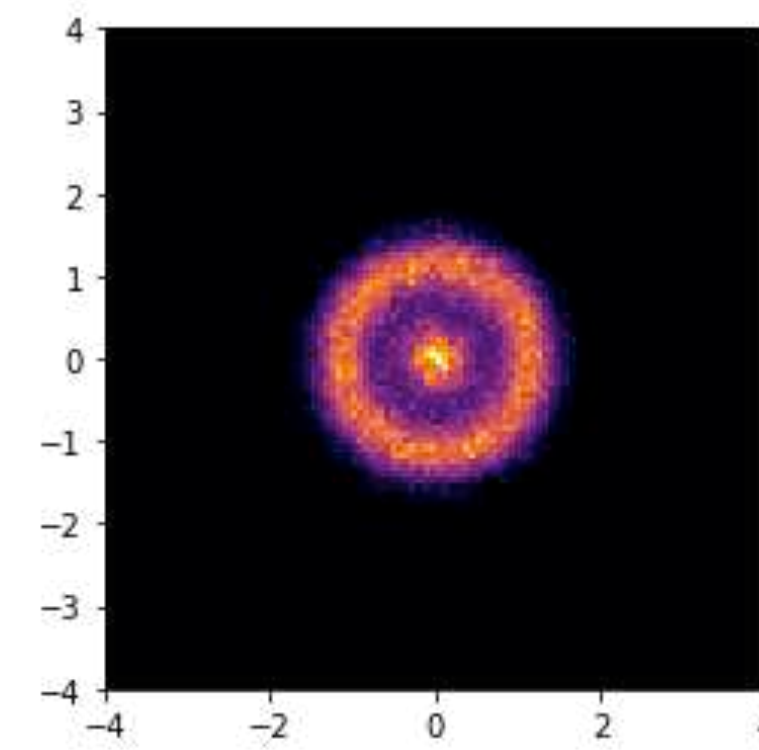
$$\mathcal{L} = \mathbb{E}_{t \sim \mathcal{U}(0,1)} \mathbb{E}_{\mathbf{x}_0 \sim \mathcal{N}(0,I)} \mathbb{E}_{\mathbf{x}_1 \sim \exp(-\beta E)/Z} \left| \mathbf{x}_1 - \mathbf{x}_0 - \mathbf{v}(\mathbf{x}, t) \right|^2$$

$$Z = \mathbb{E}_{\mathbf{x} \sim q(\mathbf{x})} \left[e^{-\beta E(\mathbf{x}) - \ln q(\mathbf{x})} \right] \quad \ln q(\mathbf{x}) = \ln \mathcal{N}(0,I) - \int_0^1 \nabla \cdot \mathbf{v} dt$$



Base density
Gaussian samples

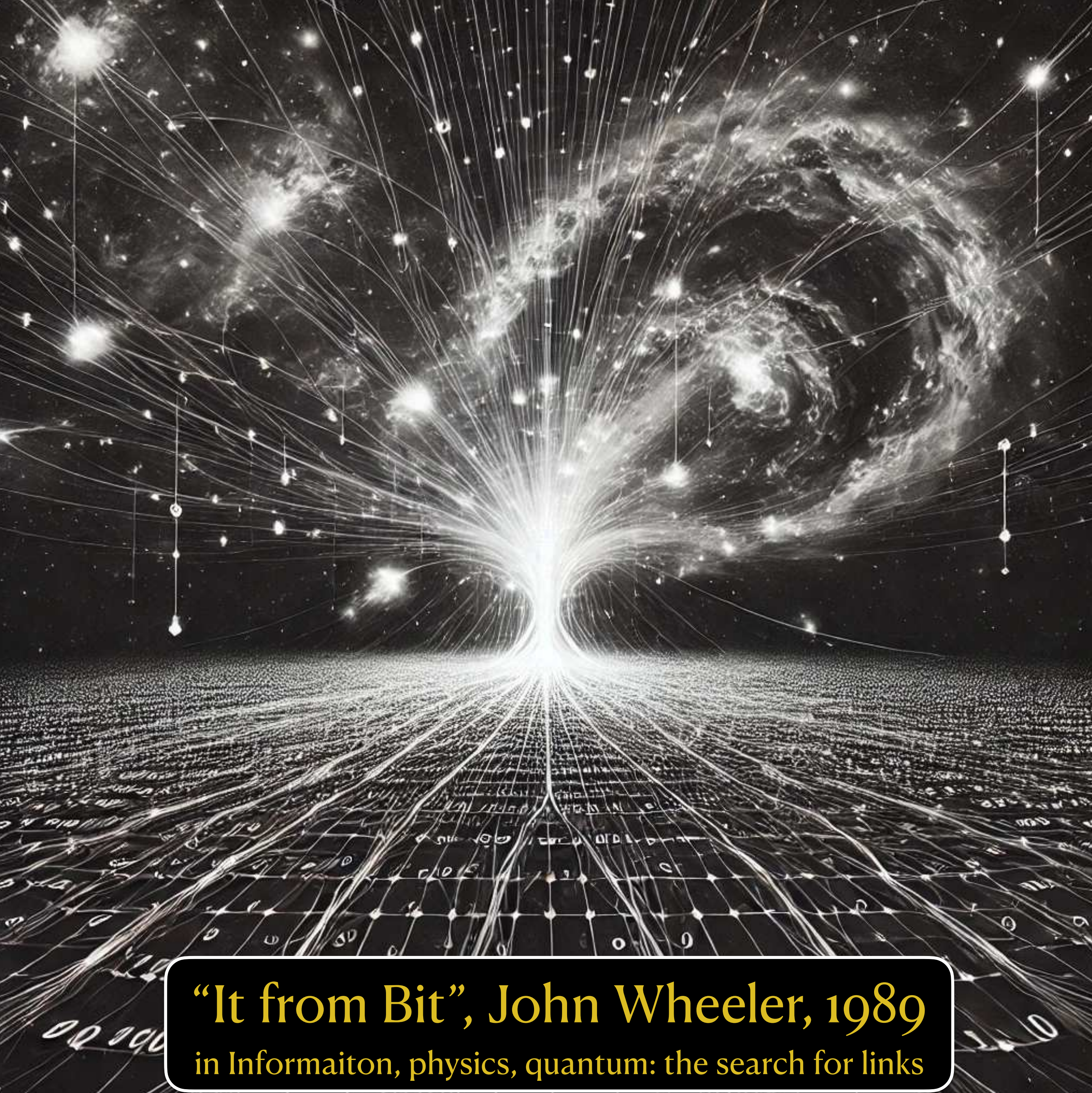
← Interpolate samples to estimate free energy differences →



Target density
Monte Carlo samples

 <https://colab.research.google.com/drive/1t-Vk37Axxp040B7uXFUNlk-zeCC2lcX3?usp=sharing>

Jarzynski PRE '02, see also likelihood-based training of flows Wirnsberger et al, 2002.04913, 2111.08696



Generative AI for **It**

“It from Bit”, John Wheeler, 1989
in Informaiton, physics, quantum: the search for links

Hamiltonian dynamics as symplectic gradient flow

Hamiltonian equations

$$\begin{cases} \dot{p} = -\frac{\partial H}{\partial q} \\ \dot{q} = +\frac{\partial H}{\partial p} \end{cases}$$

Phase space variables

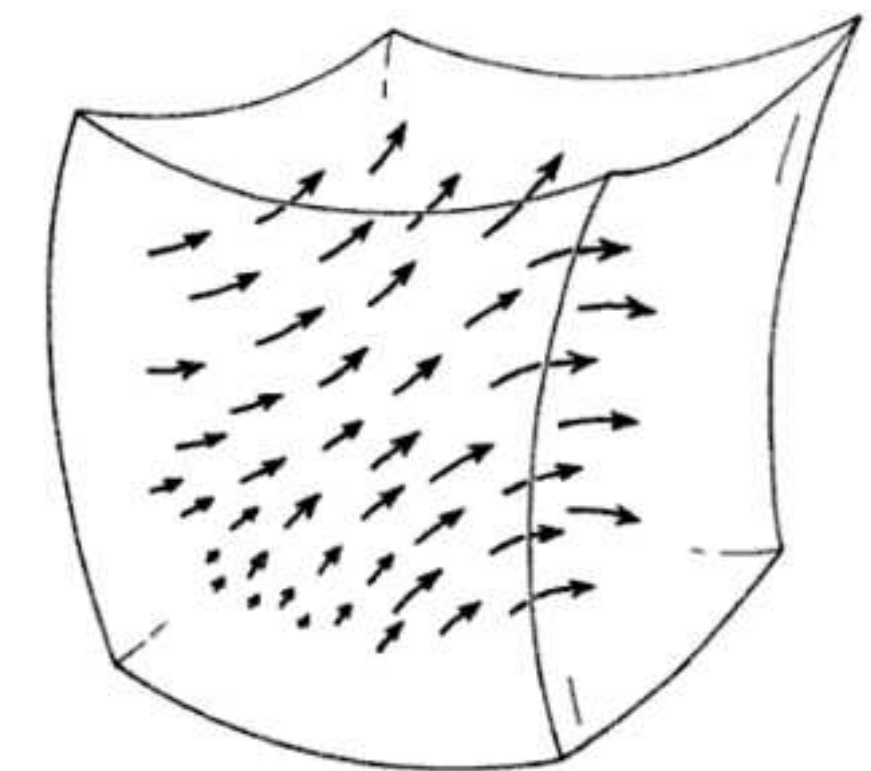
$$\mathbf{x} = (p, q)$$

Symplectic gradient flow

$$\dot{\mathbf{x}} = \nabla_{\mathbf{x}} H(\mathbf{x}) J$$

Symplectic metric

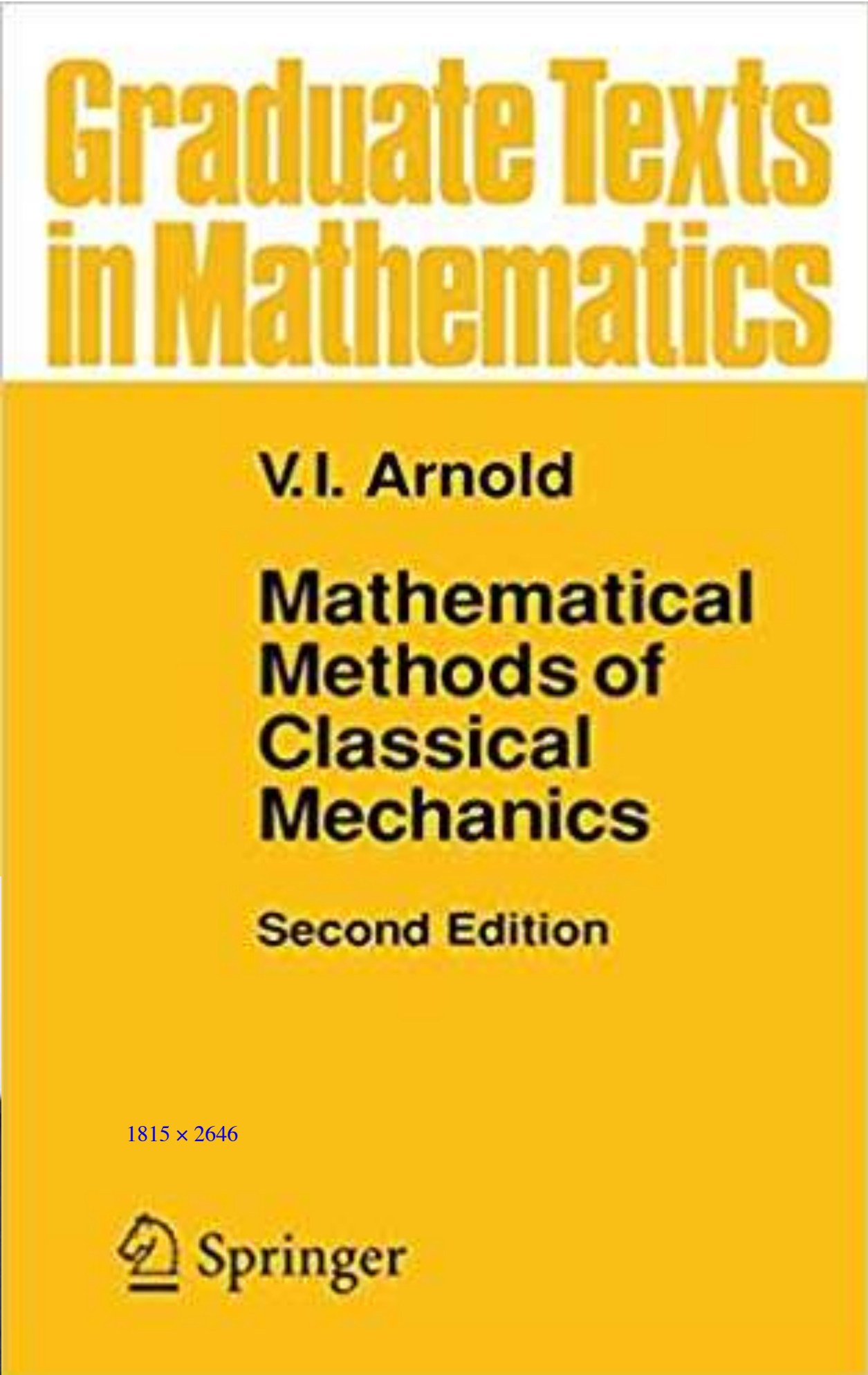
$$J = \begin{pmatrix} & I \\ -I & \end{pmatrix}$$



