



## Comparison among different models of track growth and experimental data

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### Abstract

The lengths of major axes of the alpha-particle track openings in CR-39 detectors were calculated by using four different models of track growth. Special computer programs were prepared for each of the models under comparison. Some differences were found among the results from the various models considering the wide ranges of incident angles, energies and removed layers during chemical etching. Systematical comparisons with experimental data for various incident angles and energies up to 5 MeV were also performed.

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### 1. Introduction

Development of tracks in materials during etching has attracted much attention for a long time (e.g., Henke and Benton, 1971; Paretzke et al., 1973; Somogyi and Szalay, 1973; Somogyi, 1980; Fromm et al., 1988; Fews and Henshaw, 1982; Ditlov, 1995; Nikezic and Kostic, 1997; Nikezic and Yu, 2003a, 2004). Formation of a track as a result of the incidence of an ion is due to the simultaneous acting of the etching solution with two etching rates, i.e., the variable rate along the path of the incident particle (called the track etch rate  $V_t$ ) and a constant rate in all other regions (called the bulk etch rate  $V_b$ ). All models of track growth are based on these two parameters which were introduced by Fleischer et al. (1975).

During etching, the track development passes through several phases. At the beginning of etching, the track is

conical with a sharp tip. If etching passes the end point of the particle trajectory, the track tip becomes rounded. Therefore, these two main phases can be distinguished as the sharp conical one and the rounded one. If etching is prolonged, the track can become spherical (Durrani and Bull, 1987). The track opening also passes through different phases, which were analyzed in details by Somogyi and Szalay (1973). The track opening is circular when the incident angle is  $90^\circ$ . For oblique incidence, the track opening is elliptical, elliptical + circular or circular depending on the etching condition, incident energy and incident angle (for a constant  $V_t$ ), or semi-elliptical or an even more complex “egg-like” geometrical figure (for a varying  $V_t$ ).

In the present paper, we compared four different models of track growth given by Somogyi and Szalay (1973), Fews and Henshaw (1982), Fromm et al. (1988) and Nikezic and Yu (2003a,b). All of these four models enabled calculations of the length of the major axis of the alpha-particle track openings, which was the quantity compared in the present work. Some results of our model have already been published (Nikezic and Yu, 2002, 2003b), but without comparison with experimental data.

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Many other works have made partial comparisons between models of the track growth and experiment data. For example, Dörschel et al. (2003a) compared the calculated and experimentally measured profiles of tracks from alpha particles and  $^7\text{Li}$  ions in the CR-39 detector. Comparisons were made for one incident energy (5.9 MeV for alpha particles and 7.95 MeV for  $^7\text{Li}$  ions) and one incident angle ( $40^\circ$  with respect to the normal to the detector surface) for two removed layers. Good agreement between the experimental and calculated profiles was found in all examined cases. Calculated and measured profiles for protons (with energy of 1.27 MeV), alpha particles (with energy 4 MeV),  $^7\text{Li}$  ions (with energy of 4.82 MeV) and  $^{12}\text{C}$  ions (with energy of 14.8 MeV) were also compared by Dörschel et al. (2003b) and again good agreements were found for two removed layers. The calculations in both these two papers were based on the model of Fromm et al. (1988).

However, there has not been systematic comparison for a whole range of incident angles and energies of alpha particles. In the present paper, systematical comparison between models and experimental results is given for whole ranges of incident angles and energies.

