

## PAPER



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Facet-engineered CeO<sub>2</sub>/graphene composites for  
enhanced NO<sub>2</sub> gas-sensing†

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The sensitivity of semiconductor gas sensors depends on the exposed crystal planes. High-energy crystal planes usually have an open surface structure and more active sites which favor the absorption of gas molecules and thus might improve the gas sensing performance. However, it is commonly difficult for the high-energy crystal planes to be exposed. In this paper, the morphology and exposed facets of CeO<sub>2</sub> nanoparticles in the CeO<sub>2</sub> nanoparticles/graphene composites are tailored simply by changing the volume ratio of ethylene glycol to deionized water (EG/H<sub>2</sub>O) in the solution during fabrication. CeO<sub>2</sub> nanocubes enclosed by the {100} facets are produced with an EG/H<sub>2</sub>O ratio of 1:1, otherwise CeO<sub>2</sub> nanograins enclosed by the {111} facets are obtained, which is closely related to the changes in the energetics on the surface. Furthermore, it is found that the CeO<sub>2</sub>{100}/graphene composites deliver substantially enhanced gas sensing performance for NO<sub>2</sub>, as compared to CeO<sub>2</sub>{111}/graphene composites. First-principles calculations illustrate that the electrons flow from graphene to the CeO<sub>2</sub>{111} facets resulting in electron depletion on graphene. In contrast, the electrons flow from the CeO<sub>2</sub>{100} surface to graphene resulting in electron accumulation on graphene and the band gap of CeO<sub>2</sub>{100}/graphene is nearly zero implying a metallic behavior. The rapid charge exchange between NO<sub>2</sub> and CeO<sub>2</sub>{100}/graphene composites leads to the improved gas sensing properties. The results present us a clear physical picture for the tunable facets and morphology as well as the enhanced performances of nanoparticles/graphene composites.

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## 1. Introduction

Nitrogen dioxide (NO<sub>2</sub>) is one of the common and toxic air pollutants from combustion and automotive emission causing serious diseases such as chronic bronchitis, emphysema, and respiratory irritation.<sup>1–3</sup> Therefore, it is crucial to develop high-performance sensors to detect NO<sub>2</sub> in air quickly and reliably. Many sensing devices including metal oxide semiconductor gas

sensors, solid electrolyte gas sensors, optic fiber gas sensors, and surface acoustic wave (SAW) gas sensors have been developed<sup>4–7</sup> and among them, metal oxide semiconductors have attracted much attention. However, they normally operate at a high temperature resulting in high power consumption<sup>8</sup> and so low-cost, high-sensitivity, and fast-response sensing materials are much more desirable especially in portable devices.

Essentially, the gas sensing performance of metal oxides is mainly restricted by the lower electrical conductivity at room temperature.<sup>9</sup> To solve this problem, the composites of metal oxides and another material with high electrical conductivity could be developed. As well known, graphene possesses a large specific surface area, high electrical conductance,<sup>10</sup> low signal-to-noise level, and stable response to NO<sub>2</sub> at room temperature.<sup>11–14</sup> If metal oxide nanoparticles can be uniformly dispersed on graphene sheets, the electrical conductance and the charge exchange between metal oxide and NO<sub>2</sub> can be significantly improved, and the structure is expected to have excellent gas sensing characteristics.<sup>9</sup> The composites combining graphene with SnO<sub>2</sub>, ZnO, WO<sub>3</sub>, Co<sub>3</sub>O<sub>4</sub>, and Cu<sub>2</sub>O have been shown to deliver good sensing performance for NO<sub>2</sub> at room temperature.<sup>15–19</sup>

Compared to other metal oxides, CeO<sub>2</sub> has more active gas sensing properties due to its polymorphism and mixed valence,

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and it becomes one of the potential candidate materials in the field of gas sensing.<sup>20</sup> Fu *et al.* successfully synthesized the composites of CeO<sub>2</sub> nanoparticles and graphene, and illustrated their fast response to 100 ppm NO<sub>2</sub>, but the sensitivity is only 10.39% at room temperature.<sup>20</sup> It is mandatory to improve the sensitivity through modification of CeO<sub>2</sub> nanoparticles. Since gas sensing is a surface process, the sensitivity depends on the exposed crystal planes.<sup>21–23</sup> Usually, high-energy crystal planes have an open surface structure and more active sites<sup>24–29</sup> which favor the absorption of gas molecules. Furthermore, the electronic states are also facet dependent and hence, the energy band alignment of the heterojunction and charge transfer across the interface in composites may be adjusted by changing the facets of metal oxides.<sup>30</sup> Accordingly, the gas sensing performance of CeO<sub>2</sub>/graphene composites might be enhanced by controllable faceting of CeO<sub>2</sub> nanoparticles. As for CeO<sub>2</sub>, the {100} plane is an open surface and should exhibit more active properties than {111} and {110} planes.<sup>31–36</sup> However, the surface energy of the CeO<sub>2</sub>{100} plane is higher and the atomic configuration is less stable than the {111} and {110} facets.<sup>37</sup> Consequently, CeO<sub>2</sub>{100} planes diminish rapidly during grain growth<sup>38</sup> and synthesis of CeO<sub>2</sub> with controllable facets is still challenging.

In this work, the hydrothermal method is adopted to fabricate the composites of CeO<sub>2</sub> nanoparticles and graphene, and it is found that the morphology and facets of the CeO<sub>2</sub> nanoparticles in the composites can be tailored simply by changing the volume ratio of ethylene glycol (EG) to deionized water (H<sub>2</sub>O) in the solution. Compared to CeO<sub>2</sub>{111} facets/graphene composites, CeO<sub>2</sub>{100} facets/graphene composites exhibit higher sensitivity and more rapid response to NO<sub>2</sub>. In particular, the sensitivity to 100 ppm NO<sub>2</sub> is 33% at room temperature, and it is enhanced by 200% as compared to the reported results (10.39%).<sup>20</sup> It is demonstrated that the gas sensing performance of CeO<sub>2</sub>/graphene composites could be enhanced by controllable faceting of CeO<sub>2</sub> nanoparticles. First-principles calculations are performed to elucidate the mechanisms for the controllable faceting as well as the enhanced performance.