Hamiltonian dynamics as symplectic gradient flow

Hamiltonian eq

$$\begin{cases} \dot{p} = - \\ \dot{q} = + \end{cases}$$

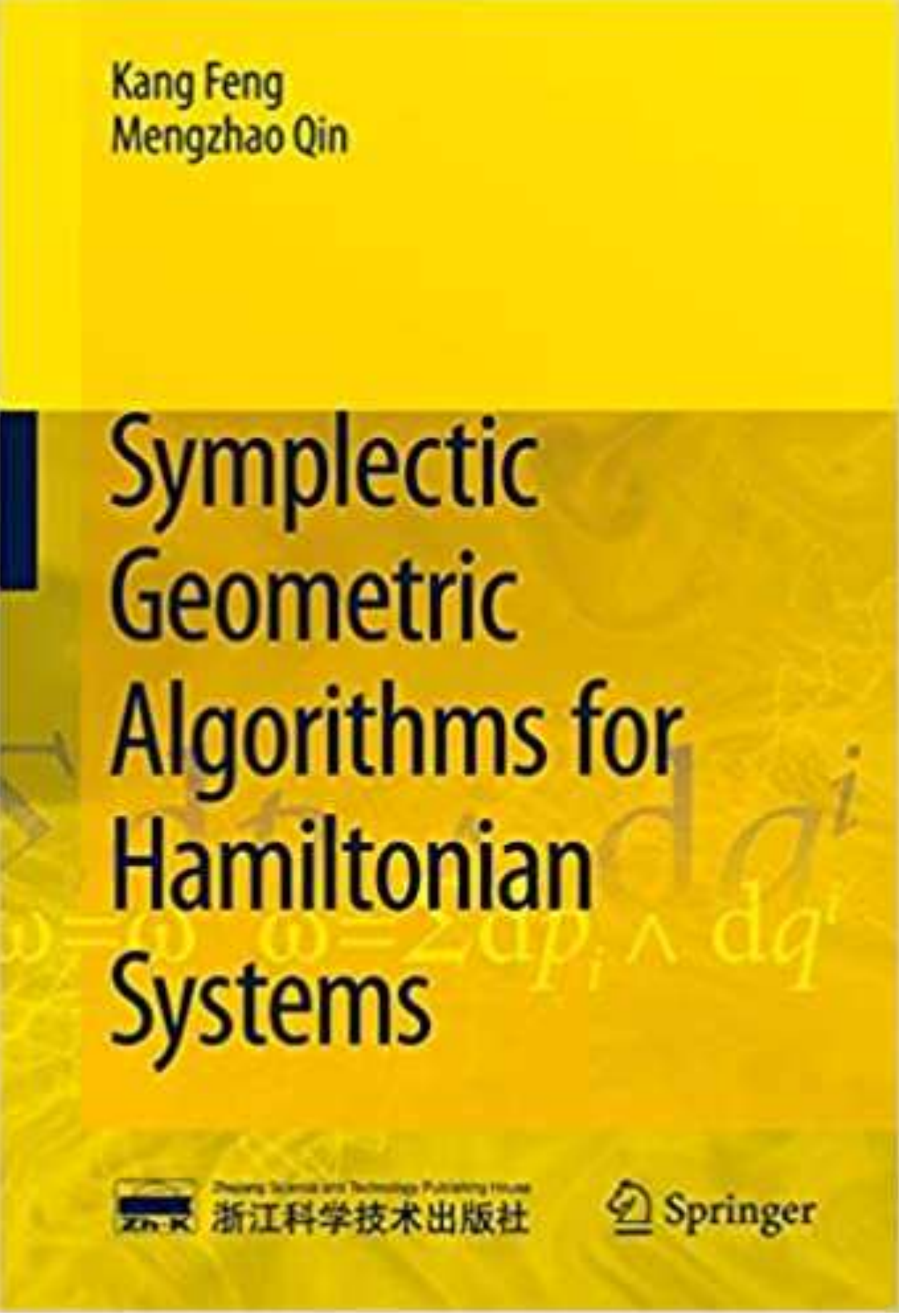


space va

$$= (p, q)$$

lectic m

$$\begin{pmatrix} & I \\ -I & \end{pmatrix}$$



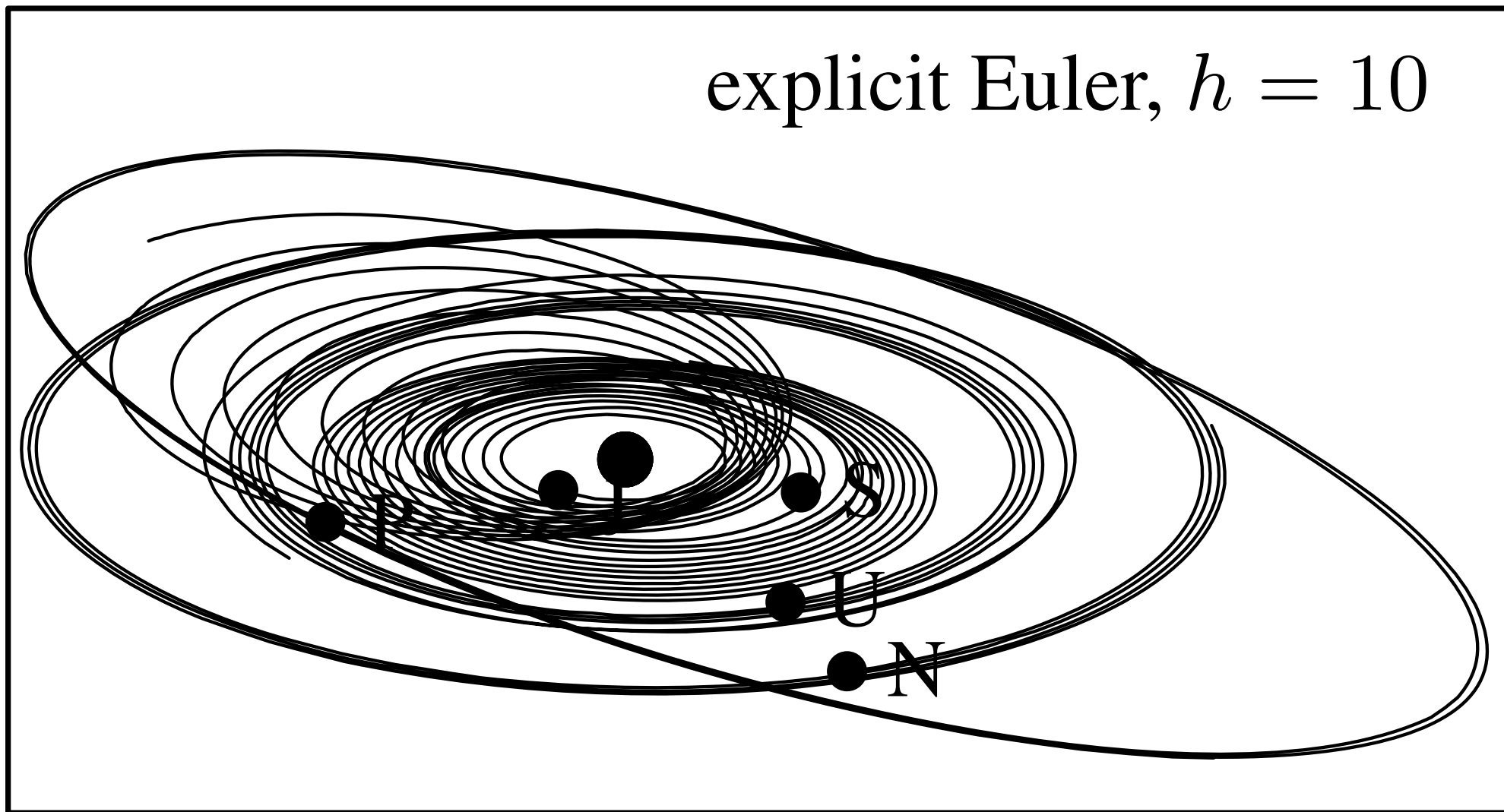
c gradient flow

$$\nabla_x H(x) J$$

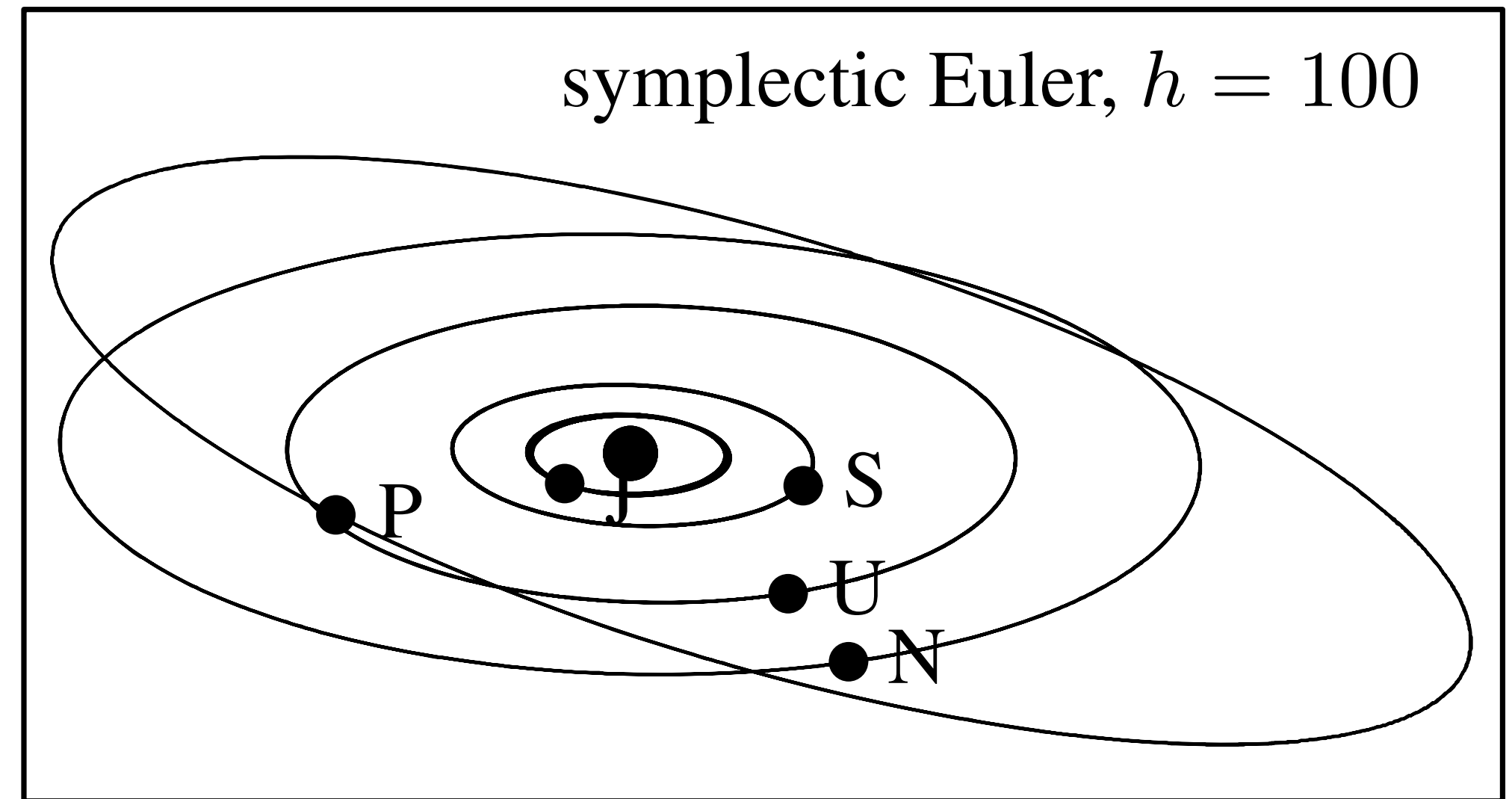


Symplectic Integrators

explicit Euler, $h = 10$



symplectic Euler, $h = 100$



Canonical transformation for Moon-Earth-Sun 3-body problem

634 THÉORIE DU MOUVEMENT DE LA LUNE.

$$\begin{aligned}
 (E_{11}) \quad & + \left(\frac{3}{8} e_1^2 - \frac{3}{4} \gamma_1^2 e_1^2 - \frac{3}{2} e_1^2 - \frac{411}{16} e_1^2 e^2 \right) \frac{n^2}{n_1^2} \\
 & + \left(\frac{219}{64} e_1^2 - \frac{99}{4} \gamma_1^2 e_1^2 - \frac{619}{32} e_1^2 - \frac{9843}{128} e_1^2 e^2 \right) \frac{n^2}{n_1^2} \\
 & + \frac{189}{128} e_1^2 \frac{n^2}{n_1^2} - \frac{65337}{1024} e_1^2 \frac{n^2}{n_1^2} - \frac{5}{64} e_1^2 \frac{n^2}{n_1^2} \cdot \frac{a_1^2}{a^2} \cos \theta_1 (t + \epsilon) \\
 & - \frac{99}{128} e_1^2 \frac{n^2}{n_1^2} \cos 2 \theta_1 (t + \epsilon), \\
 \theta = \theta_1 (t + \epsilon) \\
 (E_{21}) \quad & - \left[\left(\frac{3}{4} - \frac{3}{2} \gamma_1^2 + \frac{3}{8} e_1^2 - \frac{15}{8} e^2 + \frac{3}{4} \gamma_1^2 + \frac{15}{4} \gamma_1^2 e^2 - \frac{171}{64} e_1^2 - \frac{15}{16} e_1^2 e^2 \right) \frac{n^2}{n_1^2} \right. \\
 & + \left(\frac{3}{8} - \frac{3}{4} \gamma_1^2 + \frac{21}{16} e_1^2 - \frac{411}{16} e^2 \right) \frac{n^2}{n_1^2} \\
 & + \left(\frac{219}{64} - \frac{99}{4} \gamma_1^2 + \frac{1399}{128} e_1^2 - \frac{9843}{128} e^2 \right) \frac{n^2}{n_1^2} \\
 & \left. + \frac{189}{128} \frac{n^2}{n_1^2} - \frac{65229}{1024} \frac{n^2}{n_1^2} - \frac{5}{64} \frac{n^2}{n_1^2} \cdot \frac{a_1^2}{a^2} \right] \sin \theta_1 (t + \epsilon) \\
 & + \left[\left(\frac{9}{64} - \frac{9}{16} \gamma_1^2 - \frac{45}{128} e_1^2 - \frac{45}{64} e^2 \right) \frac{n^2}{n_1^2} + \frac{9}{64} \frac{n^2}{n_1^2} + \frac{675}{512} \frac{n^2}{n_1^2} \right] \sin 2 \theta_1 (t + \epsilon) \\
 & - \frac{9}{256} \frac{n^2}{n_1^2} \sin 3 \theta_1 (t + \epsilon), \\
 a = a_1 \left\{ 1 + \left[\left(\frac{2}{3} e_1^2 - 3 \gamma_1^2 e_1^2 - \frac{15}{4} e_1^2 - \frac{15}{4} e_1^2 e^2 + \frac{3}{2} \gamma_1^2 e_1^2 + \frac{15}{2} \gamma_1^2 e_1^2 \right. \right. \right. \\
 & \left. \left. + \frac{15}{2} \gamma_1^2 e_1^2 e^2 + \frac{101}{32} e_1^2 + \frac{75}{8} e_1^2 e^2 \right) \frac{n^2}{n_1^2} \right. \right. \\
 & \left. + \left(\frac{3}{4} e_1^2 - \frac{3}{2} \gamma_1^2 e_1^2 - \frac{15}{8} e_1^2 - \frac{411}{8} e_1^2 e^2 \right) \frac{n^2}{n_1^2} \right. \\
 & \left. + \left(\frac{219}{32} e_1^2 - \frac{99}{2} \gamma_1^2 e_1^2 - \frac{1819}{64} e_1^2 - \frac{9843}{64} e_1^2 e^2 \right) \frac{n^2}{n_1^2} \right. \\
 & \left. + \frac{189}{64} e_1^2 \frac{n^2}{n_1^2} - \frac{27249}{512} e_1^2 \frac{n^2}{n_1^2} - \frac{5}{32} e_1^2 \frac{n^2}{n_1^2} \cdot \frac{a_1^2}{a^2} \right] \cos \theta_1 (t + \epsilon) \\
 & - \frac{9}{16} e_1^2 \frac{n^2}{n_1^2} \cos 2 \theta_1 (t + \epsilon) \left. \right\}, \\
 (G_{11}) \quad & \gamma = \gamma_1 - \left[\left(\frac{3}{8} \gamma_1^2 e_1^2 - \frac{3}{4} \gamma_1^2 e_1^2 - \frac{3}{4} \gamma_1^2 e_1^2 - \frac{15}{16} \gamma_1^2 e_1^2 e^2 \right) \frac{n^2}{n_1^2} \right. \\