## 2. Calculation method

Four different computer programs were separately prepared for the four track growth models mentioned above. Details of the models can be found in the corresponding references so they are not repeated here. Comparisons were made for alpha-particle tracks in the CR-39 detector. The input parameters,  $V_t$  and  $V_b$ , given below were adopted for all models. The bulk etch rate was taken as  $V_b = 1.2 \mu\text{m/h}$  (Ho et al., 2003). The variable track etch rate  $V_t$  was taken from Yu et al. (2005) as

$$V = 1 + e^{-a_1 R' + b} - e^{-a_2 R' + b}, \quad (1)$$

where  $R'$  was the residual range of alpha particles, and the constants were  $a_1 = 0.06082 \mu\text{m}^{-1}$ ,  $a_2 = 0.8055 \mu\text{m}^{-1}$  and  $b = 1.119$ . This function was obtained by measuring the depths of sub-micron tracks in a CR-39 detector generated by alpha particles with energies from 1 to 4.5 MeV, from replicas using atomic force microscopy (Yu et al., 2005). Therefore, this function can be useful for energies up to  $\sim 5$  MeV (or  $R' \approx 28.5 \mu\text{m}$  in the CR-39 detector).

## 3. Experiment

The CR-39 detectors used in the present work were purchased from Page Mouldings (Pershire) Limited (Worcestershire, England). The cut detectors were systematically irradiated with alpha particles with energies ranging from 1 to 5 MeV, with steps of 1 MeV. An  $^{241}\text{Am}$  alpha-particle source (main initial energy  $E_0 = 5.48$  MeV) was employed,

and the stopping medium between the source and the detectors was air. The alpha-particle energy incident on a piece of detector with a distance from the source was determined with an alpha spectroscopy system (ORTEC Model 5030). The incident angle (with respect to the detector surface) was also varied from  $40^\circ$  up to  $90^\circ$ , with steps of  $10^\circ$ . Irradiation with angles smaller than  $40^\circ$  will be close to the critical angle. Fromm et al. (1988) found that there were many cases for which a great number of tracks were missing when the incident angle was smaller than  $40^\circ$ .

At least three pieces of CR-39 detectors were irradiated with each combination of incident energy and angle. The irradiated detectors were etched in an aqueous 6.25 N NaOH solution at a temperature of  $70^\circ$  for 15 h. Under these etching conditions, the bulk etch rate has been found to be about  $1.2 \mu\text{m/h}$  (Ho et al., 2003). The lengths of the major axes were measured by the Image Processing and Analysis System (Leica Imaging Systems QW in standard V2.3) with a magnification of 1000. More than 20 tracks were measured and analyzed for each combination of incident energy and angle.

## 4. Results

Computations of the major axes have been performed for incident angles from  $40^\circ$  to  $90^\circ$  with steps of  $10^\circ$ , and for four different etching times, viz., 5, 10, 15 and 20 h. Since  $V_b = 1.2 \mu\text{m/h}$  was taken, the removed layers used in calculations were 6, 12, 18 and  $24 \mu\text{m}$ .

The results are given in Figs. 1–6. The same symbols are used in all figures: open circles for results from the model of Fromm et al., pluses for the Fewes–Henshaw model, triangles for the Somogyi–Szalay model and solid lines for the Nikežic–Yu model. The thickness of the removed layers is given close to the corresponding groups of data.

Fig. 1 shows the results for the incident angle of  $90^\circ$ , where the major axis of the track opening is given as a function of the incident energy, while the etching condition is used as a parameter. One can observe good agreement among different models. Agreement is particularly good for smaller incident energies. The results from the models of Fewes–Henshaw and of Fromm et al. are very close at all energies, with differences between 0.5% and 1%. The Somogyi–Szalay model agrees very well with others, but deviates slightly at larger energies. The largest difference occurs at the largest considered energy of 7 MeV and amounts up to  $1.5 \mu\text{m}$ , which is 7–20% larger in comparison to the Fewes–Henshaw model. The Nikežic–Yu model falls between the models of Fewes–Henshaw and Fromm et al. on one side and the Somogyi–Szalay model on the other side. In fact, the same pattern appears in all investigated cases.