## 2. Experimental and theoretical section

### 2.1 Material preparation

Graphene oxide (GO) was synthesized from pure graphite powders by the modified Hummers method<sup>39</sup> and graphene was prepared by reduction of GO with H<sub>2</sub> at 400 °C for 1 h. 10 mg of graphene were dissolved in a solvent consisting of ethylene glycol (EG) and deionized water (H<sub>2</sub>O) with different volume ratios (EG + H<sub>2</sub>O = 0 mL + 40 mL, 10 mL + 30 mL, 20 mL + 20 mL, 30 mL + 10 mL, or 40 mL + 0 mL) under sonication for 2 h. 0.5 mmol Ce(NO<sub>3</sub>)<sub>3</sub>·6H<sub>2</sub>O, 1.5 mmol NaOH, and 0.3 g of PVP were added to the solvent and stirred magnetically for 3 h. The mixture was transferred to a 100 mL Teflon-lined autoclave and heated to 180 °C for 24 h. Afterwards, the precipitate was washed three times with distilled water and ethanol and dried at 70 °C in a vacuum oven

overnight to obtain the CeO<sub>2</sub>/graphene nanocomposite. Pure CeO<sub>2</sub> nanoparticles were prepared in the same manner at different EG/H<sub>2</sub>O volume ratios (0:40, 10:30, 20:20, 30:10, and 40:0), respectively.

### 2.2 Characterization

The X-ray diffraction (XRD) patterns were acquired between 10° and 90° (SHIMADZU XRD-7000S) with Cu K<sub>α</sub> radiation. Raman spectra were recorded on a Horiba HR800 spectrometer with a 633 nm laser as the excitation source. The morphology of the samples was characterized using scanning electron microscopy (SEM, FEI Quanta 600S) and high-resolution transmission electron microscopy (HR-TEM, JEOL JEM 2100F). The Mott–Schottky curves were measured on an electrochemical analyzer (CHI760E, Shanghai). The voltage in the Mott–Schottky measurement was scanned from −0.4 V to 0.4 V *versus* SCE (saturated calomel reference electrode) at a frequency of 10 000 Hz and 0.5 M NaSO<sub>3</sub> was used as the electrolyte. A Keithley 4200 Semiconductor Characterization System was used to obtain the current–voltage (*I*–*V*) curves.

### 2.3 Gas sensing measurements

5 mg of CeO<sub>2</sub>/graphene powder were dispersed in 1 mL of ethyl alcohol under sonication for 30 min and then 0.2 mL of the suspension was added dropwise onto SiO<sub>2</sub> with interdigitated gold electrode arrays with a pipette. The substrate was heated to 60 °C for 15 min to evaporate the solvent before the gas sensing measurement. The sensor was installed in a test chamber with an inlet and an outlet. The electrical resistance was measured at room temperature at a relative humidity (RH) of 26%. The gas concentration was controlled by injection and the chamber was purged with air to recover the resistance.

### 2.4 Theoretical model and calculations

Density-functional theory (DFT) calculations were conducted using the Vienna ab initio simulation package (VASP).<sup>40–42</sup> The kinetic energy cutoff for the plane wave basis set was 500 eV. Brillouin zone integration on grids with 3 × 3 × 1 and 11 × 11 × 1 Monkhorst–Pack *k*-points was implemented for geometrical optimization and calculation of the density of states, respectively. Geometrical optimization was performed by using the conjugate gradient scheme until the Hellman–Feynman force was smaller than 0.05 eV Å<sup>−1</sup>. The vacuum space perpendicular to the CeO<sub>2</sub>/graphene is around 15 Å, which was enough to separate the interaction between periodic images. To model the interface, a 2 × 2 CeO<sub>2</sub>(111) slab (O–Ce–O, nine layers) and a 3 × 3 graphene sheet were stacked with a lattice mismatch of 3.7%, while a 2 × 2 CeO<sub>2</sub>(100) slab (eight layers), and a 3 × 2 graphene sheet were stacked with a lattice mismatch of 4.6%. A 3 × 3 slab was employed to calculate the surface energy and adsorption energy. The slab thickness was 10 Å and the vacuum space between the slabs was 20 Å. The surface energy was calculated using the following equation:

$$E = \frac{E_{\text{surf}} - nE_{\text{bulk}}}{A} \quad (1)$$

in which  $E_{\text{surf}}$  is the total energy of the slab,  $E_{\text{bulk}}$  is the average energy of each  $\text{CeO}_2$  unit in the bulk, and  $n$  is the number of  $\text{CeO}_2$  units in each slab. The surface area of the slab is denoted by  $A$ .

The absorption energy of ethylene glycol on the  $\text{CeO}_2$  nanoparticles was calculated as follows:

$$E_{\text{ads}} = E_{(\text{EG}+\text{CeO}_2)} - E_{\text{EG}} - E_{\text{CeO}_2} \quad (2)$$

where  $E_{\text{EG}+\text{CeO}_2}$  is the total energy of the relaxed ethylene glycol- $\text{CeO}_2$  system,  $E_{\text{EG}}$  is the energy of free ethylene glycol, and  $E_{\text{CeO}_2}$  is the total energy of the relaxed  $\text{CeO}_2$  slab.

The formation energy of the  $\text{CeO}_2$ /graphene hybrid structure was calculated as:

$$E = E_{\text{CeO}_2/\text{graphene}} - E_{\text{CeO}_2} - E_{\text{graphene}} \quad (3)$$

in which  $E_{\text{CeO}_2/\text{graphene}}$  is the total energy of the relaxed  $\text{CeO}_2$ /graphene system,  $E_{\text{CeO}_2}$  is the total energy of the relaxed  $\text{CeO}_2$  slab, and  $E_{\text{graphene}}$  is the total energy of the relaxed graphene slab.

### 3. Results and discussion

#### 3.1 Characterization of $\text{CeO}_2$ /graphene composites

The synthesis process is schematically illustrated in Fig. 1 and the detailed experimental procedures are described in the Experimental section. The  $\text{CeO}_2$ /graphene nanocomposites are characterized by XRD, Raman spectra, SEM, TEM, and HR-TEM. As shown in Fig. 2, the main XRD peaks can be assigned to the (111), (200), (220), (311), (222), (400), (331), (420) and (422) planes of  $\text{CeO}_2$  in a cubic fluorite-type structure ( $a = 5.41 \text{ \AA}$ ,  $Fm\bar{3}m$  space group, JCPDS card no. 34-394). The XRD peaks from  $\text{CeO}_2$  decrease gradually from samples (a) to (e) and the full width at half-maximum (FWHM) increases gradually due to suppressed growth of the  $\text{CeO}_2$  nanoparticles with the introduction of ethylene glycol. According to Scherrer's equation, the average size of the  $\text{CeO}_2$  nanoparticles is about

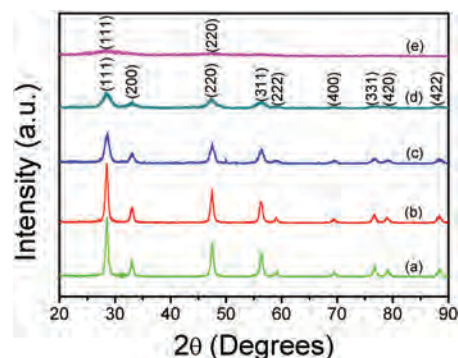


Fig. 2 XRD patterns of the  $\text{CeO}_2$ /graphene composites prepared with different EG/ $\text{H}_2\text{O}$  volume ratios: (a) 0 : 40, (b) 10 : 30, (c) 20 : 20, (d) 30 : 10, and (e) 40 : 0.