 & \left. + \frac{3}{16} \gamma_1^2 e_1^2 \frac{n^2}{n_1^2} + \frac{219}{128} \gamma_1^2 e_1^2 \frac{n^2}{n_1^2} \right] \cos \theta_1 (t + \epsilon), \\
 (H_{11}) \quad & \gamma = \gamma_1 - \left[\left(\frac{3}{8} \gamma_1^2 e_1^2 - \frac{3}{4} \gamma_1^2 e_1^2 - \frac{3}{4} \gamma_1^2 e_1^2 - \frac{15}{16} \gamma_1^2 e_1^2 e^2 \right) \frac{n^2}{n_1^2} \right. \\
 & \left. + \frac{3}{16} \gamma_1^2 e_1^2 \frac{n^2}{n_1^2} + \frac{219}{128} \gamma_1^2 e_1^2 \frac{n^2}{n_1^2} \right] \cos \theta_1 (t + \epsilon).
 \end{aligned}$$

640 THÉORIE DU MOUVEMENT DE LA LUNE.

$$\begin{aligned}
 & + \left(\frac{13}{64} + \frac{187}{32} \gamma^2 - \frac{237}{128} e^2 + \frac{195}{128} e^2 - \frac{1389}{32} \gamma^2 - \frac{599}{64} \gamma^2 e^2 + \frac{2805}{64} \gamma^2 e^2 \right. \\
 & \left. - \frac{103173}{1024} e^2 - \frac{3105}{256} e^2 e^2 \right) \frac{n^2}{n^2} \\
 & + \left(\frac{79}{16} + \frac{55}{48} \gamma^2 - \frac{1063}{48} e^2 + \frac{2133}{32} e^2 \right) \frac{n^2}{n^2} + \left(\frac{153}{8} + \frac{3245}{96} \gamma^2 - \frac{73159}{768} e^2 + \frac{240085}{512} e^2 \right) \frac{n^2}{n^2} \\
 & + \frac{22441}{288} \frac{n^2}{n^2} + \frac{99916415}{442368} \frac{n^2}{n^2} + \frac{4431}{2048} \frac{n^2}{n^2} \cdot \frac{a^2}{n^2} \left. \right\} \\
 \text{De ces valeurs de L, G, H, on déduit} \\
 \frac{da}{dt} = \frac{1}{an} \left\{ 2 + \left(\frac{1969}{32} - \frac{1629}{8} \gamma^2 + \frac{34985}{128} e^2 + \frac{28635}{64} e^2 \right) \frac{n^2}{n^2} \right. \\
 & \left. + \left(\frac{415}{2} - \frac{2745}{4} \gamma^2 + \frac{31449}{16} e^2 + \frac{43299}{16} e^2 \right) \frac{n^2}{n^2} + \frac{61185}{64} \frac{n^2}{n^2} + \frac{1532167}{576} \frac{n^2}{n^2} \right\}, \\
 \frac{dG}{dt} = -\frac{1}{an} \left\{ \left(\frac{527}{8} - \frac{3633}{16} \gamma^2 - \frac{9991}{128} e^2 + 480 e^2 \right) \frac{n^2}{n^2} \right. \\
 & \left. + \left(\frac{2757}{8} - \frac{2493}{2} \gamma^2 - \frac{7161}{16} e^2 + \frac{36459}{8} e^2 \right) \frac{n^2}{n^2} + \frac{104117}{64} \frac{n^2}{n^2} + \frac{277537}{48} \frac{n^2}{n^2} \right\}, \\
 \frac{dH}{dt} = -\frac{1}{an} \left\{ \left(\frac{15}{16} + \frac{15}{16} \gamma^2 - \frac{1809}{32} e^2 + \frac{225}{32} e^2 \right) \frac{n^2}{n^2} \right. \\
 & \left. + \left(\frac{167}{8} - 66 \gamma^2 - \frac{2625}{8} e^2 + \frac{4509}{16} e^2 \right) \frac{n^2}{n^2} + \frac{895}{16} \frac{n^2}{n^2} + \frac{176531}{576} \frac{n^2}{n^2} \right\}, \\
 \frac{de}{dt} = \frac{1}{a^2 n e} \left\{ 1 - e^2 + \left(\frac{1991}{64} - \frac{1113}{16} \gamma^2 - \frac{40571}{128} e^2 + \frac{28065}{128} e^2 \right) \frac{n^2}{n^2} + \frac{3323}{24} \frac{n^2}{n^2} + \frac{62483}{96} \frac{n^2}{n^2} \right\}, \\
 \frac{d\gamma}{dt} = -\frac{1}{a^2 n e} \left\{ 1 - \frac{1}{2} e^2 - \frac{1}{8} e^2 - \frac{1}{16} e^2 \right. \\
 & \left. + \left(\frac{1901}{64} - \frac{1113}{16} \gamma^2 - \frac{3831}{8} e^2 + \frac{28065}{128} e^2 \right) \frac{n^2}{n^2} + \frac{3323}{24} \frac{n^2}{n^2} + \frac{62483}{96} \frac{n^2}{n^2} \right\}, \\
 \frac{dL}{dt} = \frac{1}{a^2 n e} \frac{141}{8} e^2 \frac{n^2}{n^2}, \\
 \frac{d\gamma}{dt} = \frac{1}{a^2 n \gamma} \frac{183}{32} \gamma^2 \frac{n^2}{n^2},
 \end{aligned}$$



Charles Delaunay



More than 1800 pages of this, ~20 years of efforts (1846-1867)

How to find useful canonical transformations for more complex systems?