The results for the incident angle of  $80^\circ$  are given in Fig. 2. There are no qualitative differences from the results shown in Fig. 1. The difference, up to  $1.5 \mu\text{m}$  (8–24%), between the models of Somogyi–Szalay and Fewes–Henshaw

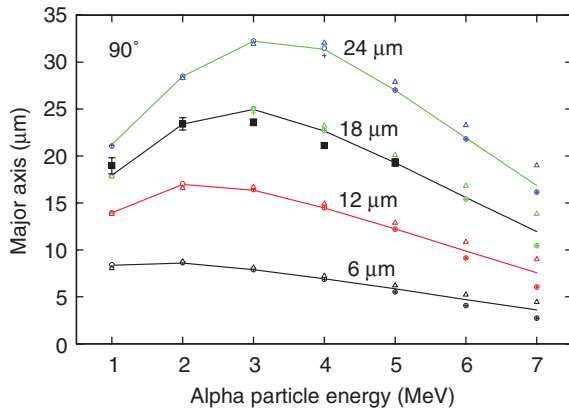


Fig. 1. Comparison of the lengths of the major axis of tracks from alpha particles with different incident energies calculated by four different track growth models for the incident angle  $90^\circ$  with experimental data. Open circles: results from the model of Fromm et al. (1988); Pluses: results from the model of Fewes and Henshaw (1982); Triangles: results from the model of Somogyi and Szalay (1973); Solid lines: results from the model of Nikezic and Yu (2003 a, b); Solid squares: experimental data with error bars showing one standard deviations. The removed layers are 6, 12, 18 and  $24\ \mu\text{m}$  as shown.

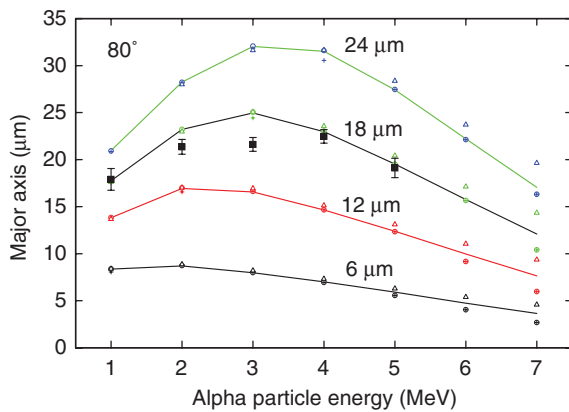


Fig. 2. Same as Fig. 1, but for the incident angle  $80^\circ$ .

is again largest at 7 MeV. The Nikezic–Yu model also falls between those two. The differences become smaller and all models converge for smaller energies.

For smaller incident angles (see Figs. 3–5), the discrepancies among the models increase, but not significantly. For the incident angle of  $40^\circ$  (Fig. 6), some data are missing because the programs do not produce results for some of the models. The poorest agreement for the incident angle of  $40^\circ$  (Fig. 6) occurs at the right end of the curves (for incident energies of 5, 6 and 7 MeV) where the Fewes–Henshaw and Fromm et al. models gave smaller values than the Somogyi–Szalay and Nikezic–Yu models. This is consequence of different

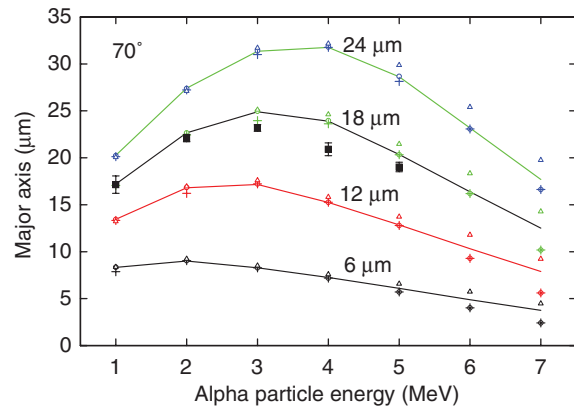


Fig. 3. Same as Fig. 1, but for the incident angle  $70^\circ$ .