20 nm, 15 nm, 10 nm, and 5 nm for samples (a) to (d), respectively. Fig. 3 shows the Raman spectra of graphene,  $\text{CeO}_2$ , and  $\text{CeO}_2$ /graphene composites. Two prominent peaks corresponding to the G and D bands of graphene are observed. The G band is usually assigned to the  $E_{2g}$  vibration mode and the D band is associated with structural defects with the  $A_{1g}$  symmetry.<sup>43,44</sup> The intensity ratios of D to G bands ( $I_{\text{D}}/I_{\text{G}}$ ) are 1.03, 1.13, 1.06, 1.05, and 1.04 for samples (a) to (e), respectively, which are larger than that of graphene (0.97) due to partial modification of the functional groups and formation of new  $\text{sp}^2$  clusters during reduction.<sup>45</sup> The new peak at  $463 \text{ cm}^{-1}$  from  $\text{CeO}_2$ /graphene composites corresponds to symmetrical stretching of the Ce-O bond, and is sensitive to the structural disorder in the oxygen sub-lattice.<sup>46</sup> The peak red-shifts from  $465 \text{ cm}^{-1}$  in  $\text{CeO}_2$  to  $463 \text{ cm}^{-1}$  in  $\text{CeO}_2$ /graphene composites as a result of charge transfer between them.

The morphology of the  $\text{CeO}_2$ /graphene nanocomposites is characterized by SEM, TEM and HR-TEM, and the results are shown in Fig. S1 (ESI<sup>†</sup>) and Fig. 4. As shown in Fig. S1 (ESI<sup>†</sup>),  $\text{CeO}_2$  nanoparticles are dispersedly attached on graphene

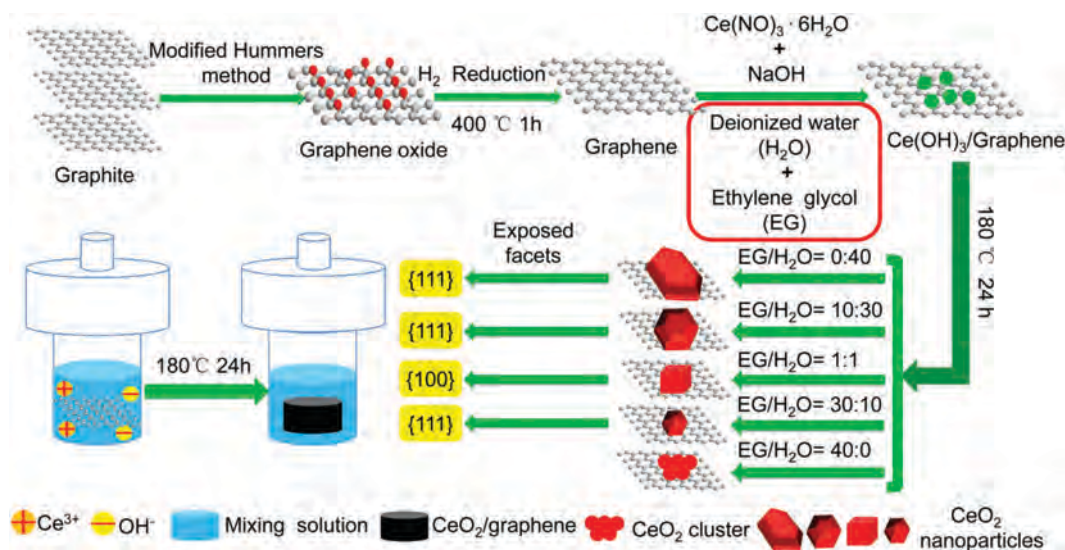


Fig. 1 Schematic illustration of the synthesis of  $\text{CeO}_2$ /graphene composites.

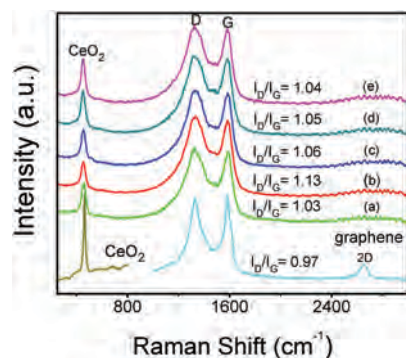


Fig. 3 Raman spectra of the CeO<sub>2</sub>, graphene and CeO<sub>2</sub>/graphene nanocomposites prepared with different EG/H<sub>2</sub>O volume ratios: (a) 0 : 40, (b) 10 : 30, (c) 20 : 20, (d) 30 : 10, and (e) 40 : 0.

sheets, and almost no isolated CeO<sub>2</sub> nanoparticles are evidenced. Since the graphene sheets were synthesized by reduction of graphene oxide, there are commonly a great number of oxygen groups, vacancies and other structural defects.<sup>9</sup> As a result, the nucleation and growth of CeO<sub>2</sub> nanoparticles are preferred on graphene sheets because of the strong interaction between them. This leads to good adhesion between CeO<sub>2</sub> nanoparticles and graphene sheets. Fig. 4 shows the SEM, TEM and HR-TEM images of the CeO<sub>2</sub>/graphene nanocomposites prepared with different EG/H<sub>2</sub>O volume ratios. If deionized water (40 mL H<sub>2</sub>O) is used, the CeO<sub>2</sub> nanocrystals have an irregular octahedral shape with an average diameter of 20 nm (Fig. 4a<sub>1</sub> and a<sub>2</sub>).