Canonical transformations and deep learning

Change of variables

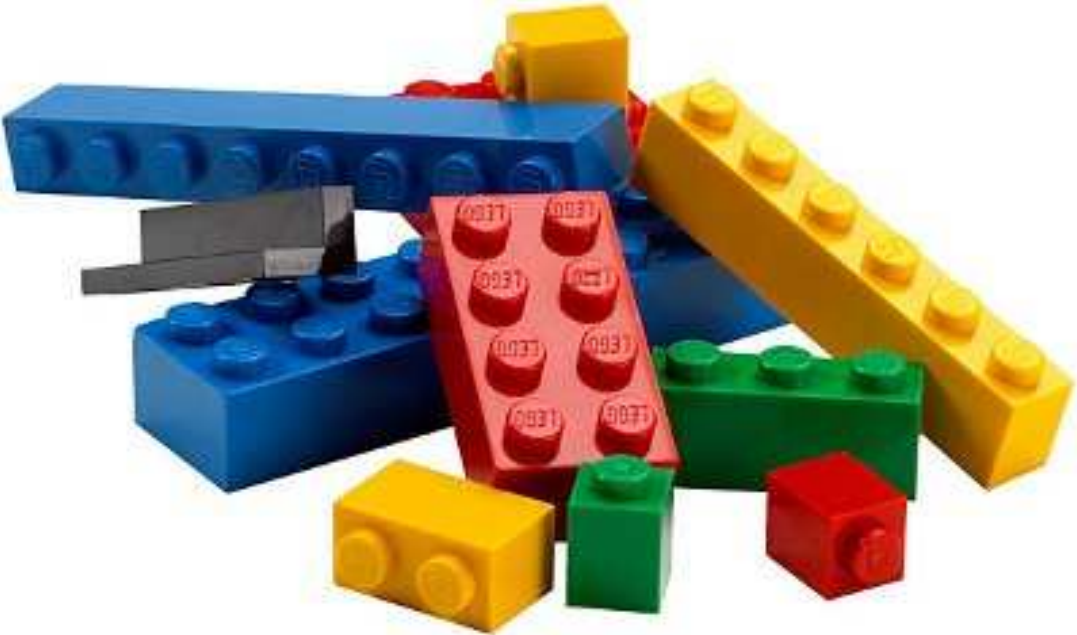
$$X = (p, q) \xleftrightarrow{\hspace{10em}} Z = (P, Q)$$

$$(\nabla_X Z) \begin{pmatrix} & I \\ -I & \end{pmatrix} (\nabla_X Z)^T = \begin{pmatrix} & I \\ -I & \end{pmatrix}$$

symplectic condition

$$Z = \mathcal{T}(X)$$

$$\mathcal{T} = \mathcal{T}_1 \circ \mathcal{T}_2 \circ \mathcal{T}_3 \circ \dots$$



Compose symplectic layers to form a deep neural network
and learn them either from data or variationally without data

Canonical transformations and generative models

Canonical transformation deforms **phase space density** $p(X) = e^{-\beta H(X)}$

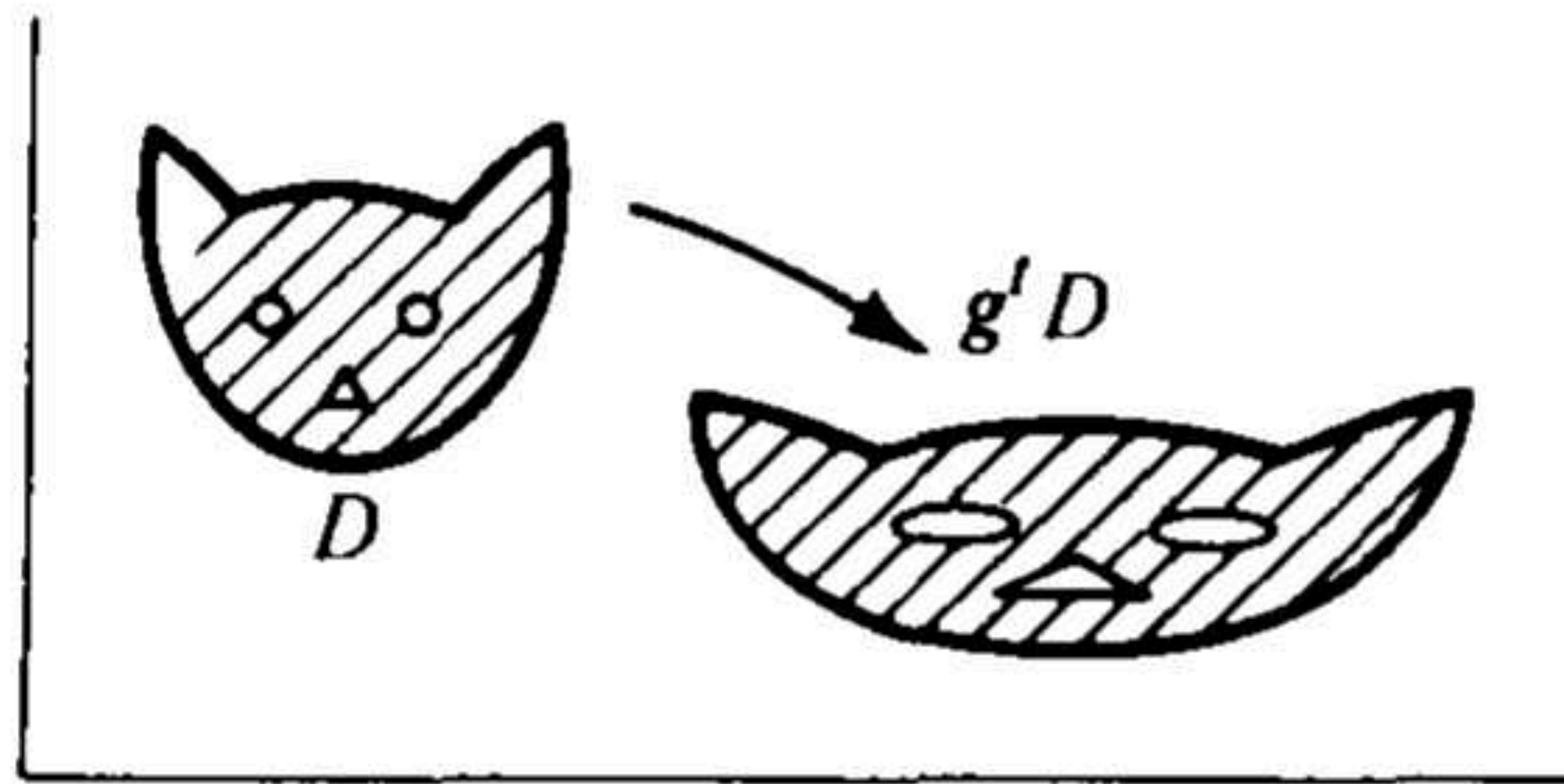
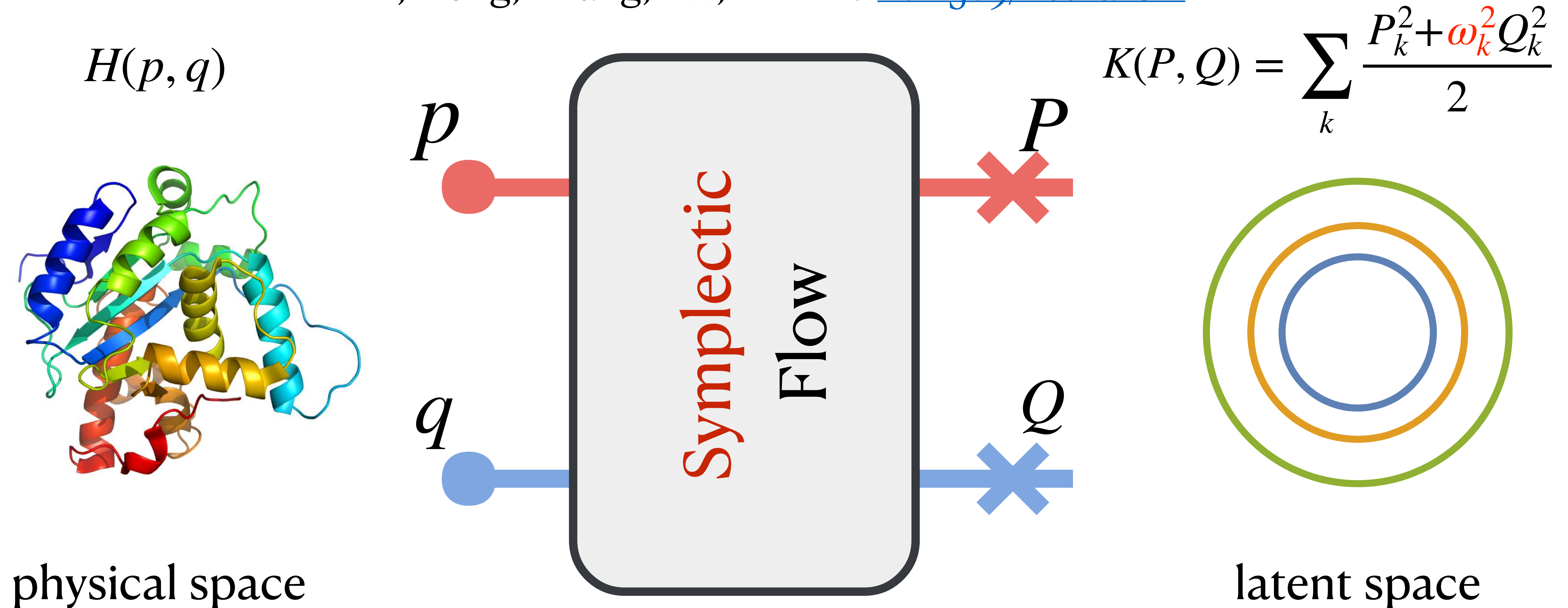


Image from Arnold '78

Flow models are great at transforming probability densities

Neural Canonical Transformations

Li, Dong, Zhang, LW, PRX '20 [li012589/neuralCT](https://arxiv.org/abs/1912.05890)

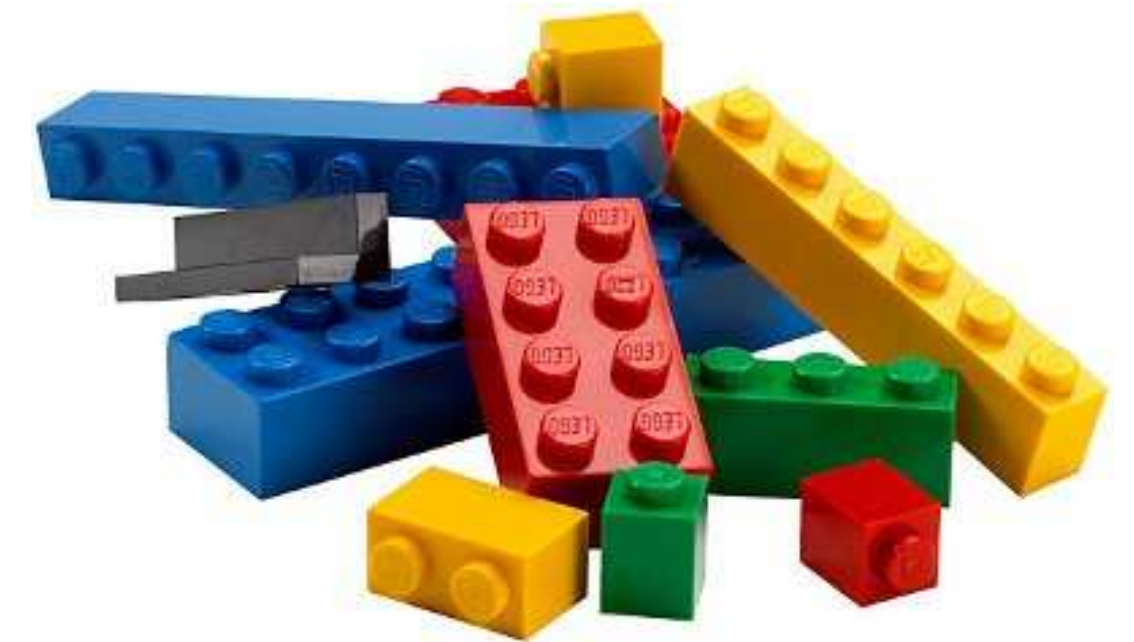
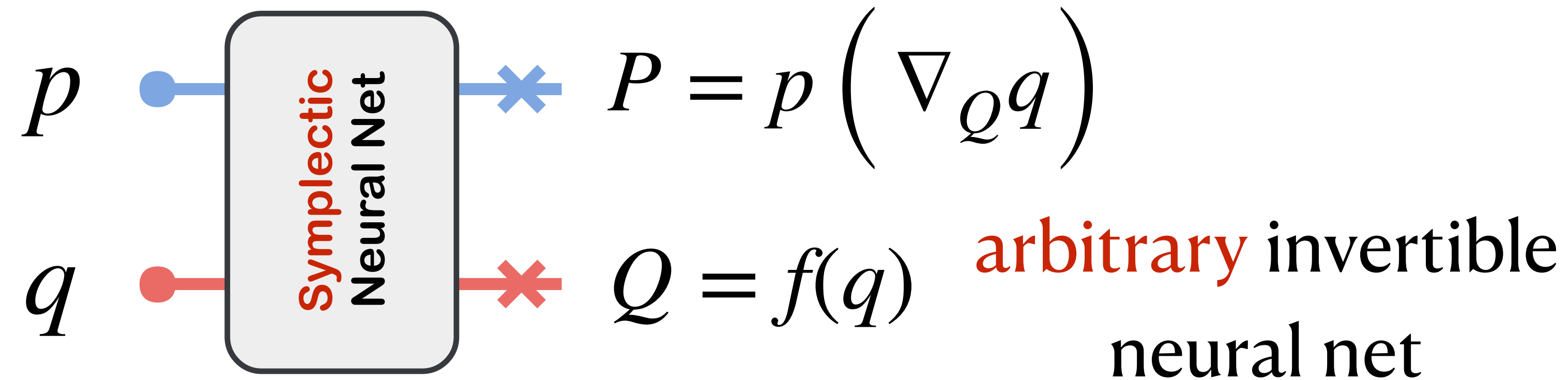