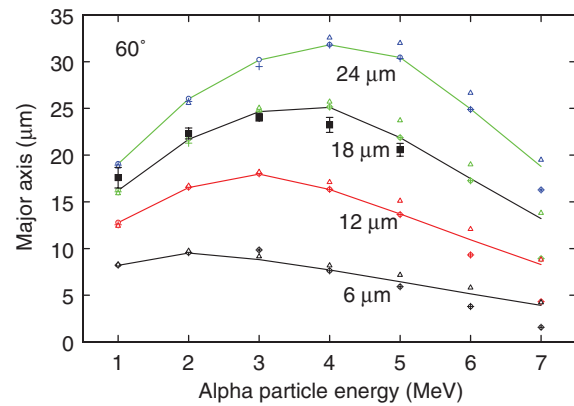


Fig. 4. Same as Fig. 1, but for the incident angle  $60^\circ$ .

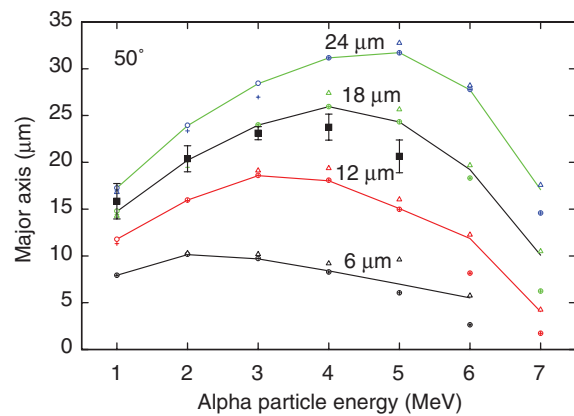


Fig. 5. Same as Fig. 1, but for the incident angle  $50^\circ$ .

treatments of the tracks whose developments do not start on the initial surface of the detector. In the lower energy region, agreement is also very good for the incident angle of  $40^\circ$ .

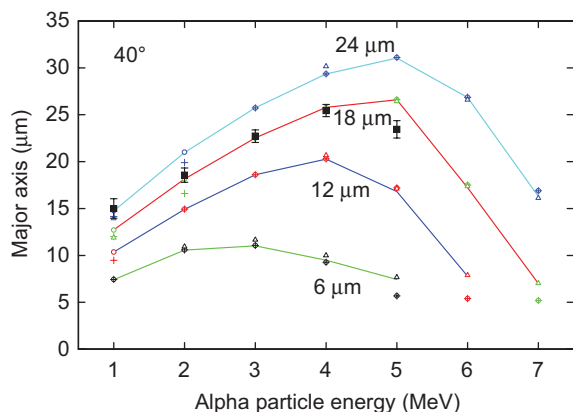


Fig. 6. Same as Fig. 1, but for the incident angle 40°.

Comparisons with experimental results are also given in Figs. 1–6. Experimental data are represented by solid squares and these correspond to a removed layer of 18 μm. There are no observed systematic discrepancies between the calculated and the experimental values. In all examined cases, the discrepancies are smaller than 20%; and in some cases these are even below 5%.

## 5. Conclusions

Comparison of track growth models through the prediction of the lengths of the major axis has shown small differences among the four examined models. The models from Fewes–Henshaw and Fromm et al. agree within less than 1% and the results overlap for all the examined incident energies, incident angles and removed layers. The Somogyi–Szalay model gives consistently larger values for the major axis for larger incident energies while the Nikezic–Yu model consistently falls between the two extreme cases. The general conclusion is that these models give practically the same values for the major axis of the track opening for wide ranges of incident angles and removed layer, for incident alpha energies up to 5 MeV. For incident alpha energies larger than 5 MeV, discrepancies among models become apparent. These are likely due to the inherent different approaches of the different models, and will be the subject of a separate investigation in the future. Finally, the values calculated using different models of track growth agree with the experimental data to within 20%.

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