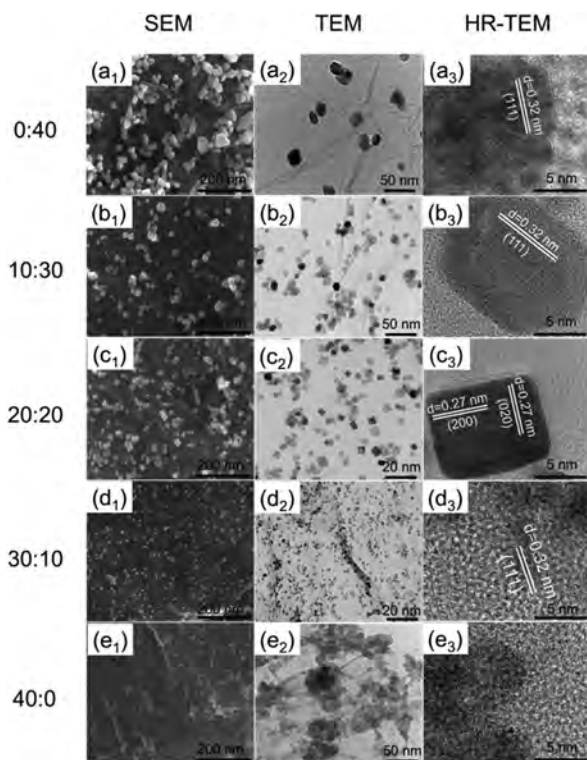


Fig. 4 SEM, TEM and HR-TEM images of the CeO<sub>2</sub>/graphene nanocomposites prepared with different EG/H<sub>2</sub>O volume ratios: (a) 0 : 40, (b) 10 : 30, (c) 20 : 20, (d) 30 : 10, and (e) 40 : 0.

Clear lattice fringes with a separation of 0.32 nm observed from the side surface correspond to the {111} facets of CeO<sub>2</sub> (Fig. 4a<sub>3</sub>). When the volume ratio of EG/H<sub>2</sub>O is increased to 10 : 30, the CeO<sub>2</sub> nanoparticles evolve into an irregular truncated octahedral morphology with an average size of 15 nm (Fig. 4b<sub>1</sub> and b<sub>2</sub>) and they are also mainly enclosed by the {111} facets (Fig. 4b<sub>3</sub>). When the volume ratio of EG/H<sub>2</sub>O is further adjusted to 20 : 20, CeO<sub>2</sub> nanoparticles become nanocubes with an average size of 10 nm (Fig. 4c<sub>1</sub> and c<sub>2</sub>). As shown in the HR-TEM image in Fig. 4c<sub>3</sub>, the lattice spacing perpendicular to the sides of cubes is 0.27 nm corresponding to the {200} planes of CeO<sub>2</sub>. This means that CeO<sub>2</sub> nanocubes enclosed by the {200} planes are synthesized. If the volume ratio of EG/H<sub>2</sub>O is increased to 30 : 10, the CeO<sub>2</sub> becomes irregular truncated octahedrons again but with a reduced size of 5 nm (Fig. 4d<sub>1</sub> and d<sub>2</sub>) and the CeO<sub>2</sub> nanocrystals are enclosed by the {111} facets (Fig. 4d<sub>3</sub>). If the EG/H<sub>2</sub>O volume ratio is increased to 40 : 0, that is, only EG in the solution, CeO<sub>2</sub> nanoclusters in lieu of nanoparticles are obtained (Fig. 4e<sub>1</sub>–e<sub>3</sub>). If the volume ratio of EG/H<sub>2</sub>O is changed to 10 : 10, 15 : 15, or 25 : 25 (that is, the same volume of EG and H<sub>2</sub>O), the cubic morphology of the CeO<sub>2</sub> particles enclosed by {100} facets is maintained (Fig. 5). The EDS maps in Fig. S2 (ESI<sup>†</sup>) show that C, O, and Ce are homogeneously distributed throughout the samples confirming the formation of the CeO<sub>2</sub>/graphene composite.

### 3.2 The role of EG in controlling the facets and morphology of CeO<sub>2</sub> nanoparticles

In order to understand the role of ethylene glycol in the formation of the CeO<sub>2</sub>{111} and {100} facets, density functional theory (DFT) calculations are performed to determine the surface energy of the {111} and {100} facets as well as the absorption energy of ethylene glycol. Fig. S3 (ESI<sup>†</sup>) presents the results. The surface energy of the {100} facets (5.1 J m<sup>-2</sup>) is higher than that of the {111} facets (0.70 J m<sup>-2</sup>) by 600%. Growth on the {100} facets is more rapid and so the {100} facets disappear gradually resulting in the formation of irregular octahedral CeO<sub>2</sub> nanoparticles

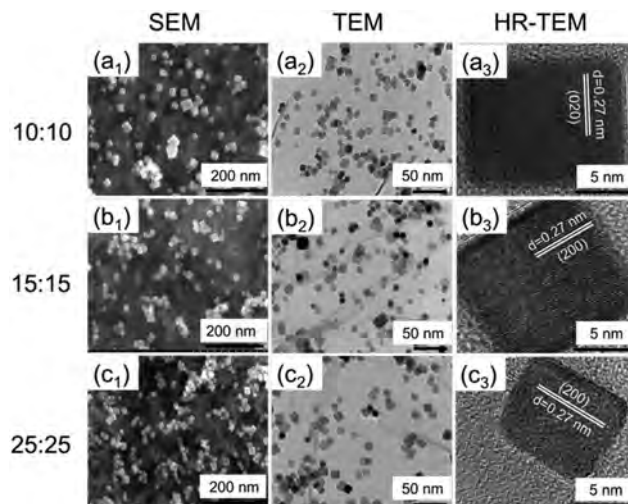


Fig. 5 SEM, TEM and HR-TEM images of the CeO<sub>2</sub>/graphene nanocomposites prepared with different EG/H<sub>2</sub>O volume ratios: (a) 10 : 10, (b) 15 : 15, and (c) 25 : 25.