Learn harmonic frequencies of the base to identify slow collective modes

See Bondesan et al 1906.04645, Ishikawa et al 2103.00372 for investigations on integrability

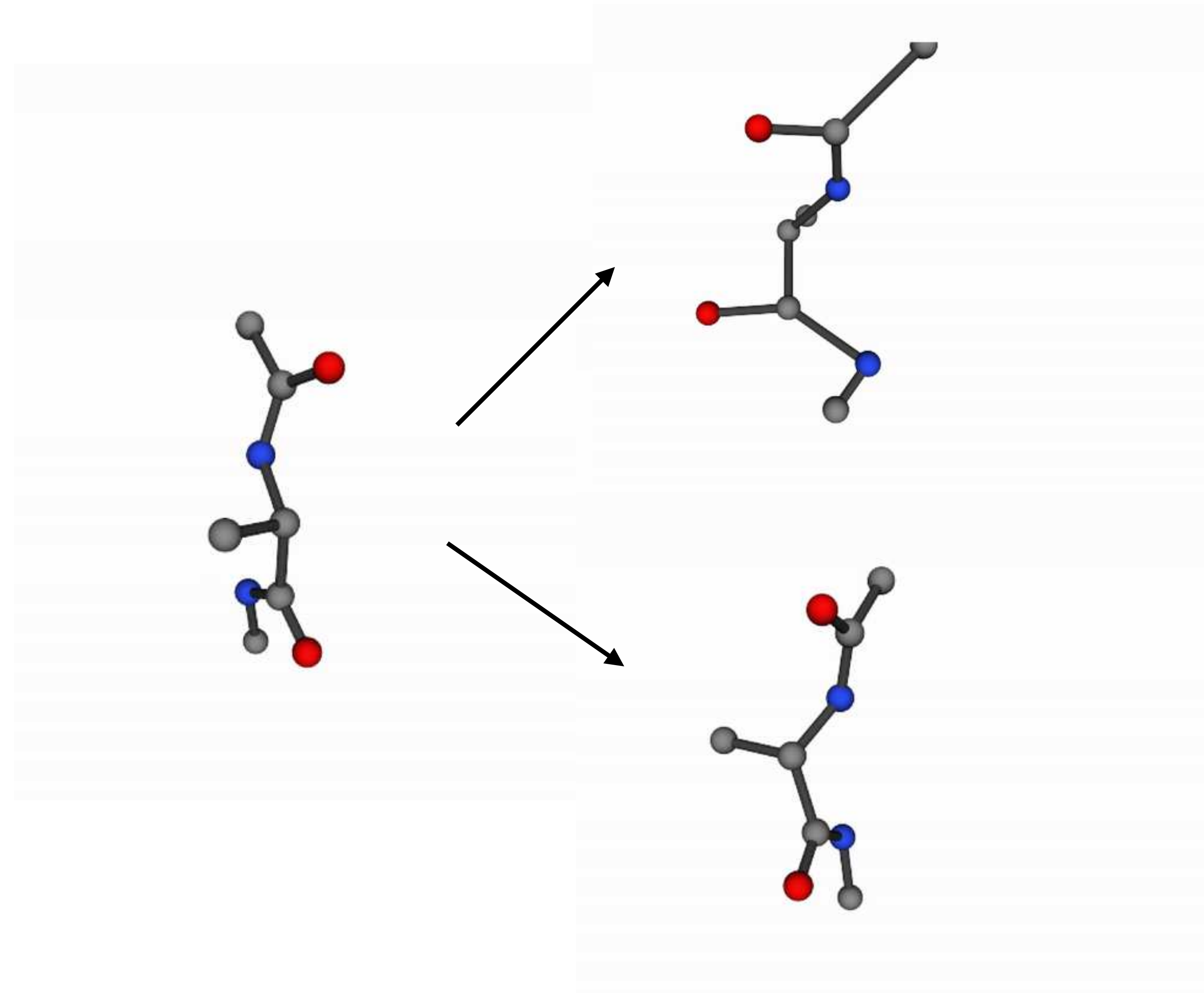
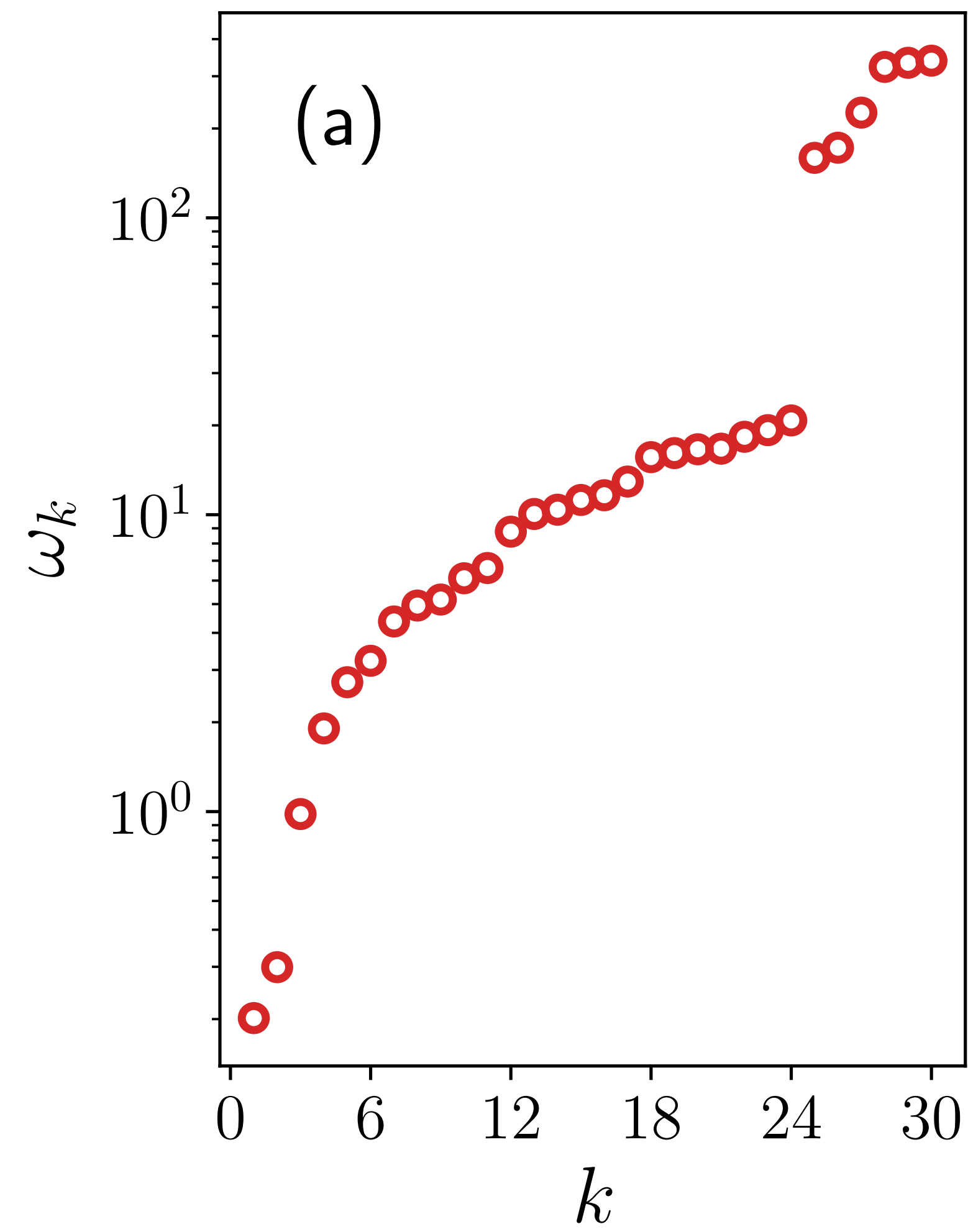
Symplectic primitives

- **Neural point transformations**



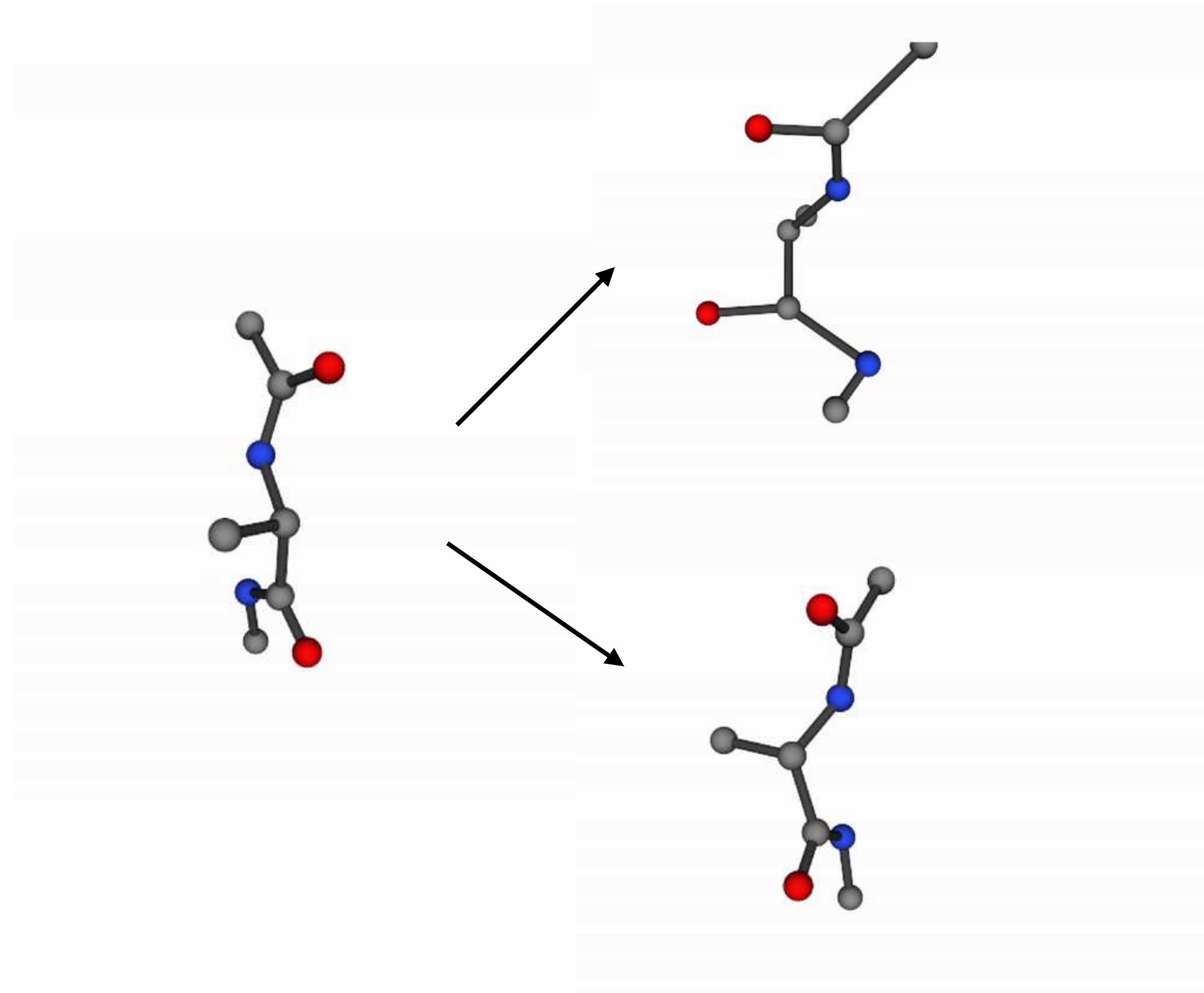
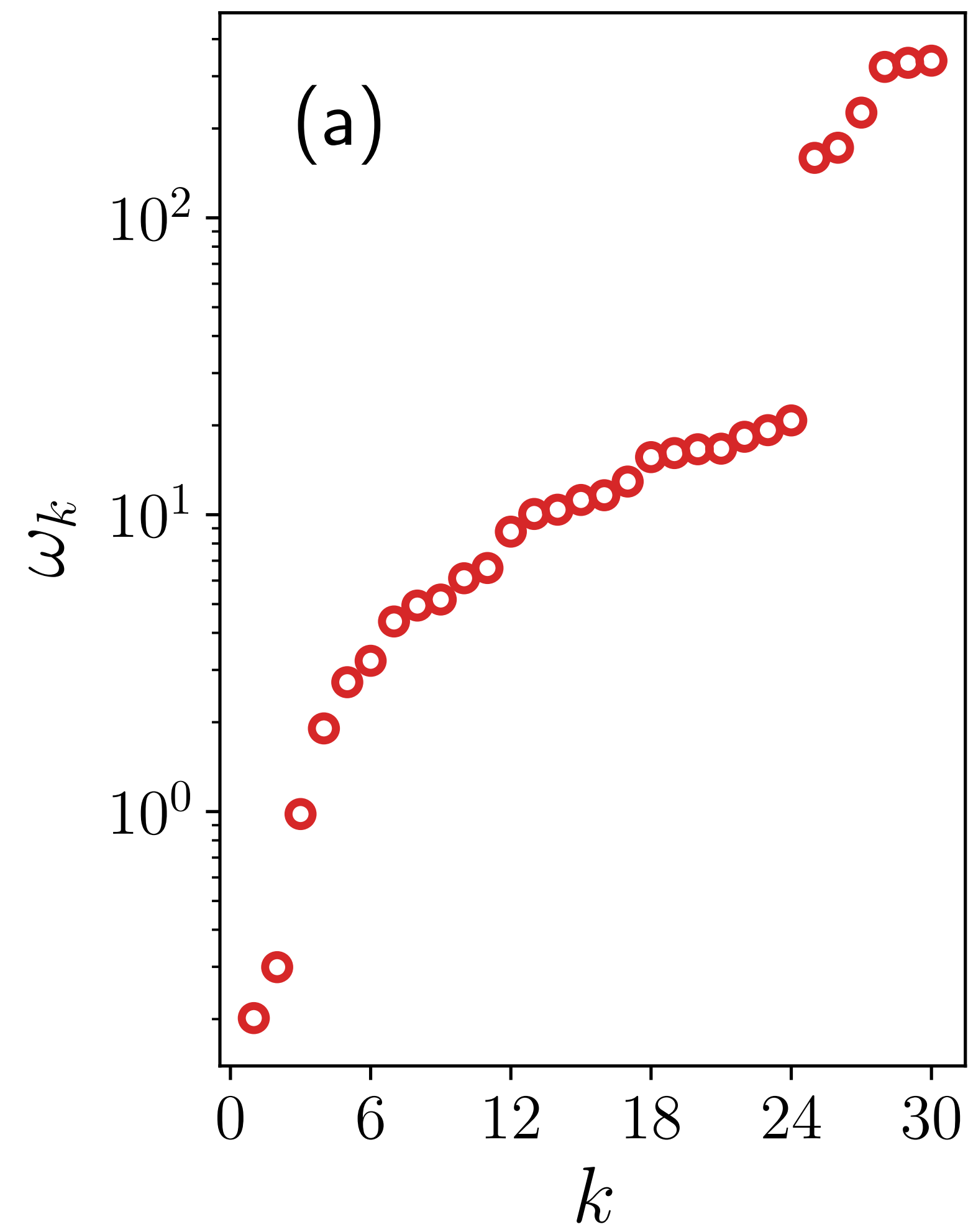
- Linear transformation: Symplectic Lie group $Sp(2n)$
- Continuous-time flow: Symplectic generating functions via Hamiltonian dynamics See also Bondesan, Lamacraft, 1906.04645
Neural ODE, Chen et al, 1806.07366, Monge-Ampère flow, Zhang et al 1809.10188

Li, Dong, Zhang, LW, PRX '20

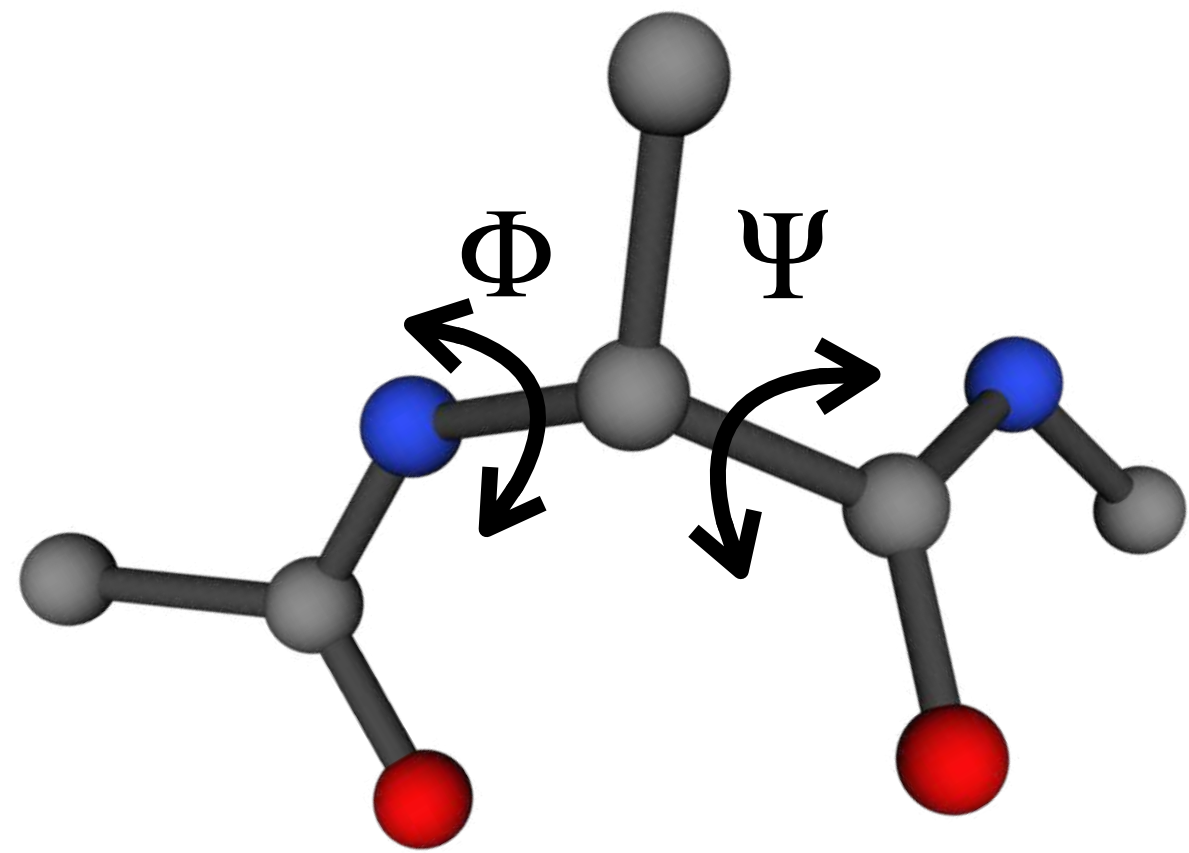


Neural canonical transformation identifies nonlinear slow modes

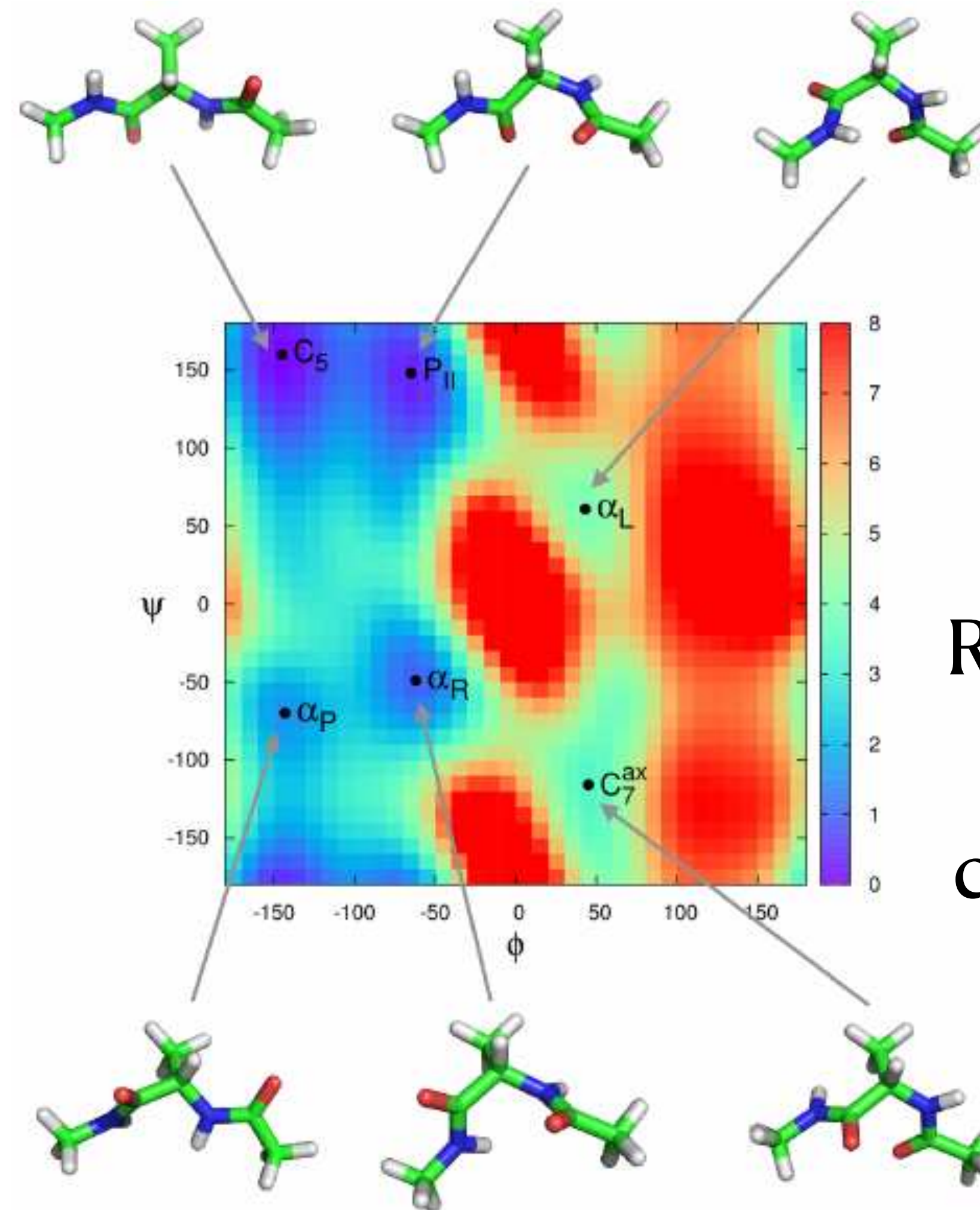
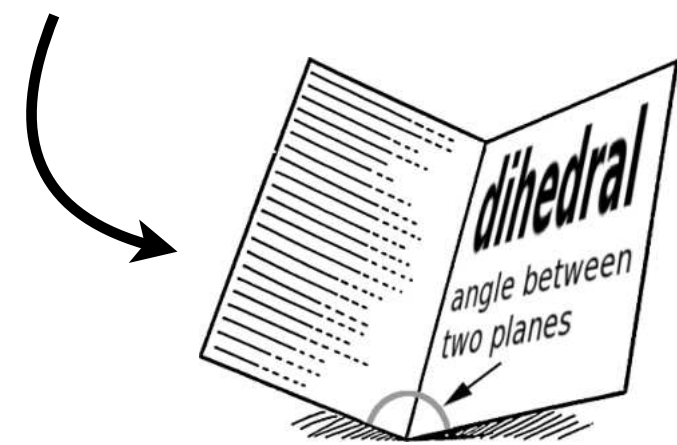
Li, Dong, Zhang, LW, PRX '20