with the {111} facets being exposed (Fig. 4a<sub>1</sub>–a<sub>3</sub>). Furthermore, the absolute value of adsorption energy for each ethylene glycol molecule on the {100} facets (0.7 eV) is larger than that on the {111} facets (0.047 eV). Hence, the ethylene glycol molecules preferentially absorb on the {100} facets and the surface energy of the {100} facets decreases from 5.1 J m<sup>-2</sup> to -1.43 J m<sup>-2</sup>, which is even smaller than that of the clean {111} facets (0.7 J m<sup>-2</sup>). Consequently, the growth on the {111} facets becomes more rapid, and CeO<sub>2</sub> nanoparticles surrounded by the {100} facets are fabricated by using the appropriate volume ratio of EG/H<sub>2</sub>O (for instance, 20 : 20) in the hydrothermal process. This promotes the formation of nanocubes with {100} facets exposed (Fig. 4c<sub>1</sub>–c<sub>3</sub>). However, adsorption of EG on the {111} facets becomes strong at large concentrations of ethylene glycol. Growth on all facets is then suppressed and the surface energy of the {111} facets (-9.71 J m<sup>-2</sup>) becomes the smallest again giving rise to the formation of small CeO<sub>2</sub> nanoparticles surrounded by the {111} facets (Fig. 4d<sub>1</sub>–d<sub>3</sub>). At a sufficiently large concentration of ethylene glycol, for example, at a volume ratio of EG/H<sub>2</sub>O of 40 : 0, there are enough ethylene glycol molecules absorbed on the facets and growth on all exposed facets is inhibited leading to nano-clusters (Fig. 4e<sub>1</sub>–e<sub>3</sub>). So the concentration of ethylene glycol is crucial to control the exposed facets as well as the morphology and size of the CeO<sub>2</sub> nanoparticles.

### 3.3 The gas sensing performances of CeO<sub>2</sub>/graphene composites

Fig. 6 displays the Mott–Schottky (MS) curves. They are approximately linear with a negative slope for both CeO<sub>2</sub> nanoparticles and CeO<sub>2</sub>/graphene composites, indicative of p-type semiconducting behaviour. The carrier density can be calculated by the Mott–Schottky equation:

$$\frac{1}{C_{SC}^2} = \frac{2}{e\epsilon\epsilon_0 N} \left( E - E_{fb} - \frac{KT}{e} \right) \quad (4)$$

where  $C_{SC}$  is the capacitance of the space charge region,  $e$  is the electronic charge,  $\epsilon$  is the relative dielectric constant of the CeO<sub>2</sub>,

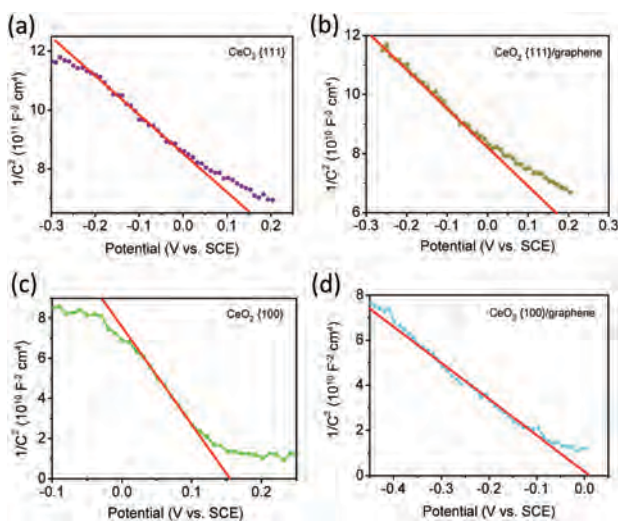


Fig. 6 Mott–Schottky plots of the CeO<sub>2</sub>{111} plane, CeO<sub>2</sub>{100} plane, CeO<sub>2</sub>{111}/graphene and CeO<sub>2</sub>{100}/graphene nanocomposites.

$\epsilon_0$  is the permittivity under vacuum,  $N$  is the hole density,  $E$  is the applied potential,  $E_{fb}$  is the flat band potential,  $K$  is the Boltzmann constant, and  $T$  is the absolute temperature.<sup>47</sup> The densities of holes in the CeO<sub>2</sub> nanoparticles with exposed {111} facets, CeO<sub>2</sub>{111} facets/graphene composites, CeO<sub>2</sub> nanoparticles with exposed {100} facets, and CeO<sub>2</sub>{100} facets/graphene composites were about  $1.18 \times 10^{19}$  cm<sup>-3</sup>,  $2.87 \times 10^{19}$  cm<sup>-3</sup>,  $3.37 \times 10^{19}$  cm<sup>-3</sup> and  $7.67 \times 10^{19}$  cm<sup>-3</sup>, respectively. Hence, the carrier concentration is increased by 100% in the CeO<sub>2</sub>/graphene nanocomposites.

The sensitivity of the samples to NO<sub>2</sub>,  $(R_{gas} - R_{air})/R_{air}$ , is determined, in which  $R_{gas}$  and  $R_{air}$  denote the resistance of the sample exposed to NO<sub>2</sub> gas and air, respectively. First of all, the response curves of the CeO<sub>2</sub>/graphene sensor were measured at different temperatures, and the results are shown in Fig. S4 (ESI<sup>†</sup>). As the working temperature is increased, the response and recovery times are shortened, but the sensitivity is greatly lowered. At room temperature, adsorption of NO<sub>2</sub> continues during the gas exposure, while desorption is negligible owing to the strong bonding between NO<sub>2</sub> and CeO<sub>2</sub>/graphene nanocomposites.<sup>48</sup> Prior to the equilibrium between adsorption and desorption, the sensor resistance decreases continuously with increasing NO<sub>2</sub> coverage.<sup>48</sup> However, the desorption of NO<sub>2</sub> is enhanced at higher working temperature, the balance between adsorption and desorption in both gas exposure and recovering process can be quickly reached.<sup>48</sup> So the response and recovery times are shortened at higher working temperatures, and the sensitivity is lowered. The better room-temperature sensing performance of CeO<sub>2</sub>/graphene nanocomposites is expected for monitoring inflammable and explosive environments.<sup>49</sup> Fig. 7a shows a sensitivity of 50 ppm NO<sub>2</sub> at room temperature. During exposure to NO<sub>2</sub>, the resistance of CeO<sub>2</sub>/graphene nanocomposites decreases sharply. NO<sub>2</sub> is an electron acceptor due to the unpaired electrons on the nitrogen atoms.<sup>48</sup> The electrons will transfer from CeO<sub>2</sub>/graphene composites to NO<sub>2</sub>, thereby increasing the concentration of holes in CeO<sub>2</sub>/graphene composites and reducing the electrical resistance. During subsequent exposure to air, NO<sub>2</sub> molecules desorb from the surface and the resistance goes up.<sup>48</sup> The sensitivity values of the samples a–e are 9.88%, 10.57%, 24.82%, 17.78%, and 7.76%, respectively, being 1.79, 1.92, 4.51, 3.23, and 1.41 times that of pure CeO<sub>2</sub> nanoparticles (5.50%). In particular, sample c with CeO<sub>2</sub> nano-cubes surrounded by the {100} facets delivers the best sensing performance. The sensitivity is enhanced by 350% and the response and recovery times by which 90% of the resistance is restored in the adsorption and desorption processes are about 30 s and 85 s, respectively (Fig. 7b). As shown in Fig. 7c, the sensitivity of pure CeO<sub>2</sub> and CeO<sub>2</sub>/graphene composites decreases with NO<sub>2</sub> concentrations in the range from 200 ppm to 10 ppm. Sample c with CeO<sub>2</sub> nano-cubes surrounded by the {100} facets always shows the highest sensitivity and 10 ppm NO<sub>2</sub> gas can be detected with a large resistance variation of 14%. It is enhanced by 1400% and 140% compared to pure CeO<sub>2</sub> nanoparticles and CeO<sub>2</sub>{111} facets/graphene composites (sample d), respectively. Fig. 7d shows the response time to NO<sub>2</sub> in the concentration range of 10–200 ppm. The response time decreases with increasing concentrations of NO<sub>2</sub> and in particular, from 75 s

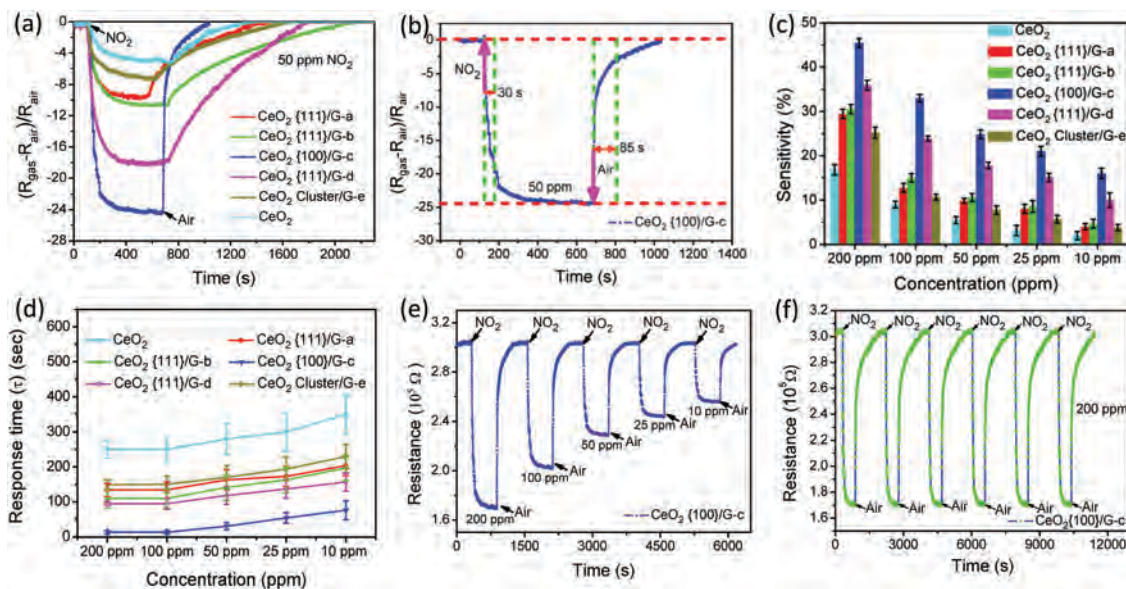


Fig. 7 (a) Resistance variation of  $\text{CeO}_2$ ,  $\text{CeO}_2$ /graphene composites a–e with respect to 50 ppm  $\text{NO}_2$ ; (b) response and recovery curves of  $\text{CeO}_2$ /graphene-c in the presence of 50 ppm  $\text{NO}_2$  gas at room temperature; (c) sensitivity of pure  $\text{CeO}_2$ ,  $\text{CeO}_2$ /graphene composites a–e for 10–200 ppm  $\text{NO}_2$  at room temperature; (d) response time to 10–200 ppm  $\text{NO}_2$ . (e) Response and recovery curves of  $\text{CeO}_2$ (100)/graphene-c for 10–200 ppm  $\text{NO}_2$  at room temperature; (f) electrical resistance of  $\text{CeO}_2$ /graphene-c versus time in the presence of 200 ppm  $\text{NO}_2$ . (G: graphene)

to 15 s for sample c ( $\text{CeO}_2$ {100}/graphene composites). It means that the  $\text{NO}_2$  molecules are more tightly bound on the  $\text{CeO}_2$ {100} surface than on the  $\text{CeO}_2$ {111} surface. As for  $\text{CeO}_2$ {111} facets/graphene composites, the sensing performance decreases in the following order: sample d > sample b > sample a, indicating that smaller  $\text{CeO}_2$  nano-particles increase the sensitivity of  $\text{CeO}_2$ /graphene composites. Fig. 7e and Fig. S5 (ESI<sup>†</sup>) show the response and recovery curves of  $\text{CeO}_2$ {100}/graphene (sample c) and  $\text{CeO}_2$ {111}/graphene sensors (sample a, b, d, and e) to  $\text{NO}_2$  gas of different concentrations at room temperature, respectively. Clearly, the resistance of the  $\text{CeO}_2$ {100}/graphene sensor decreases quickly during exposure to  $\text{NO}_2$  but reverts to the initial value fast upon exposure to air, indicating good response and recovery performance. Fig. 7f shows six response cycles of sample c during exposure to 200 ppm  $\text{NO}_2$ . The cycling curves nearly coincide with each other demonstrating stable sensing. The long-term stability of the sensor was evaluated through exposure to 200 ppm  $\text{NO}_2$  for a period of ten days. As shown in Fig. S6 (ESI<sup>†</sup>), the sensor exhibits good stability and the error is smaller than 5%.

Selectivity is also an important parameter for gas sensors. The responses of the  $\text{CeO}_2$ {100}/graphene composites to  $\text{NO}_2$ ,  $\text{NH}_3$ , and  $\text{O}_2$  have been measured. The results are shown in Fig. S7 (ESI<sup>†</sup>), and the responses to  $\text{NO}_2$ ,  $\text{O}_2$ ,  $\text{NH}_3$ ,  $\text{H}_2$ ,  $\text{CO}$ , and  $\text{C}_2\text{H}_2$  reported by others<sup>20</sup> are also summarized for comparison. It is clearly that the  $\text{CeO}_2$ {100}/graphene composites exhibit the highest sensitivity to  $\text{NO}_2$  among them, and thus an excellent selectivity toward  $\text{NO}_2$  at room temperature.