Neural canonical transformation identifies nonlinear slow modes



slow motion of the two torsion angles

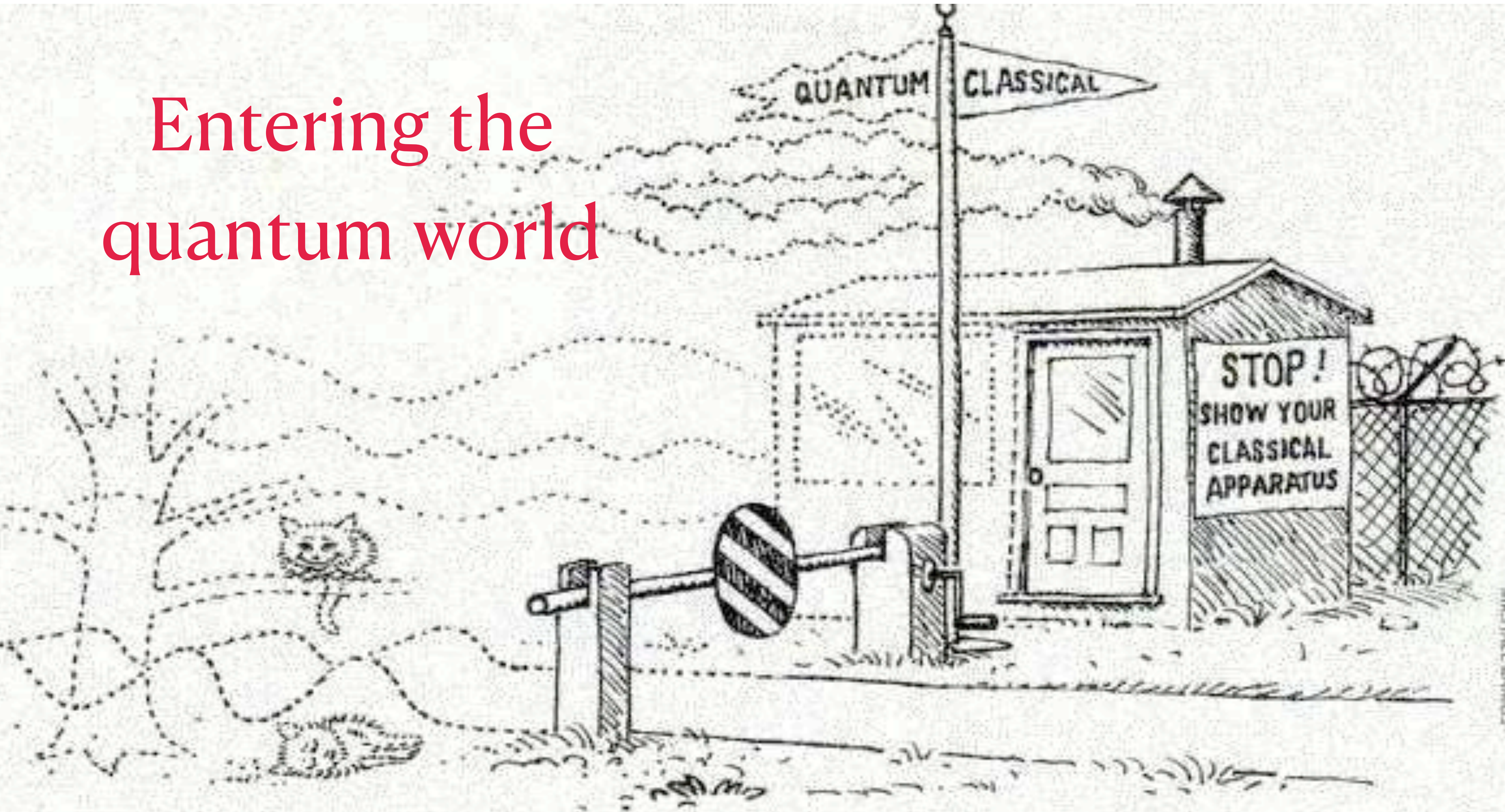


Ramachandran plot of stable conformations

**Dimensional reduction to slow collective variables
useful for control, prediction, enhanced sampling...**

check the paper 1910.00024, PRX '20 for more examples & applications

Entering the quantum world



Neural *canonical* transformations

Li, Dong, Zhang, LW, PRX '20

Xie, Zhang, LW, JML '21

classical world

Symplectic transformation

Probability density

p

Kullback-Leibler divergence

$\mathbb{KL}(p || q)$

quantum world

Unitary transformation

Density matrix

ρ

Quantum relative entropy

$S(\rho || \sigma)$

Neural *canonical* transformations

Li, Dong, Zhang, LW, PRX '20

Xie, Zhang, LW, JML '21

classical world

Symplectic transformation

Probability density

p

Kullback-Leibler divergence

$$\mathbb{KL}(p || q)$$

quantum world

Unitary transformation

Density matrix

ρ

Quantum relative entropy

$$S\left(\rho || \frac{e^{-\beta H}}{Z}\right)$$

The variational free energy principle

Gibbs–Bogolyubov–Feynman–Delbrück–Molière

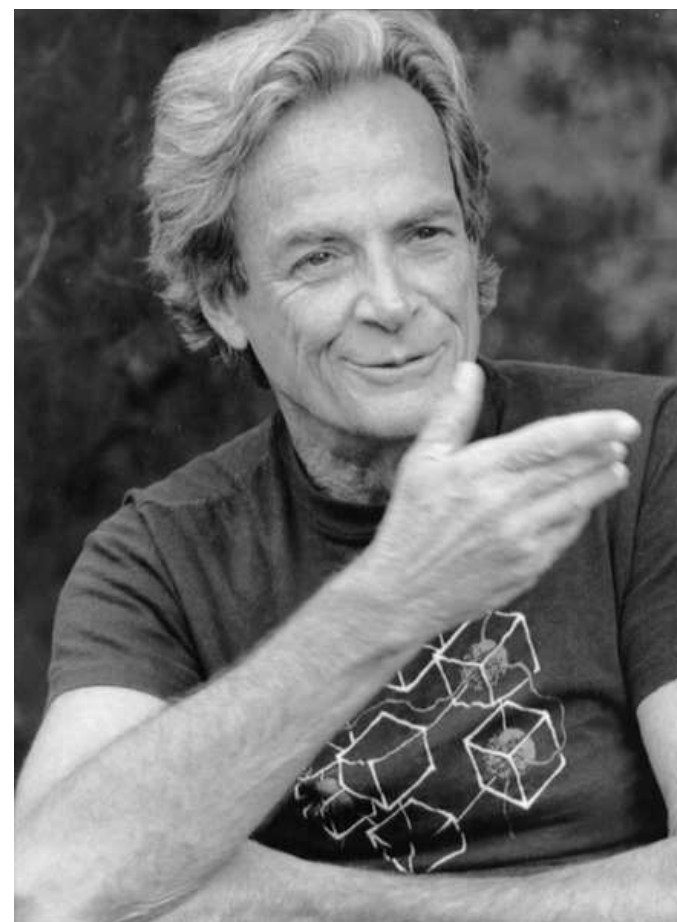
$$\min F[\rho] = \text{Tr}(H\rho) + k_B T \text{Tr}(\rho \ln \rho) \geq F$$



variational density matrix

energy

entropy



Difficulties in Applying the Variational Principle to Quantum Field Theories¹

Richard P. Feynman

ρ ?

Generative models !

¹transcript of Professor Feynman's talk in 1987

1981

14318 citations

Simulating Physics with Computers

Richard P. Feynman

“... it’s a wonderful problem,
because it doesn't look so easy.”

1987

22 citations

**Difficulties in Applying the Variational
Principle to Quantum Field Theories¹**

Richard P. Feynman

“...it is no damn good at all!”

Learn 1 Sensitivity to High Frequencies.

Express 2 Only Gaussian Trial States

Sample 3 We Still Have To Do a Functional Integral

Variational density matrices as generative models

Learnable unitary transformation
generated by point transformation

Learnable probabilistic model
for occupation probability

$\sqrt{\text{flow}}$

JML '22, SciPost Physics'23

See Cranmer et al 1904.05903

Saleh et al, 2308.16468

Siciliano et al 2407.03802

$$\rho = \sum_n U |n\rangle p_n \langle n| U^\dagger$$

VAN
PRL '19

$$\text{Tr}\rho = 1$$

$$\rho \succ 0$$

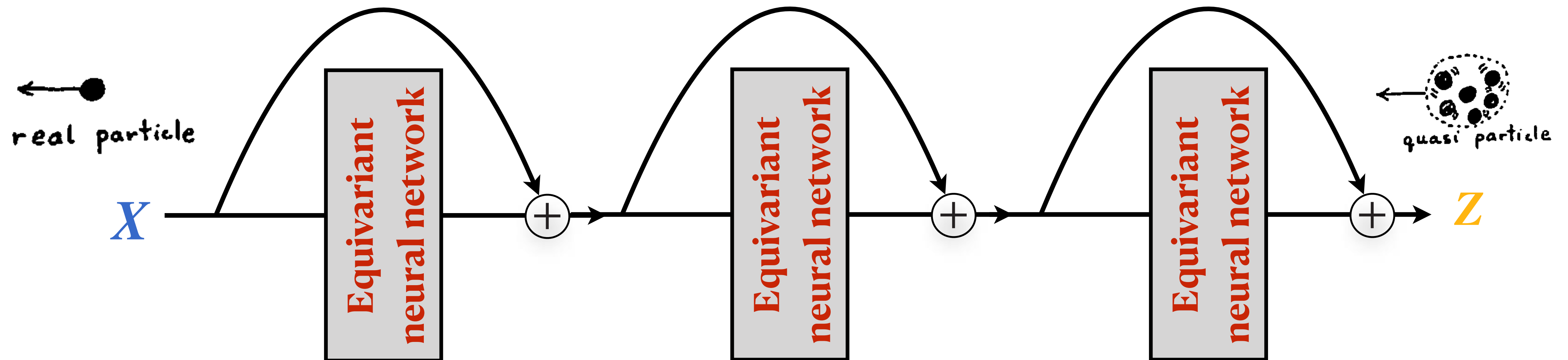
$$\rho^\dagger = \rho$$

Many-body “base” states e.g.
Fermi sea, Hartree-Fock states,
harmonic crystal, ...

The physics of $\sqrt{\text{flow}}$

Xie, Zhang, LW, JML '22

$$\langle \mathbf{X} | \mathbf{U} | \mathbf{n} \rangle = \langle \mathbf{Z} | \mathbf{n} \rangle \cdot \left| \det \left(\frac{\partial \mathbf{Z}}{\partial \mathbf{X}} \right) \right|^{\frac{1}{2}}$$