### 3.4 Gas-sensing mechanism

The gas-sensing properties of metal oxides are mainly dependent on the exposed facets.<sup>50</sup> As shown in Fig. S8 (ESI<sup>†</sup>), the exposed

{100} polar facets are the real reason for the good sensing performance of  $\text{CeO}_2$  nanoparticles to  $\text{NO}_2$  gas at room temperature. Structurally, each Ce atom is eightfold-coordinated with O atoms and each O atom is fourfold-coordinated with Ce atoms in the bulk, but they are unsaturated on the surface. Each O or Ce atom on the  $\text{CeO}_2$ {111} surface has one dangling bond and two on the {100} surface. Therefore, there are more dangling bonds on the {100} surface than on the {111} surface.<sup>51</sup> Consequently, the {100} facet is more active in adsorbing  $\text{NO}_2$  at room temperature. What is more, the {100} polar surface usually contains more oxygen vacancies than {111} nonpolar facets.<sup>51</sup> The oxygen vacancies can significantly enhance the adsorption of oxygen molecules, and electrons will transfer from the oxygen vacancies in  $\text{CeO}_2$  to the oxygen molecules, resulting in more oxygen species (especially  $\text{O}_2^-$ ). These oxygen species will react with  $\text{NO}_2$ , resulting in a more abrupt change in the conductivity of the sensor.<sup>52</sup> The graphene sheets in the composites improve the electrical conductivity of  $\text{CeO}_2$  and thus the gas sensing performance at room temperature substantially.<sup>9</sup> Furthermore, the overlapped and crumpled regions on graphene sheets provide a large surface area, thus facilitate the effective access of gases onto  $\text{CeO}_2$ /graphene composites.<sup>53</sup> Finally, a strong electronic interaction occurs at the interface of graphene and  $\text{CeO}_2$ , which will modulate the charge exchange during gas-sensing,<sup>54</sup> which is further confirmed using first-principles calculations.

### 3.5 Effect of facets on gas sensing performance

Fig. 8a and Fig. S9 (ESI<sup>†</sup>) show the work functions of the  $\text{CeO}_2$ {111} plane, the  $\text{CeO}_2$ {100} plane terminated with cerium ( $\text{CeO}_2$ {100}–Ce), the  $\text{CeO}_2$ {100} plane terminated with oxygen ( $\text{CeO}_2$ {100}–O), and graphene, respectively. The work function of the  $\text{CeO}_2$ {111} facet (6.14 eV) is larger than that of graphene (4.54 eV) and the electrons will transfer from graphene to the

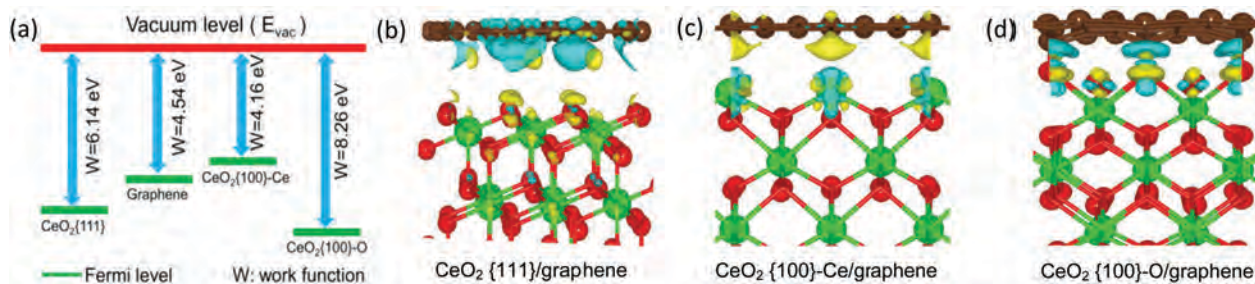


Fig. 8 (a) The work function of the  $\text{CeO}_2\{111\}$ , graphene,  $\text{CeO}_2\{100\}\text{-Ce}$ , and  $\text{CeO}_2\{100\}\text{-O}$  plane; 3D charge density differences of (b)  $\text{CeO}_2\{111\}$ /graphene hybrid structure, (c)  $\text{CeO}_2\{100\}\text{-Ce}$ /graphene hybrid structure and (d)  $\text{CeO}_2\{100\}\text{-O}$ /graphene hybrid structure. Yellow and cyan represent electron accumulation and depletion, respectively. The C, Ce, and O atoms are represented by brown, green and red balls, respectively.

$\text{CeO}_2\{111\}$  facet, resulting in electron depletion on graphene and electron accumulation on the  $\text{CeO}_2\{111\}$  plane, as evidenced from the 3D charge density difference in Fig. 8b. The energy bands of  $\text{CeO}_2\{111\}$  planes are bent down and hole depletion occurs at the interface (Fig. S10, ESI<sup>†</sup>). This leads to Schottky contact. Structurally, the  $\text{CeO}_2\{100\}$  plane can be either oxygen or cerium terminated.<sup>51</sup> The formation energy of the hybrid structure of graphene and the  $\text{CeO}_2\{100\}\text{-Ce}$  plane ( $-5.14$  eV) is lower than that with O termination ( $-4.39$  eV), and the  $\text{CeO}_2\{100\}\text{-Ce}$ /graphene hybrid structure is the stable one. Because the Fermi level of the  $\text{CeO}_2\{100\}\text{-Ce}$  plane is higher than that of graphene (Fig. 8a), the electrons are transferred from the  $\text{CeO}_2\{100\}\text{-Ce}$  surface to graphene, leading to electron depletion on the  $\text{CeO}_2\{100\}\text{-Ce}$  plane and electron accumulation on graphene, as confirmed by the 3D charge density difference in Fig. 8c. The polarity of the  $\text{CeO}_2\{100\}\text{-Ce}$  surface is compensated to some degree and the  $\text{CeO}_2\{100\}\text{-Ce}$  surface becomes more stable, which is different from the previously reported results.<sup>55,56</sup> Supposing that the  $\text{CeO}_2\{100\}$  plane is oxygen terminated, the Fermi level will be much lower than that of graphene (Fig. 8a), and the electrons will be transferred from graphene to  $\text{CeO}_2\{100\}$  (Fig. 8d), resulting in Schottky contact. This is not consistent with the measured  $I$ - $V$  curves (Fig. 10). So the heterogeneous interface of graphene and the  $\text{CeO}_2\{100\}\text{-Ce}$  plane is preferential.

To gain more insights into their electronic origin, the total and projected density of states of the  $\text{CeO}_2\{111\}$  plane, the  $\text{CeO}_2\{100\}\text{-Ce}$  plane,  $\text{CeO}_2\{111\}$  facets/graphene composites, and  $\text{CeO}_2\{100\}\text{-Ce}$  facets/graphene composites are calculated based on DFT. As shown in Fig. S11a and b (ESI<sup>†</sup>), the valence band maximum (VBM) and the conduction band minimum (CBM) of the  $\text{CeO}_2\{111\}$  plane are dominated by the O 2p and Ce 4f states, respectively. The Fermi level is close to the VBM (O 2p) characteristic of a p-type semiconductor and in agreement with the  $M$ - $S$  curve. The VBM and CBM of the  $\text{CeO}_2\{100\}\text{-Ce}$  facets are dominated by the O 2p and Ce 4f states, respectively and overlapped at the Fermi level, suggesting a metallic behavior (Fig. S11c and d, ESI<sup>†</sup>). When the hetero-interface is formed between graphene and the  $\text{CeO}_2\{111\}$  plane, the VBM and the CBM are mainly determined by the C 2p and Ce 4f states, respectively (Fig. 9a and b). The band gap between the occupied C 2p state and the unoccupied Ce 4f state is about 0.4 eV, which is larger than the thermal energy at room temperature (0.026 eV). Therefore, it is difficult for the electrons

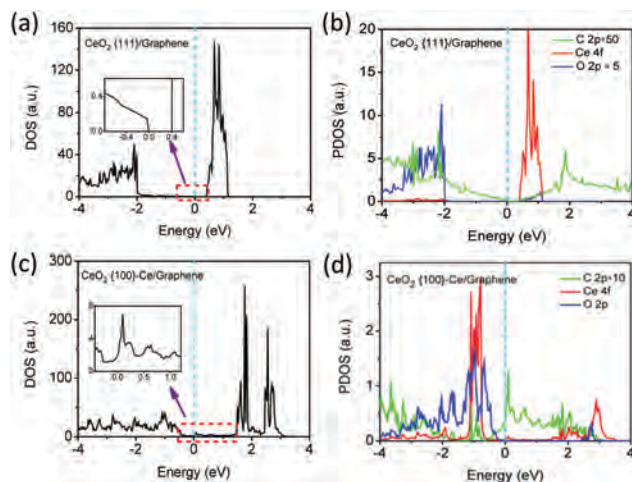


Fig. 9 The DOS and PDOS of (a and b)  $\text{CeO}_2\{111\}$ /graphene and (c and d)  $\text{CeO}_2\{100\}\text{-Ce}$ /graphene composites, respectively. The dashed lines indicate the Fermi level.

to transfer from the VBM of graphene to the CBM of  $\text{CeO}_2\{111\}$  at room temperature. However, as shown in Fig. 9c and d, the VBM and CBM of the  $\text{CeO}_2\{100\}\text{-Ce}$ /graphene hetero-interface are determined by the O 2p and C 2p states, respectively, and they overlap at the Fermi level, suggesting a metallic behavior. It is easy for the electrons to transfer from the VBM of the  $\text{CeO}_2\{100\}$  plane to the CBM of graphene at room temperature. Owing to their high carrier mobility in graphene, rapid charge exchange takes place between  $\text{NO}_2$  and  $\text{CeO}_2\{100\}$  facets/graphene composites. Consequently, the sensitivity, response and recovery are enhanced.

The current *versus* voltage ( $I$ - $V$ ) curves of the  $\text{CeO}_2$  nanoparticles, graphene sheets,  $\text{CeO}_2\{111\}$  facets/graphene composites, and  $\text{CeO}_2\{100\}\text{-Ce}$  facets/graphene composites are shown in Fig. 10. The slope of  $\text{CeO}_2$ /graphene composites is larger than that of pure  $\text{CeO}_2$  nanoparticles as a result of reduced resistance. The  $I$ - $V$  curve of  $\text{CeO}_2\{111\}$  facets/graphene composites is non-linear, revealing the Schottky contact, but that of  $\text{CeO}_2\{100\}\text{-Ce}$  facets/graphene composites is linear, suggesting the Ohmic contact. This confirms the above theoretical prediction, and the hybrid structures of graphene and  $\text{CeO}_2$  nanocubes surrounded by the  $\{100\}$  facets are crucial to the enhanced  $\text{NO}_2$  sensing performance.

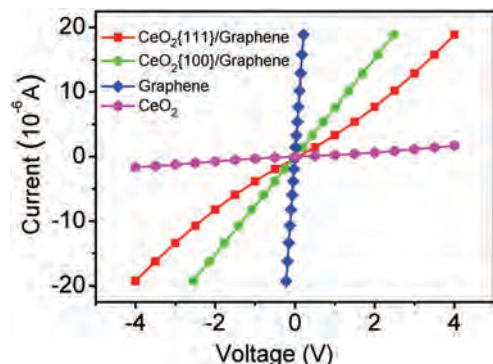


Fig. 10 C–V curves of CeO<sub>2</sub> nanoparticles, graphene sheets, CeO<sub>2</sub>{111}/graphene, and CeO<sub>2</sub>{100}/graphene nanocomposites.

## 4. Conclusion

The hydrothermal method is adopted to fabricate CeO<sub>2</sub>/graphene nanocomposites. It is found that the shape and facets of the CeO<sub>2</sub> nanocrystals can be adjusted simply by changing the volume ratio of ethylene glycol to deionized water in the solution. When the volume ratio of ethylene glycol to deionized water is 1:1, cubic CeO<sub>2</sub> nanoparticles surrounded by high-energy {100} facets are obtained, otherwise CeO<sub>2</sub> nanograins enclosed by the {111} facets are obtained, which is closely related to the changes in the energetics on surface. As compared to CeO<sub>2</sub>{111}/graphene composites, the CeO<sub>2</sub>{100}/graphene nanocomposites exhibit an excellent response to NO<sub>2</sub>, for example, the sensitivity during exposure to 10 ppm NO<sub>2</sub> is increased by 1400% as compared to pure CeO<sub>2</sub>. First-principles calculations reveal that electrons transfer from graphene sheets to the CeO<sub>2</sub>{111} planes producing electron depletion on graphene at the interface. In contrast, the electrons move from the CeO<sub>2</sub>{100} planes to graphene sheets and, the VBM and the CBM overlap at the Fermi level producing a metallic behavior. Therefore, NO<sub>2</sub> can easily capture electrons from the CeO<sub>2</sub>{100} facets/graphene nanocomposites at room temperature and the sensitivity of NO<sub>2</sub> detection is significantly enhanced. The results present us a clear physical picture for the tunable facets and morphology as well as the enhanced performances of metal oxide nanoparticles/graphene composites.

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