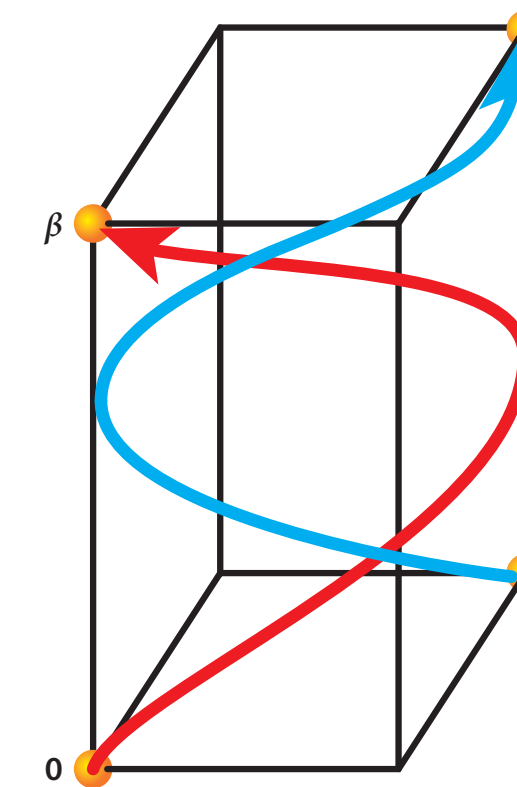
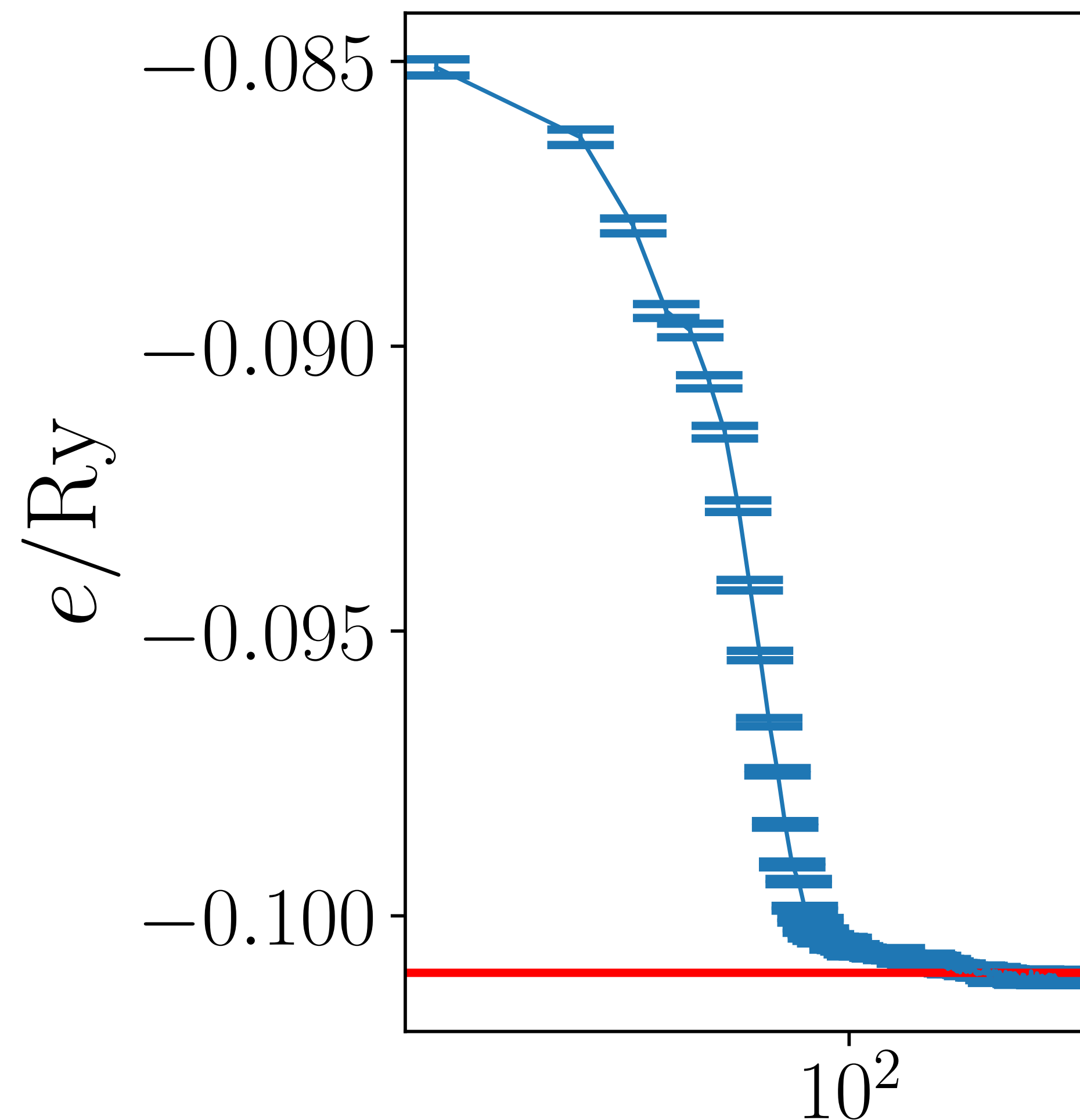
$\mathbf{X} \leftrightarrow \mathbf{Z}$: unitary backflow between particle and quasiparticle coordinates

e.g. backflow transformation Feynman & Cohen 1956 $\mathbf{z}_i = \mathbf{x}_i + \sum_{j \neq i} \eta(|\mathbf{x}_i - \mathbf{x}_j|) (\mathbf{x}_j - \mathbf{x}_i)$

Benchmarks on finite-temperature Coulomb gas

Xie, Zhang, LW, SciPost Physics '23

$r_s = 10, T/T_F = 0.0625, N = 33$



Severe sign problem

r_s	Θ	$\langle sign \rangle$	E_{tot}^{exact}	E_{tot}
4.0	0.0625	-0.00055(62)	-0.5(1)	-0.1023(7)
10.0	0.0625	-0.002(1)	-0.16(2)	-0.1010(1)

Brown et al, PRL '13 Restricted PIMC
see also Schoof et al PRL '15, Malone et al PRL '16

Point Transformations and the Many Body Problem*

M. EGER† AND E. P. GROSS

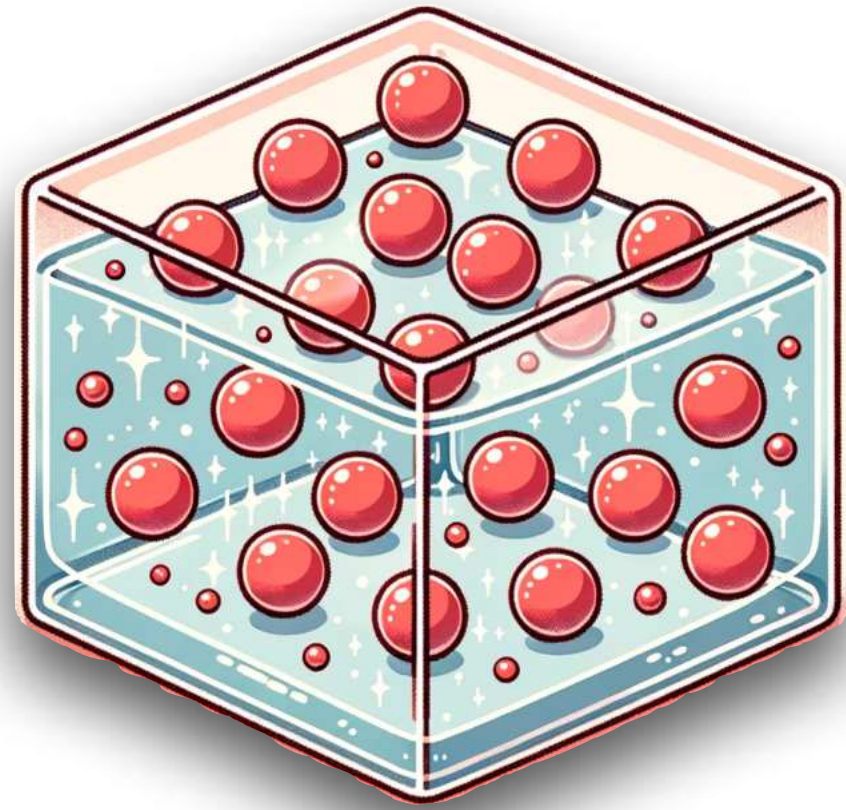
Brandeis University, Waltham, Massachusetts

An investigation is made of possible uses of many dimensional coordinate transformations in the quantum many-body problem. The transformed Hamiltonian is quadratic in the momenta with a space dependent metric. The original potential energy undergoes alteration and an additional “metric” potential energy appears. A relatively complete analysis of the transformed original potential is made, and the coordinate transformation can be used to suppress undesirable features of the original potential. For bosons one can attempt to directly map a complete set of noninteracting states onto approximate eigenstates of the system with interactions. Contact is made with a theory of weakly interacting bosons. In the general case it emerges that a given transformation uniquely fixes all the spatial correlation functions, which can be explicitly computed. The extended point transform can then be used as a link between diverse experimental quantities. The full use of the transformation to compute from first principles requires adequate approximations to the Jacobian and the inverse transform. These problems are not studied.

√flow

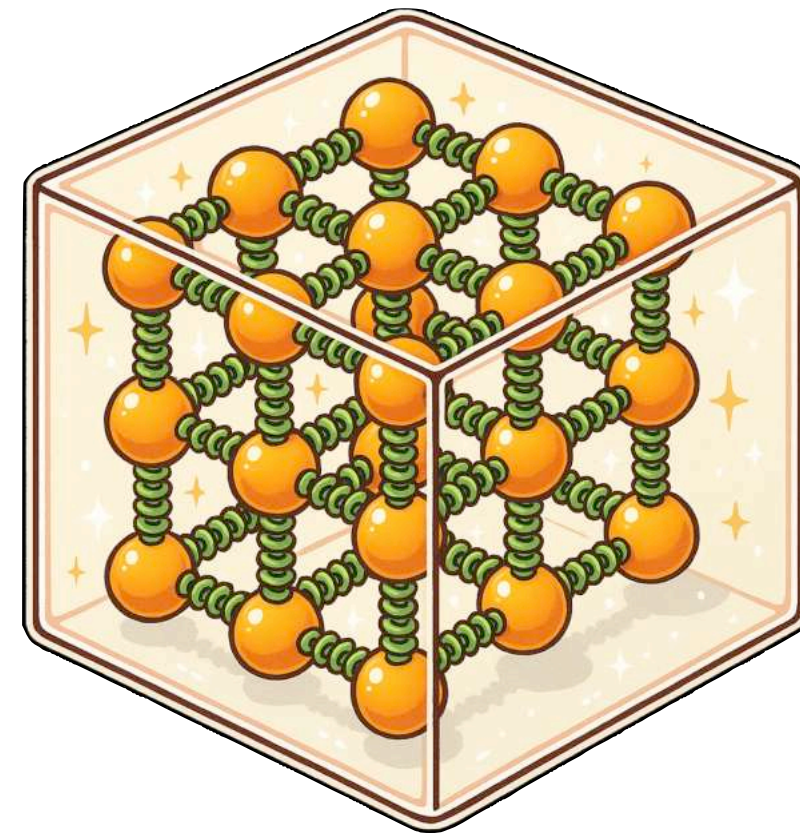
materializes this dream

The deep variational free energy approach



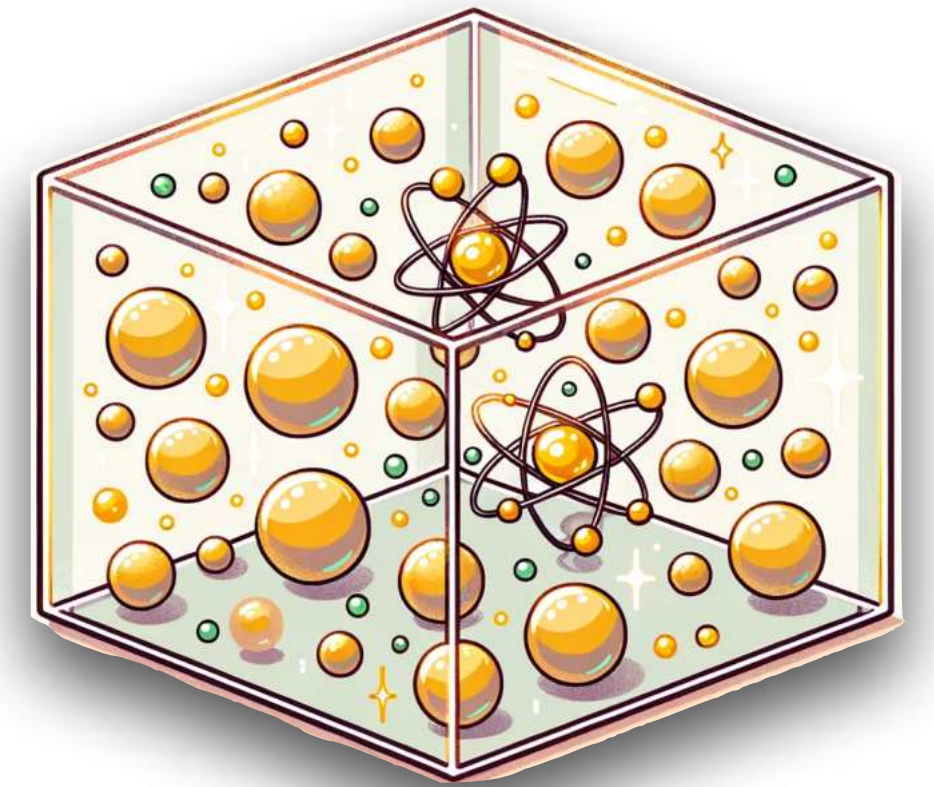
Electron liquid
thermodynamics

JML '22 and SciPost Physics '23



Anaharmonic quantum
solid structure

JCP '24 and PRL '25



Dense hydrogen
equation of state
PRL '23 and PRB '25

A computational framework taking in account of electron correlation, thermal effect, and anharmonic lattices for free energy, entropy, and excitation spectra

Generative AI for **It**

①

$$p(X|y) \propto p(X)p(y|X)$$

Matter inverse design
Exploiting intuitions in data

②

$$F[\rho] = E - TS$$

Nature's cost function
Variational free energy is finally practical

Turning physics problems into stochastic optimization

Leverages the deep learning engine

